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ANCHORED NASH INEQUALITIES AND HEAT KERNEL BOUNDS FOR STATIC AND DYNAMIC DEGENERATE ENVIRONMENTS

JEAN-CHRISTOPHE MOURRAT, FELIX OTTO

ABSTRACT. We introduce anchored versions of the Nash inequality. They allow to control the L^2 norm of a function by Dirichlet forms that are not uniformly elliptic. We then use them to provide heat kernel upper bounds for diffusions in degenerate static and dynamic random environments. As an example, we apply our results to the case of a random walk with degenerate jump rates that depend on an underlying exclusion process at equilibrium.

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KEYWORDS: Nash inequality, heat kernel, diffusion in dynamic random medium.

1. INTRODUCTION

In the study of reversible random walks in degenerate static random environments, heat kernel estimates have played a key role. Consider a random walk evolving on a super-critical (infinite) percolation cluster. As was pointed out in [39], the problem of showing that the walk satisfies a central limit theorem for almost every realisation of the environment (i.e. a “quenched” CLT) is reduced to showing a spatially averaged (instead of uniform) sublinear control of the corrector if the heat kernel is known to satisfy a diffusive upper bound. For the case of percolation, such a bound was obtained in [33, 7]. This bound indeed became a primary ingredient in the subsequent proofs that the random walk on percolation clusters and other degenerate i.i.d. environments satisfies a quenched central limit theorem [39, 9, 32, 15, 31, 2]. (A notable exception is a geometric argument introduced in [9], which in two dimensions does not use any a priori heat kernel bounds.) We refer to [12] (in particular, Section 4.3 therein) and [28] for reviews. Related ideas of regularity theory lead to a proof of the local central limit theorem in [8].

The approach to heat kernel bounds in [33] is based on studying isoperimetric inequalities on boxes, with an effective dimension that is slowly tuned up to d as the size of the box increases. In [7], upper and lower bounds are proved (including optimal off-diagonal control) using a fine decomposition of the space into good and bad boxes, where a box is said to be good if the Poincaré inequality (with the standard scaling) holds in that box. The case of i.i.d. conductances with a polynomial tail near zero is investigated in [17] using a combination of probabilistic arguments and Harnack inequalities.

One can also consider environments that do not possess strong decorrelation properties at large distances. Assuming a certain moment condition (from above and below) on the conductances, it was shown in [3] that the random walk satisfies a quenched central limit theorem. The proof relies on weighted Sobolev inequalities, and on Moser’s iteration scheme. Harnack inequalities and a local CLT were then established in [4], and Gaussian heat kernel bounds in [5]. Related ideas were developed in [22] to prove a quenched invariance principle for Brownian motion subject to an incompressible drift. In a different direction, diffusions on clusters of

correlated percolation models such as random interacements are studied in [37] via isoperimetric inequalities.

Some assumptions on the law of the conductivities are necessary for diffusive heat kernel bounds to hold. Indeed, even in the setting of i.i.d. conductances, where the quenched CLT is known to hold in full generality [15, 31, 2], diffusive heat kernel upper bounds will not hold if the conductivities are allowed to take very small values [10]. The basic mechanism is that a portion of space surrounded by a region of low conductivity acts as a trap for the random walk. This phenomenon was analysed with precision in [16, 13, 14].

The aim of this work is to introduce a new method to prove diffusive heat kernel bounds for degenerate environments. In a celebrated work [35], Nash introduced the inequality that now bears his name for this purpose, for uniformly elliptic differential operators. This was then extended to the setting of Markov chains in [19]. Our first main result, obtained in Section 2, is an anchored version of the Nash inequality. The standard Nash inequality gives a control of the L^2 norm of a function f in terms of $\|\nabla f\|_2$ and $\|f\|_1$. We show that one can control the L^2 norm of f in terms of $\|w\nabla f\|_2$, $\|f\|_1$ and $\| |x|^{p/2} f \|_2$, where w is a weight function that can for instance be taken to be the square root of the conductance field (in particular, it is not bounded away from 0 in cases of interest). Note that the translation invariance of the standard Nash inequality is broken by the term $\| |x|^{p/2} f \|_2$.

We then show in Section 3 that despite this extra term $\| |x|^{p/2} f \|_2$, one can obtain diffusive heat kernel upper bounds for static degenerate environments satisfying a certain flexible assumption. In particular, this assumption covers bounded stationary conductances with inverse having a moment of order $p > d$. When the conductances are independent, a moment of order $p > 1/4$ suffices. We thus recover certain results of [4, 17]. We can also cover cases with vanishing conductances, see Remark 3.8. As a by-product of our method, we also obtain an estimate on the behaviour of the heat kernel away from the diagonal. Stronger information can be obtained from the on-diagonal bounds we prove using the result of [24] recalled here in Theorem 3.5, which is itself inspired by [26] (see also [21, 20]).

Our main motivation is to derive heat kernel bounds for *dynamic* degenerate random environments. We are not aware of previous results in this context, although quenched central limit theorems have been proved in some cases (see in particular [23]). The dynamics of the environment can facilitate anomalous behaviour of the heat kernel [18]. As a simple example, one can construct a dynamic environment such that the probability of return to the origin decays from 1 to almost zero, and then climb back up to almost 1/4. Indeed, take z a neighbour of the origin, and consider that initially, the edge connecting the origin to z is open, and these two sites are disconnected from the rest of the graph. The probability for the walk to be on z goes to 1/2 as time passes. Then change the environment so that the site z is isolated from the rest, and otherwise every edge is open. The probability for the walk to be at the origin decays to 0. Then move back to the initial situation with the origin and z connected, and disconnected from the rest: the probability for the walk to be at the origin moves back to being almost 1/4.

Let us denote by $p_{s,t}(x,y)$ the probability for the walk started at x at time s to be at y at time $t \geq s$. It is more convenient to aim first for bounds on

$$\mathcal{E}_t := \sum_{x \in \mathbb{Z}^d} (p_{0,t}(0,x))^2,$$

in particular because, as opposed to $p_{0,t}(0,0)$, this quantity is monotone in t . In Section 4, we provide a general criterion that allows to prove diffusive bounds of

the form

$$(1.1) \quad \mathcal{E}_t \leq \frac{\mathcal{X}}{t^{d/2}},$$

where \mathcal{X} is a random variable. In order to illustrate the relevance of our criterion, we study in Section 5 an example of a random walk whose (symmetric) jump rates $(a_t(e))$ are a local function of an underlying simple exclusion process at equilibrium. We show that in this case, (1.1) holds with a random variable \mathcal{X} that has finite moments of every order, regardless of the density of particles in the exclusion process, or of the specific definition of the jump rates. (In particular, it may be that at any given time, the set of sites that are reachable by the random walk is finite.) We also show that for every $r \geq 1$, $\sup_t \mathbb{E}[(t^{d/2} p_{0,t}(0,0))^r]$ is finite, and that for every $\varepsilon > 0$, $\mathbb{E}[(\sup_t t^{d/2-\varepsilon} p_{0,t}(0,0))^r]$ is also finite. In order for these results to hold, we need only little information on the underlying dynamics. Roughly speaking, we only use the fact that the jump rates are very unlikely to vanish over an extended period of time. In particular, we do not use the fact that the simple exclusion process has a polynomial mixing rate of order $t^{-d/2}$ (as was shown in [11, 27]).

An immediate consequence of the diffusive heat kernel bound is the transience of the random walk in dimension $d \geq 3$. This extends [36, Theorem 1.1 (ii)] to conservative (and therefore much less mixing) environments.

1.1. Notation. Throughout this paper, we focus on the discrete space \mathbb{Z}^d , $d \geq 1$. We say that $x, y \in \mathbb{Z}^d$ are *neighbours*, and write $x \sim y$, if $|x - y| = 1$, where $|\cdot|$ is the Euclidean distance. This turns \mathbb{Z}^d into a graph, and we denote by \mathbb{B} the associated set of (unoriented) edges. For any positive integer r , we denote by B_r the box $\{-r, \dots, r\}^d$, and \mathbb{B}_r the set of edges with both end-points in B_r . For any $p \in [1, \infty]$, we denote by $\|\cdot\|_{L_r^p}$ the norm of the space $L^p(B_r)$ or $L^p(\mathbb{B}_r)$ (which of the two will be clear from the context), i.e.

$$\|f\|_{L_r^p} = \left(\sum_{x \in B_r} |f(x)|^p \right)^{1/p} \quad \text{or} \quad \|f\|_{L_r^p} = \left(\sum_{e \in \mathbb{B}_r} |f(e)|^p \right)^{1/p},$$

with the usual interpretation as a supremum if $p = \infty$. We also write

$$\|f\|_p = \left(\sum_{x \in \mathbb{Z}^d} |f(x)|^p \right)^{1/p} \quad \text{or} \quad \|f\|_p = \left(\sum_{e \in \mathbb{B}} |f(e)|^p \right)^{1/p}.$$

Finally, we introduce the shorthand notation $|\cdot|_* := |\cdot| \vee 1$.

2. ANCHORED NASH INEQUALITY

The aim of this section is to prove the following theorem.

Theorem 2.1 (Anchored Nash inequality). *Let $p \in (d, +\infty)$, $q \in (d, +\infty]$, and $\theta \in [\theta_c, 1]$, where $\theta_c \in [0, 1)$ is defined by*

$$(2.1) \quad \frac{1}{\theta_c} = 1 + \frac{dp + 2p}{dp + 2d} \left(\frac{q}{d} - 1 \right).$$

Define $\alpha, \beta, \gamma \in [0, 1)$ by

$$(2.2) \quad \alpha = (1 - \theta) \frac{d}{d+2} + \theta \frac{p}{p+2}, \quad \beta = (1 - \theta) \frac{2}{d+2}, \quad \gamma = \theta \frac{2}{p+2}.$$

There exists $C < \infty$ such that for any functions $f : \mathbb{Z}^d \rightarrow \mathbb{R}$ and $w : \mathbb{B} \rightarrow \mathbb{R}_+$,

$$(2.3) \quad \|f\|_2 \leq C (\mathcal{M}_q \|w \nabla f\|_2)^\alpha \|f\|_1^\beta \| |x|_*^{p/2} f \|_2^\gamma,$$

where

$$(2.4) \quad \mathcal{M}_q^q = \mathcal{M}_q^q(w) = \sup_{r \in \mathbb{N}} \frac{1}{|\mathbb{B}_r|} \sum_{e \in \mathbb{B}_r} w^{-q}(e).$$

Remark 2.2. By setting $q = +\infty$ and $\theta = 0$, we recover the usual Nash inequality, provided that \mathcal{M}_∞ is finite (i.e. that w^{-1} is bounded). For $\theta > 0$ (and thus $\gamma > 0$), the inequality is no longer translation invariant. Indeed, a crucial aspect of the inequality is that in the definition of \mathcal{M}_q , the supremum is taken only over boxes that are centred around the origin. The price one has to pay for this is the appearance of the extra term $\| |x|_*^{p/2} u \|_2^\gamma$, which is of course not translation invariant either.

Remark 2.3. Clearly, $\alpha + \beta + \gamma = 1$, as must be the case by homogeneity.

Remark 2.4. Our proof applies as well to the setting of continuous space. Considerations of scale invariance then show that the exponents α , β and γ have to be of the form (2.2) for some θ .

Remark 2.5. As far as our applications are concerned, the precise value of θ_c is unimportant, the only crucial point being that $\theta_c < 1$.

Before we turn to the proof of Theorem 2.1, we recall some classical functional inequalities for the reader's convenience.

Theorem 2.6 (Poincaré-Sobolev inequalities). *Let $1 < p < d$ and let p^* be such that $d/p^* = d/p - 1$. There exists a constant $C < \infty$ such that for every r and $f : B_r \rightarrow \mathbb{R}$,*

$$(2.5) \quad \|f - \bar{f}_r\|_{L^{p^*}} \leq C \|\nabla f\|_{L^p},$$

where we write

$$\bar{f}_r = \frac{1}{|B_r|} \sum_{x \in B_r} f(x).$$

Proof. For every $x, y \in B_r$, we give ourselves a path $\gamma_{x,y}$ of nearest-neighbour points of \mathbb{Z}^d from x to y that stays as close as possible to the segment joining x and y . To be precise, by symmetry we may assume that $y \geq x$ (i.e. each coordinate of y is larger than the corresponding coordinate of x). In this case, we construct a path $x_0 = x, x_1, \dots, x_l = y$ as follows. We write $I_{x,y}$ for the line segment joining x and y in \mathbb{R}^d . Assuming that x_0, \dots, x_j are already built, we define

$$i = \arg \min \{ \text{dist}(x_j + \mathbf{e}_{i'}, I_{x,y}), 1 \leq i' \leq d \},$$

(take the smallest one in case of ties), and put $x_{j+1} = x_j + \mathbf{e}_i$. Such a path stays at distance at most \sqrt{d} from $I_{x,y}$, and has no self-intersections. We have

$$|f(y) - f(x)| \leq \sum_{e \in \gamma_{x,y}} |\nabla f|(e),$$

with obvious notation. Since

$$(2.6) \quad f(x) - \bar{f}_r = \frac{1}{|B_r|} \sum_{y \in B_r} (f(x) - f(y)),$$

the triangle inequality yields

$$|f(x) - \bar{f}_r| \leq \frac{1}{|B_r|} \sum_{\substack{y \in B_r \\ e \in \gamma_{x,y}}} |\nabla f|(e).$$

One can check that there exists a constant C (independent of r) such that for every $x \in B_r$ and $e \in \mathbb{B}$,

$$|\{y \in B_r : e \in \gamma_{x,y}\}| \leq Cr \left(\frac{r}{1 + |x - \underline{e}|} \right)^{d-1}.$$

As a consequence,

$$(2.7) \quad |f(x) - \bar{f}_r| \lesssim \sum_{e \in \mathbb{B}_r} (1 + |x - \underline{e}|)^{-(d-1)} |\nabla f|(e).$$

The desired result now follows from the Hardy-Littlewood-Sobolev inequality, see e.g. [6, Theorems 1.5 and 1.7]. \square

Remark 2.7. Theorem 2.6 can be generalised in the following way. Under the same conditions, if $g : B_r \rightarrow \mathbb{R}$ is such that $\bar{g}_r = 1$, then

$$\|f - \overline{(fg)}_r\|_{L_r^{p^*}} \leq C \|g\|_{L_r^\infty} \|\nabla f\|_{L_r^p}.$$

Indeed, it suffices to observe that

$$\left| \overline{(fg)}_r - \bar{f}_r \right| = \left| \frac{1}{|B_r|} \sum_{x \in B_r} g(x) (f(x) - \bar{f}_r) \right| \leq \|g\|_{L_r^\infty} \frac{1}{|B_r|^{1/p^*}} \|f - \bar{f}_r\|_{L_r^{p^*}},$$

and to use the triangle inequality and Theorem 2.6 (with $\|g\|_{L_r^\infty} \geq 1$).

In dimension 2, we will also need to use inequality (2.5) in the critical case $p^* = 2$, $p = 1$. The following theorem gives the result in any dimension.

Theorem 2.8 (Isoperimetric inequality). *Let $p^* = d/(d-1)$. There exists $C < \infty$ such that for every r and $f : B_r \rightarrow \mathbb{R}$,*

$$\|f - \bar{f}_r\|_{L_r^{p^*}} \leq C \|\nabla f\|_{L_r^1}.$$

Proof. Let m be a median of f . It suffices to show that

$$(2.8) \quad \|f - m\|_{L_r^{p^*}} \leq C \|\nabla f\|_{L_r^1}.$$

Indeed, if $m > \bar{f}_r$, then by the defining property of the median,

$$\frac{|B_r|}{2} \leq |\{x \in B_r : f(x) \geq m\}| \leq \frac{1}{(m - \bar{f}_r)^{p^*}} \|f - \bar{f}_r\|_{L_r^{p^*}}^{p^*}.$$

If $m < \bar{f}_r$, the symmetric argument shows that in every case,

$$|m - \bar{f}_r| \leq \frac{2}{|B_r|^{1/p^*}} \|f - \bar{f}_r\|_{L_r^{p^*}}.$$

By the triangle inequality, it thus suffices to show (2.8). Without loss of generality, we may assume that $m = 0$. Let $g = f \vee 0$. Since $|\nabla g| \leq |\nabla f|$, it suffices to show that

$$(2.9) \quad \|g\|_{L_r^{p^*}} \leq C \|\nabla g\|_{L_r^1} \quad (g \text{ s.t. } |\{x \in B_r : g(x) = 0\}| \geq |B_r|/2).$$

For $t \geq 0$, let $\mathcal{L}_t = \{x \in B_r : g(x) > t\}$. Since $g = \int_0^\infty \mathbf{1}_{\mathcal{L}_t} dt$, we have by Minkowski's inequality

$$\|g\|_{L_r^{p^*}} \leq \int_0^\infty |\mathcal{L}_t|^{1/p^*} dt.$$

The isoperimetric inequality [38, Theorem 3.3.15] ensures that there exists $C < \infty$ such that for every $A \subseteq B_r$ with $|A| \leq |B_r|/2$, we have $|A|^{(d-1)/d} \leq C|\partial A|$, where

$$\partial A = \{e = \{x, y\} \in \mathbb{B}_r : x \in A \text{ and } y \notin A\}.$$

Hence,

$$\begin{aligned} \|g\|_{L_r^{p^*}} &\leq C \int_0^\infty |\partial_t \mathcal{L}_t| dt \\ &= C \int_0^\infty \sum_{x \sim y: g(x) < t \leq g(y)} 1 dt \\ &= C \sum_{x \sim y: g(x) < g(y)} (g(y) - g(x)) = C \|\nabla g\|_{L_r^1}, \end{aligned}$$

which concludes the proof. \square

Remark 2.9. Under the conditions of Theorem 2.8, if $g : B_r \rightarrow \mathbb{R}$ is such that $\bar{g}_r = 1$, then

$$\|f - \overline{(fg)}_r\|_{L_r^{p^*}} \leq C \|g\|_{L_r^\infty} \|\nabla f\|_{L_r^1}.$$

Indeed, the proof is identical to that of Remark 2.7.

Proof of Theorem 2.1. The proof will be divided into two steps. We show the result for $\theta = 1$ in the first step, and for $\theta = \theta_c$ in the second step. The full result then follows by interpolation: it suffices to take $\nu \in [0, 1]$, write $\|f\|_2 = \|f\|_2^\nu \|f\|_2^{1-\nu}$ and apply the result with $\theta = 1$ and $\theta = \theta_c$ on each of the two terms.

Step 1. We show that the theorem holds for $\theta = 1$. Let p' be such that

$$(2.10) \quad \frac{1}{p'} = \frac{1}{d} + \frac{1}{2}.$$

Let R be an even positive integer. We define a function $g : B_R \rightarrow \mathbb{R}_+$ that is 0 in $B_{R/2}$, and is constant on the complementary ‘‘annulus’’ $A_R := B_R \setminus B_{R/2}$. The constant is specified by imposing $\bar{g}_R = 1$. We use the Poincaré-Sobolev inequality in the form given by Remark 2.7 in dimension $d \geq 3$, and the isoperimetric inequality in the form given by Remark 2.9 when $d = 2$, to derive

$$\|f - \overline{(fg)}_R\|_{L_R^2} \lesssim \|\nabla f\|_{L_R^{p'}}.$$

In view of the definition of g , we have

$$\|\overline{(fg)}_R\|_{L_R^2} \lesssim |B_R|^{-1/2} \|f\|_{L^1(A_R)} \leq \|f\|_{L^2(A_R)} \leq R^{-p/2} \| |x|_*^{p/2} f \|_{L_R^2},$$

where we used Jensen’s inequality for the second inequality. By Hölder’s inequality (and recalling (2.10)), we also have

$$\|\nabla f\|_{L_R^{p'}} \leq \|w^{-1}\|_{L_R^d} \|w \nabla f\|_{L_R^2}.$$

By Jensen’s inequality (and since $q \geq d$),

$$\|w^{-1}\|_{L_R^d} \leq |B_R|^{1/d-1/q} \|w^{-1}\|_{L_R^q} \leq |B_R|^{1/d} \mathcal{M}_q,$$

with \mathcal{M}_q as in (2.4). We thus get

$$\|f\|_{L_R^2} \lesssim |B_R|^{1/d} \mathcal{M}_q \|w \nabla f\|_{L_R^2} + R^{-p/2} \| |x|_*^{p/2} f \|_{L_R^2}.$$

Clearly,

$$(2.11) \quad \|f\|_{L^2(\mathbb{Z}^d \setminus B_R)} \leq R^{-p/2} \| |x|_*^{p/2} f \|_{L^2},$$

so we established

$$(2.12) \quad \|f\|_{L^2} \lesssim |B_R|^{1/d} \mathcal{M}_q \|w \nabla f\|_{L^2} + R^{-p/2} \| |x|_*^{p/2} f \|_{L^2},$$

uniformly over R even positive integer. It suffices to change the implicit multiplicative constant in (2.12) to ensure that the inequality

$$(2.13) \quad \|f\|_{L^2} \lesssim R \mathcal{M}_q \|w \nabla f\|_{L^2} + R^{-p/2} \| |x|_*^{p/2} f \|_{L^2}$$

remains valid for every real $R \geq 1$. Since $\|f\|_{L^2} \leq \| |x|_*^{p/2} f \|_{L^2}$, the inequality extends to all $R > 0$. Optimizing over R , we obtain

$$\|f\|_{L^2} \lesssim (\mathcal{M}_q \|w\nabla f\|_{L^2})^{p/(p+2)} \| |x|_*^{p/2} f \|_{L^2}^{2/(p+2)},$$

as desired.

Step 2. We now show that Theorem 2.1 holds for $\theta = \theta_c$.

Let Q be a box of size $r \in \mathbb{N}$. We use again p' as in (2.10). By the Poincaré-Sobolev (or isoperimetric) inequality,

$$\|f - \bar{f}_Q\|_{L^2(Q)} \lesssim \|\nabla f\|_{L^{p'}(Q)},$$

where we denote by \bar{f}_Q the average of f over the box Q . By (2.10) and Hölder's inequality, we have

$$\|\nabla f\|_{L^{p'}(Q)} \leq \|w^{-1}\|_{L^d(Q)} \|w\nabla f\|_{L^2(Q)},$$

while

$$\|\bar{f}_Q\|_{L^2(Q)} \leq |Q|^{-1/2} \|f\|_{L^1(Q)},$$

so that

$$\|f\|_{L^2(Q)} \lesssim \|w^{-1}\|_{L^d(Q)} \|w\nabla f\|_{L^2(Q)} + |Q|^{-1/2} \|f\|_{L^1(Q)},$$

and thus

$$(2.14) \quad \|f\|_{L^2(Q)}^2 \lesssim \|w^{-1}\|_{L^d(Q)}^2 \|w\nabla f\|_{L^2(Q)}^2 + |Q|^{-1} \|f\|_{L^1(Q)}^2.$$

Let $R \in \mathbb{N}$, $R \geq r$. By Jensen's inequality,

$$|Q|^{-1/d} \|w^{-1}\|_{L^d(Q)} \leq |Q|^{-1/q} \|w^{-1}\|_{L^q(Q)}.$$

For $Q \subseteq B_R$, $\|w^{-1}\|_{L^q(Q)} \leq \|w^{-1}\|_{L^q_R} \leq |B_R|^{1/q} \mathcal{M}_q$, and we have

$$(2.15) \quad \|w^{-1}\|_{L^d(Q)} \leq |Q|^{1/d-1/q} |B_R|^{1/q} \mathcal{M}_q.$$

If (Q'_i) is a finite collection of pairwise disjoint boxes, then

$$\sum_i \|f\|_{L^1(Q'_i)}^2 \leq \|f\|_{L^1(\cup Q'_i)}^2 \sum_i \|f\|_{L^1(Q'_i)} = \|f\|_{L^1(\cup Q'_i)}^2.$$

We cover the box B_R by sub-boxes (Q_i) of size r . We can do so in such a way that no point of B_R is covered by more than 2^d sub-boxes, so that

$$\sum_i \|f\|_{L^1(Q_i)}^2 \lesssim \|f\|_{L^1}^2.$$

Combining this and (2.15) into (2.14) thus yields

$$\|f\|_{L^2_R}^2 \leq \sum_i \|f\|_{L^2(Q_i)}^2 \lesssim |Q|^{2/d-2/q} |B_R|^{2/q} \mathcal{M}_q^2 \|w\nabla f\|_{L^2_R}^2 + |Q|^{-1} \|f\|_{L^1}^2,$$

where we now simply use $|Q|$ to denote the cardinality of a box of size r , without reference to a specific box. Using (2.11) (and $|Q| \geq r$), we obtain

$$\|f\|_{L^2}^2 \lesssim |Q|^{2/d-2/q} |B_R|^{2/q} \mathcal{M}_q^2 \|w\nabla f\|_{L^2}^2 + r^{-d} \|f\|_{L^1}^2 + R^{-p} \| |x|_*^{p/2} f \|_{L^2}^2.$$

By changing the implicit multiplicative constant in this inequality, we can ensure that

$$(2.16) \quad \|f\|_{L^2} \lesssim r^{1-\frac{d}{q}} R^{\frac{d}{q}} \mathcal{M}_q \|w\nabla f\|_{L^2} + r^{-d/2} \|f\|_{L^1} + R^{-p/2} \| |x|_*^{p/2} f \|_{L^2}$$

holds uniformly over $R \geq r \geq 1$. In view of (2.13), we can extend this to all $R \geq 1$, $r \geq 1$. Since we clearly have $\|f\|_{L^2} \leq \|f\|_{L^1}$ and $\|f\|_{L^2} \leq \| |x|_*^{p/2} f \|_{L^2}$, the inequality in fact holds for every $R > 0$ and $r > 0$. The conclusion will now follow by optimizing over R and r . We summarize this optimization into the following lemma, whose proof is postponed to the end of this section.

Lemma 2.10. *Let $a, a', b, c > 0$, and define $\sigma = ab + a'c + bc$. For every $A, B, D \geq 0$,*

$$\inf_{r, R > 0} \left(R^a r^{a'} A + r^{-b} B + R^{-c} D \right) \leq 3A^{bc/\sigma} B^{a'c/\sigma} D^{ab/\sigma}.$$

In our context, we let

$$a = \frac{d}{q}, \quad a' = 1 - \frac{d}{q}, \quad b = \frac{d}{2}, \quad c = \frac{p}{2},$$

then $\sigma = ab + a'c + bc$, and finally

$$\alpha = \frac{bc}{\sigma}, \quad \beta = \frac{a'c}{\sigma}, \quad \gamma = \frac{ab}{\sigma}.$$

Applying Lemma 2.10 to (2.16) leads to (2.3). In order to check that α, β and γ can be rewritten as in (2.2) for some θ , we observe that

$$(d+2)a'c + (p+2)ab = 2a'c + 2ab + d \left(1 - \frac{d}{q} \right) \frac{p}{2} + p \frac{d}{q} \frac{d}{2} = 2\sigma,$$

so that

$$(2.17) \quad \frac{d+2}{2} \beta + \frac{p+2}{2} \gamma = 1.$$

This, together with the fact that $\alpha + \beta + \gamma = 1$, ensures a representation of the form (2.2) for some θ . The value of θ can be recovered from

$$\begin{aligned} \frac{1}{\theta} &= \frac{2}{p+2} \frac{1}{\gamma} = \frac{2}{p+2} \frac{\sigma}{ab} = \frac{2}{p+2} \left[1 + \frac{2q}{d^2} \left(\left(1 - \frac{d}{q} \right) \frac{p}{2} + \frac{dp}{4} \right) \right] \\ &= \frac{2}{p+2} \left[1 + \frac{p}{2} + \left(\frac{p}{d} + \frac{p}{2} \right) \left(\frac{q}{d} - 1 \right) \right], \end{aligned}$$

and we recognize that $\theta = \theta_c$ as defined by (2.1). \square

Proof of Lemma 2.10. It suffices to consider the case when none of A, B and D are zero. We simply choose r and R so that the contributions of the three terms in the sum to be minimized are equal, that is,

$$(2.18) \quad R^a r^{a'} A = r^{-b} B = R^{-c} D.$$

The second equality gives

$$R = \left(\frac{D}{B} \right)^{1/c} r^{b/c},$$

while by the first equality,

$$\left(\frac{D}{B} \right)^{a/c} r^{ab/c + a' + b} = \frac{B}{A}.$$

This leads to

$$r^\sigma = A^{-c} B^{c+a} D^{-a},$$

and thus

$$r^{-b} B = A^{bc/\sigma} B^{1 - \frac{ab+bc}{\sigma}} D^{ab/\sigma},$$

and the lemma follows by (2.18). \square

3. STATIC ENVIRONMENTS

We now show how to use the anchored Nash inequality derived in the previous section to deduce heat kernel bounds for static, degenerate environments. We begin by defining the notion of w -moderate static environment.

Definition 3.1. Let $w : \mathbb{B} \rightarrow \mathbb{R}_+$. For a static environment $a : \mathbb{B} \rightarrow [0, 1]$, we define $p_t(x, \cdot)$ to be the unique bounded function such that $p_0(x, y) = \mathbf{1}_{x=y}$ and

$$\partial_t p_t(x, y) = \sum_{z \sim y} a(\{y, z\}) (p_t(x, z) - p_t(x, y)).$$

In probabilistic terms, $p_t(x, y)$ is the probability for the random walk in the environment a started at x to be at y at time t . We write $u_t = p_t(0, \cdot)$ and

$$\mathcal{D}_t = \|\sqrt{a} \nabla u_t\|_2^2.$$

We say that the static environment a is w -moderate if for every $t \geq 0$,

$$(3.1) \quad \|w \nabla u_t\|_2^2 \leq \mathcal{D}_t.$$

Theorem 3.2 (Upper bound for static environments). *For every $q > d$ and every $\alpha, \beta, \gamma \in (0, 1)$ as in Theorem 2.1, there exists $C < \infty$ such that if the static environment a is w -moderate, then for every $t \geq 1$,*

$$p_t(0, 0) \leq C \frac{\mathcal{M}_q^{2\alpha/\beta}}{t^{d/2}},$$

where we recall that $\mathcal{M}_q = \mathcal{M}_q(w)$ was defined in (2.4).

Before proving the theorem, we introduce some notation, and then prove an important preliminary result. For $f : \mathbb{Z}^d \rightarrow \mathbb{R}$ and $x \in \mathbb{Z}^d$, we define $\nabla_i f(x) = f(x + \mathbf{e}_i) - f(x)$. We observe that the following discrete Leibniz rule holds, for $f, g : \mathbb{Z}^d \rightarrow \mathbb{R}$:

$$\nabla_i(fg) = (\nabla_i f)g + f(\cdot + \mathbf{e}_i)\nabla_i g.$$

We write $a_i(x) = a(x, x + \mathbf{e}_i)$.

Proposition 3.3. *There exists $C < \infty$ such that*

$$(3.2) \quad \frac{d}{dt} \| |x|_*^{p/2} u_t \|_2^2 \leq C \| |x|_*^{p/2} u_t \|_2^{2(p-2)/p} \| u_t \|_2^{4/p}.$$

Proof. The proof is similar to that of [25, (81)], with some extra care required by the fact that we do not assume uniform ellipticity here. We write $\rho(x) = |x|_*$, and observe that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \| |x|_*^{p/2} u_t \|_2^2 &= \frac{1}{2} \frac{d}{dt} \sum_{x \in \mathbb{Z}^d} \rho^p u_t^2(x) \\ &= - \sum_{e \in \mathbb{B}} \nabla(\rho^p u_t) a \nabla u_t(e) \\ &= - \sum_{\substack{x \in \mathbb{Z}^d \\ 1 \leq i \leq d}} \nabla_i(\rho^p u_t) a_i \nabla_i u_t(x). \end{aligned}$$

By the discrete Leibniz rule, we have

$$\nabla_i(\rho^p u_t) = \nabla_i(\rho^p) u_t + \rho^p(\cdot + \mathbf{e}_i) \nabla_i u_t,$$

and moreover, $\nabla_i(\rho^p)(x) \lesssim \rho^{p-1}(x)$. As a consequence, the left-hand side of (3.2) is bounded by

$$\sum_{\substack{x \in \mathbb{Z}^d \\ 1 \leq i \leq d}} [-\rho^p(\cdot + \mathbf{e}_i) a_i |\nabla_i u_t|^2 + C \rho^{p-1} u_t a_i |\nabla_i u_t|] (x)$$

for some constant C . Since $a \leq 1$, we have $a \geq a^2$, so that by the Cauchy-Schwarz inequality,

$$\sum_{\substack{x \in \mathbb{Z}^d \\ 1 \leq i \leq d}} \rho^{p-1} u_t a_i |\nabla_i u_t|(x) \leq \left(\sum_{\substack{x \in \mathbb{Z}^d \\ 1 \leq i \leq d}} \rho^p a_i |\nabla_i u_t|^2(x) \right)^{1/2} \left(\sum_{\substack{x \in \mathbb{Z}^d \\ 1 \leq i \leq d}} \rho^{p-2} |u_t|^2(x) \right)^{1/2}.$$

Up to a multiplicative constant, we may replace ρ^p by $\rho^p(\cdot + \mathbf{e}_i)$ in the middle sum above. By Young's inequality, the resulting expression is bounded by

$$\frac{1}{2C} \left(\sum_{\substack{x \in \mathbb{Z}^d \\ 1 \leq i \leq d}} \rho^p(\cdot + \mathbf{e}_i) a_i |\nabla_i u_t|^2(x) \right) + \frac{C}{2} \left(\sum_{\substack{x \in \mathbb{Z}^d \\ 1 \leq i \leq d}} \rho^{p-2} |u_t|^2(x) \right),$$

for arbitrary $C > 0$. By choosing C sufficiently large, we obtain that the left-hand side of (3.2) is bounded by a constant times

$$\sum_{x \in \mathbb{Z}^d} \rho^{p-2} |u_t|^2(x) \stackrel{(\text{H\"older})}{\leq} \left(\sum_{x \in \mathbb{Z}^d} \rho^p |u_t|^2(x) \right)^{\frac{p-2}{p}} \left(\sum_{x \in \mathbb{Z}^d} |u_t|^2(x) \right)^{\frac{2}{p}},$$

and this completes the proof. \square

Proof of Theorem 3.2. Let $\mathcal{E}_t = \|u_t\|_2^2$ and $\mathcal{N}_t = \| |x|^{p/2} u_t \|_2^2$. We write \mathcal{E}'_t for the time derivative of \mathcal{E} at time t , and similarly for other quantities. Note that

$$(3.3) \quad \mathcal{E}'_t = -2\mathcal{D}_t,$$

while by Proposition 3.3,

$$(3.4) \quad \mathcal{N}'_t \lesssim \mathcal{N}_t^{(p-2)/p} \mathcal{E}_t^{2/p}.$$

Let us define

$$\Lambda_t = 1 \vee \sup_{s \leq t} s^{d/2} \mathcal{E}_s.$$

Integrating (3.4) (and since $\mathcal{N}_0 = 1$), we obtain, for every $t \geq 1/2$,

$$(3.5) \quad \mathcal{N}_t \lesssim \Lambda_t t^{(p-d)/2}.$$

For $\alpha, \beta, \gamma \in (0, 1)$ as in Theorem 2.1, we have

$$\mathcal{E}_t \lesssim (\mathcal{M}_q^2 \|w \nabla u_t\|_2^2)^\alpha \mathcal{N}_t^\gamma,$$

where we used the fact that $\|u_t\|_1 = 1$. In view of (3.1) and (3.3), it follows that

$$\mathcal{E}_t \lesssim (-\mathcal{M}_q^2 \mathcal{E}'_t)^\alpha \mathcal{N}_t^\gamma,$$

and by (3.5), for every $t \geq 1/2$,

$$(-\mathcal{M}_q^2 \mathcal{E}'_t)^\alpha \gtrsim \mathcal{E}_t \Lambda_t^{-\gamma} t^{-\gamma(p-d)/2}.$$

Integrating this relation (and recalling that Λ_t is increasing), we get

$$(3.6) \quad \mathcal{E}_t^{1-\frac{1}{\alpha}} \gtrsim \mathcal{M}_q^{-2} \Lambda_t^{-\frac{\gamma}{\alpha}} t^{1-\frac{\gamma(p-d)}{2\alpha}},$$

provided that $1 - \frac{\gamma(p-d)}{2\alpha} > 0$. In order to check this and to simplify this expression, we recall from (2.17) that

$$\frac{d}{2} \beta + \frac{p}{2} \gamma = 1 - \beta - \gamma = \alpha,$$

so that

$$(3.7) \quad \alpha - \frac{\gamma(p-d)}{2} = \frac{d}{2} \beta + \frac{d}{2} \gamma = \frac{d}{2} (1 - \alpha).$$

Since $\alpha < 1$, the inequality (3.6) is justified, and can be rewritten as

$$\mathcal{E}_t \lesssim \mathcal{M}_q^{\frac{2\alpha}{1-\alpha}} \Lambda_t^{\frac{\gamma}{1-\alpha}} t^{-\frac{d}{2}}.$$

Since Λ_t is increasing, we get that for every $s \leq t$,

$$s^{d/2} \mathcal{E}_s \lesssim \mathcal{M}_q^{\frac{2\alpha}{1-\alpha}} \Lambda_t^{\frac{\gamma}{1-\alpha}},$$

that is,

$$\Lambda_t \lesssim \mathcal{M}_q^{\frac{2\alpha}{1-\alpha}} \Lambda_t^{\frac{\gamma}{1-\alpha}}.$$

Since $\frac{\gamma}{1-\alpha} = \frac{\gamma}{\beta+\gamma} < 1$, this shows that

$$\Lambda_t \lesssim \mathcal{M}_q^{2\alpha/\beta},$$

and in particular,

$$(3.8) \quad \mathcal{E}_t \lesssim \frac{\mathcal{M}_q^{2\alpha/\beta}}{t^{d/2}}.$$

By the semi-group property,

$$p_t(0, 0) = \sum_{x \in \mathbb{Z}^d} p_{t/2}(0, x) p_{t/2}(x, 0).$$

The symmetry of $p_t(\cdot, \cdot)$ thus leads to $p_t(0, 0) \leq \mathcal{E}_{t/2}$, and in view of (3.8), this concludes the proof. \square

As a by-product of the above proof, we also obtain the following off-diagonal information.

Proposition 3.4 (Off-diagonal decay). *For every $p > d$, $q > d$ and $\alpha, \beta, \gamma \in (0, 1)$ as in Theorem 2.1, there exists $C < \infty$ such that if the static environment a is w -moderate, then for every $t \geq 1$,*

$$\| |x|^{p/2} u_t \|_2^2 \leq C \mathcal{M}_q^{2\alpha/\beta} t^{(p-d)/2}.$$

Proof. This follows directly from (3.5) and (3.8). \square

Much more precise information on the off-diagonal behaviour of the heat kernel can be obtained as an application of [24, Theorem 1.1] (see also [26, 21, 20]). We state here a particular case of this result which is most relevant to our context.

Theorem 3.5 (Improved off-diagonal decay [24]). *There exists a constant $C < \infty$ such that the following holds. Let $x, y \in \mathbb{Z}^d$ and $c(x), c(y) < \infty$. If*

$$p_t(x, x) \leq \frac{c(x)}{t^{d/2}} \quad \text{and} \quad p_t(y, y) \leq \frac{c(y)}{t^{d/2}},$$

then for all $t \geq 1 \vee |x - y|$,

$$p_t(x, y) \leq C \frac{\sqrt{c(x)c(y)}}{t^{d/2}} \exp\left(-\frac{|y-x|^2}{Ct}\right).$$

We conclude this section by giving some examples of applications of Theorem 3.2. To begin with, we show a heat kernel upper bound assuming stationarity and a moment condition. We refer to [4, Theorem 5.4] for an example (due to Noam Berger) showing that the condition is essentially sharp for a local CLT.

Proposition 3.6 (Moment condition). *Let $q > d$ and α, β, γ be as in Theorem 2.1. Assume that under the probability measure \mathbb{P} , the family of random variables $a = (a(e))_{e \in \mathbb{B}}$ is stationary with respect to translations, takes values in $(0, 1]$, and satisfies*

$$\mathbb{E} \left[a(e)^{-q/2} \right] < \infty.$$

There exists a random variable $\mathcal{X} \geq 0$ such that

$$\forall r < \frac{q\beta}{2\alpha}, \mathbb{E}[\mathcal{X}^r] < \infty$$

and

$$(3.9) \quad p_t(0,0) \leq \frac{\mathcal{X}}{t^{d/2}}.$$

Proof. Clearly, the environment a is \sqrt{a} -moderate. Hence, it suffices to check that for every $r < 1$, we have $\mathbb{E}[\mathcal{M}_q^{qr}(\sqrt{a})] < \infty$. This follows from Proposition A.1. \square

If we also assume that the conductances are independent, then we can weaken the integrability condition to $\mathbb{E}[a(e)^{-\gamma}] < \infty$ for some $\gamma > 1/4$, and thus recover some of the results of [4, Theorem 1.13], [17, Theorem 1.8].

Proposition 3.7 (Moment condition under independence). *Let $q > d$ and α, β, γ be as in Theorem 2.1. Assume that under the probability measure \mathbb{P} , the family of random variables $a = (a(e))_{e \in \mathbb{B}}$ is i.i.d., takes values in $(0, 1]$, and satisfies*

$$\mathbb{E} \left[a(e)^{-\frac{q}{4d}} \right] < \infty.$$

There exists a random variable $\mathcal{X} \geq 0$ such that

$$\forall r < \frac{q\beta}{2\alpha}, \mathbb{E}[\mathcal{X}^r] < \infty$$

and

$$(3.10) \quad p_t(0,0) \leq \frac{\mathcal{X}}{t^{d/2}}.$$

Proof. For an edge e , we denote by \underline{e}, \bar{e} the endpoints of e , conventionally chosen so that $\bar{e} - \underline{e}$ is a canonical basis vector. One can find $2d$ paths $\ell_1(e), \dots, \ell_{2d}(e)$ linking \underline{e} to \bar{e} such that no two paths have an edge in common; moreover, we can make sure that each path contains at most 9 edges (see [34, Lemma 2.1]). We define w by

$$(3.11) \quad \frac{1}{w(e)^2} = \inf_{1 \leq i \leq 2d} \sum_{b \in \ell_i(e)} \frac{1}{a(b)}.$$

We denote by $\ell(e)$ a path (among $\ell_1(e), \dots, \ell_{2d}(e)$) that realises the infimum above. By the Cauchy-Schwarz inequality,

$$\begin{aligned} |\nabla u(e)| &\leq \sum_{b \in \ell(e)} |\nabla u(b)| \\ &\leq \left(\sum_{b \in \ell(e)} \frac{1}{a(b)} \right)^{1/2} \left(\sum_{b \in \ell(e)} a(b) |\nabla u(b)|^2 \right)^{1/2} \\ &\leq w(e)^{-1} \left(\sum_{b \in \ell(e)} a(b) |\nabla u(b)|^2 \right)^{1/2}, \end{aligned}$$

so that

$$\sum_e w(e)^2 |\nabla u(e)|^2 \leq \sum_{e, b: b \in \ell(e)} a(b) |\nabla u(b)|^2.$$

Clearly, one can find $C < \infty$ such that for each b , the number of edges e such that $b \in \ell(e)$ is bounded by C . Hence, the environment a is w/C -moderate. We now

choose $q' \in (d, q)$ and check that $w(e)^{-1}$ has finite q' -th moment. By independence, for $x > 0$,

$$\mathbb{P}[w(e)^{-q'} \geq x] = \prod_{i=1}^{2d} \mathbb{P} \left[\sum_{b \in \ell_i(e)} \frac{1}{a(b)} \geq x^{2/q'} \right].$$

Since each path has length at most 9, and by stationarity,

$$\begin{aligned} \mathbb{P} \left[\sum_{b \in \ell_i(e)} \frac{1}{a(b)} \geq x^{2/q'} \right] &\leq 9 \mathbb{P} \left[\frac{1}{a(b)} \geq \frac{x^{2/q'}}{9} \right] \\ &\leq 9^{1+\frac{q}{4d}} \mathbb{E}[a(b)^{-\frac{q}{4d}}] x^{-\frac{q}{2dq'}}. \end{aligned}$$

We obtain that $\mathbb{E}[w(e)^{-q'}] < \infty$ by combining the previous two displays and the fact that $\mathbb{E}[w(e)^{-q'}] = \int_0^\infty \mathbb{P}[w(e)^{-q'} \geq x] dx$. For every $r < 1$, we thus have $\mathbb{E}[\mathcal{M}_{q'}^{q'r}(w)] < \infty$ by Proposition A.1, so the result follows by Theorem 3.2. \square

Remark 3.8. The case of percolation corresponds to assuming that $(a(e))_{e \in \mathbb{B}}$ are i.i.d. Bernoulli random variables. Obtaining heat-kernel upper bounds for super-critical percolation would require more work, which we will not pursue here (see [33, 7] for previous work on this). In [29], the following simpler situation is considered: all edges e pointing in a given direction satisfy $a(e) = 1$, while $(a(e))$ are i.i.d. Bernoulli for the other edges. In this case, one can show that the environment is w -moderate for some w such that all moments of $w(e)^{-1}$ are finite (see [29, Lemma 4]). We thus obtain (3.9) for a random variable \mathcal{X} with finite moments of every order.

4. DYNAMIC ENVIRONMENTS

We now turn our attention to *dynamic* degenerate environments. We first define the notion of (w, K) -moderate dynamic environment.

Definition 4.1. For every $e \in \mathbb{B}$, we give ourselves a piecewise continuous function

$$\begin{cases} \mathbb{R} & \rightarrow [0, 1] \\ t & \mapsto a_t(e), \end{cases}$$

and we call $a = (a_t(e))_{e \in \mathbb{B}, t \in \mathbb{R}}$ a *dynamic environment*. (It is convenient, although inessential, to consider environments defined over the whole time line.) For every $x \in \mathbb{Z}^d$ and $s \in \mathbb{R}$, we define $(p_{s,t}(x, y))_{y \in \mathbb{Z}^d, t \geq s}$ as the unique bounded function such that $p_{s,s}(x, y) = \mathbf{1}_{x=y}$ and, for every $t \geq s$,

$$\partial_t p_{s,t}(x, y) = \sum_{z \sim y} a_t(\{y, z\}) (p_{s,t}(x, z) - p_{s,t}(x, y)).$$

In probabilistic terms, $p_{s,t}(x, y)$ is the probability that the random walk evolving in the dynamic environment a and started at time s and position x reaches y at time t . We write $u_t = p_{0,t}(0, \cdot)$ and

$$\mathcal{E}_t = \|u_t\|_2^2, \quad \mathcal{D}_t = \|\sqrt{a_t} \nabla u_t\|_2^2, \quad \overline{\mathcal{D}}_t = \int_t^{+\infty} K_{s-t} \mathcal{D}_s ds.$$

Let $K : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ and, for each $e \in \mathbb{B}$, let $t \in \mathbb{R}_+ \mapsto w_t(e) \in [0, 1]$ be a measurable function. We say that the dynamic environment a is (w, K) -moderate if for every $t \geq 0$,

$$(4.1) \quad \|w_t \nabla u_t\|_2^2 \leq \overline{\mathcal{D}}_t.$$

Theorem 4.2 (Energy upper bound for dynamic environments). *Let $q > d$ and $\alpha, \beta, \gamma \in (0, 1)$ be as in Theorem 2.1. There exists $C < \infty$ such that if the dynamic environment a is (w, K) -moderate, then*

$$(4.2) \quad \mathcal{E}_t \leq C \mathcal{C}(K) \frac{\mathcal{M}_q^{2\alpha/\beta}(w)}{t^{d/2}},$$

where

$$(4.3) \quad \mathcal{M}_q^{-2}(w) = \inf_{t \geq 1} \frac{1}{t} \int_0^t \mathcal{M}_q^{-2}(w_s) \, ds$$

and

$$(4.4) \quad \mathcal{C}(K) = 1 \vee \frac{\|K\|_1^{\alpha/\beta}}{\|K\|_{L^1([0,1])}^{(1-\alpha)/\beta}}.$$

Proof. We define

$$\bar{\mathcal{E}}_t = \int_t^{+\infty} K_{s-t} \mathcal{E}_s \, ds, \quad \mathcal{N}_t = \| |x|_*^{p/2} u_t \|_2^2.$$

Proposition 3.3 remains valid for dynamic random environments, and gives us

$$(4.5) \quad \mathcal{N}'_t \lesssim \mathcal{N}'_t^{(p-2)/p} \mathcal{E}_t^{2/p}.$$

Defining $\Lambda_t = 1 \vee \sup_{s \leq t} s^{d/2} \mathcal{E}_s$, we obtain as for the static case that for every $t \geq 1$,

$$(4.6) \quad \mathcal{N}_t \lesssim \Lambda_t t^{(p-d)/2}.$$

By Theorem 2.1,

$$\mathcal{E}_t \lesssim (\mathcal{M}_q^2(w_t) \|w_t \nabla u_t\|_2^2)^\alpha \mathcal{N}_t^\gamma,$$

which by assumption (4.1) and (4.6) leads to

$$\mathcal{E}_t \lesssim (\mathcal{M}_q^2(w_t) \bar{\mathcal{D}}_t)^\alpha \Lambda_t^\gamma t^{\gamma(p-d)/2}.$$

Since \mathcal{E}_t is decreasing, we have $\bar{\mathcal{E}}_t \lesssim \|K\|_1 \mathcal{E}_t$, and moreover, $\bar{\mathcal{E}}'_t = -2\bar{\mathcal{D}}_t$. As a consequence,

$$\bar{\mathcal{E}}_t \lesssim \|K\|_1 \left(-\mathcal{M}_q^2(w_t) \bar{\mathcal{E}}'_t \right)^\alpha \Lambda_t^\gamma t^{\gamma(p-d)/2},$$

which can be rewritten as

$$(4.7) \quad -\bar{\mathcal{E}}'_t \bar{\mathcal{E}}_t^{-\frac{1}{\alpha}} \gtrsim \|K\|_1^{-\frac{1}{\alpha}} \mathcal{M}_q^{-2}(w_t) \Lambda_t^{-\frac{\gamma}{\alpha}} t^{-\frac{\gamma(p-d)}{2\alpha}}.$$

Since Γ_t is increasing and $p > d$, we obtain

$$\bar{\mathcal{E}}_t^{1-\frac{1}{\alpha}} \gtrsim \|K\|_1^{-\frac{1}{\alpha}} \Lambda_t^{-\frac{\gamma}{\alpha}} t^{-\frac{\gamma(p-d)}{2\alpha}} \int_0^t \mathcal{M}_q^{-2}(w_s) \, ds.$$

Recalling (3.7) and the definition of $\mathcal{M}_q = \mathcal{M}_q(w)$ in (4.3), we obtain that for every $t \geq 1$,

$$t^{d/2} \bar{\mathcal{E}}_t \lesssim \|K\|_1^{\frac{\alpha}{1-\alpha}} \mathcal{M}_q^{\frac{2\alpha}{1-\alpha}} \Lambda_t^{\frac{\gamma}{1-\alpha}}.$$

We now observe that, since \mathcal{E}_t is decreasing,

$$\bar{\mathcal{E}}_t \geq \mathcal{E}_{2t} \int_0^t K_s \, ds,$$

and in particular, for $t \geq 1$, $\bar{\mathcal{E}}_t \geq \|K\|_{L^1([0,1])} \mathcal{E}_{2t}$. Since moreover, $\Lambda_2 \leq 2^{d/2}$, we get

$$\Lambda_t \lesssim 1 + \frac{\|K\|_1^{\frac{\alpha}{1-\alpha}}}{\|K\|_{L^1([0,1])}} \mathcal{M}_q^{\frac{2\alpha}{1-\alpha}} \Lambda_t^{\frac{\gamma}{1-\alpha}},$$

and this proves the theorem since $\mathcal{C}(K) \geq 1$ and $\mathcal{M}_q \geq 1$ (recall that $w \leq 1$). \square

As before, the proof above reveals more detailed information about off-diagonal decay of the heat kernel.

Proposition 4.3 (Off-diagonal decay). *Let $p > d$, $q > d$ and $\alpha, \beta, \gamma \in (0, 1)$ be as in Theorem 2.1. There exists $C < \infty$ such that if the dynamic environment a is (w, K) -moderate, then for every $t \geq 1$,*

$$\| |x|_*^{p/2} u_t \|_2^2 \leq C \mathcal{C}(K) \mathcal{M}_q^{2\alpha/\beta}(w) t^{(p-d)/2}.$$

Proof. This follows from (4.6) and (4.2). \square

For a dynamic environment a , we let $a^{(t)}$ be the dynamic environment defined, for all $s \in \mathbb{R}$, by $a_s^{(t)} = a_{t-s}$.

Corollary 4.4 (Heat kernel upper bound for dynamic environments). *In the setting of Theorem 4.2, there exists $C < \infty$ such that if a is (w, K) -moderate and $a^{(t)}$ is $(w^{(t)}, K^{(t)})$ -moderate, then*

$$(4.8) \quad p_{0,t}(0, 0) \leq C \sqrt{\mathcal{C}(K)\mathcal{C}(K^{(t)})} \frac{(\mathcal{M}_q(w) \cdot \mathcal{M}_q(w^{(t)}))^{\alpha/\beta}}{t^{d/2}}$$

Corollary 4.4 will be obtained from the following lemma.

Lemma 4.5 (Space-time reversal). *For every $u \leq v \in \mathbb{R}$, $t \in \mathbb{R}$ and $x, y \in \mathbb{Z}^d$,*

$$p_{u,v}(x, y) = p_{t-v, t-u}^{(t)}(y, x).$$

Proof. For $f : \mathbb{Z}^d \rightarrow \mathbb{R}$ of compact support, we write

$$f_{u,v}(x) = \sum_{y \in \mathbb{Z}^d} p_{u,v}(x, y) f(y),$$

$$f_{u,v}^{(t)}(x) = \sum_{y \in \mathbb{Z}^d} p_{u,v}^{(t)}(x, y) f(y).$$

The lemma is equivalent to the claim that for every f, g of compact support,

$$(f_{u,v}, g) = (f, g_{t-v, t-u}^{(t)}),$$

where (\cdot, \cdot) denotes the scalar product on $L^2(\mathbb{Z}^d)$. That this relation is correct is clear if $t \mapsto a_t$ is piecewise constant. Otherwise, one can proceed via finite-volume approximations. \square

Proof of Corollary 4.4. From the lemma,

$$p_{t/2, t}(x, 0) = p_{0, t/2}^{(t)}(0, x).$$

Hence, a Cauchy-Schwarz inequality gives

$$p_{0,t}(0, 0) = \sum_{x \in \mathbb{Z}^d} p_{0, t/2}(0, x) p_{t/2, t}(x, 0) \leq \sqrt{\mathcal{E}_{t/2} \mathcal{E}_{t/2}^{(t)}}$$

(with obvious notation), and the result thus follows from Theorem 4.2. \square

We conclude this section by providing a method to build (w, K) out of a dynamic environment a , so that a is (w, K) -moderate.

Proposition 4.6 (A criterion for moderation). *There exists $c < \infty$ such that the following holds. Let $k : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ be such that $\int (1 + t^2) k_t dt < \infty$, let*

$$K_t = k_t + \int_t^{+\infty} s k_s ds,$$

and for a dynamic environment a , let

$$(4.9) \quad w_t^2 = \int_t^{+\infty} k_{s-t} a_s ds.$$

Then the dynamic environment a is (w, cK) -moderate.

Remark 4.7. Recall that in order to get non-trivial bounds from Theorem 4.2, we still need to make sure that w^{-q} is integrable for some $q > d$ (the fact that $K \in L^1$ is automatic from the assumption on k). This can be obtained if one knows that it is unlikely for a_s to remain close to 0 during a long time interval.

Proof. The condition $\int(1+t^2)k_t dt < \infty$ ensures that K is integrable. We need to estimate

$$\|w_t \nabla u_t\|_2^2 = \sum_{e \in \mathbb{B}} w_t^2 (\nabla u_t)^2(e) = \int_t^{+\infty} k_{s-t} a_s (\nabla u_t)^2(e) ds.$$

This quantity is bounded by

$$\int_t^{+\infty} k_{s-t} \underbrace{\sum_{e \in \mathbb{B}} a_s (\nabla u_s)^2(e)}_{=\mathcal{D}_s} ds + \int_t^{+\infty} k_{s-t} \sum_{e \in \mathbb{B}} a_s (\nabla u_s - \nabla u_t)^2(e) ds.$$

Up to a constant (and since $a_s \leq 1$), the second sum above is bounded by

$$\begin{aligned} \sum_{x \in \mathbb{Z}^d} (u_s - u_t)^2(x) &= \sum_{e \in \mathbb{B}} \left(\int_t^s \nabla^* a_{s'} \nabla u_{s'}(e) ds' \right)^2 \\ &\stackrel{(\text{Jensen})}{\leq} (s-t) \sum_{e \in \mathbb{B}} \int_t^s (\nabla^* a_{s'} \nabla u_{s'}(e))^2 ds', \end{aligned}$$

which, up to a constant, is bounded by

$$(s-t) \int_t^s \sum_{e \in \mathbb{B}} a_{s'} (\nabla u_{s'}(e))^2 ds' = (s-t) \int_t^s \mathcal{D}_{s'} ds'.$$

We have thus shown that

$$\begin{aligned} \|w_t \nabla u_t\|_2^2 &\lesssim \int_t^{+\infty} k_{s-t} \mathcal{D}_s ds + \int_t^{+\infty} (s-t) k_{s-t} \int_t^s \mathcal{D}_{s'} ds' ds \\ &\lesssim \int_t^{+\infty} k_{s-t} \mathcal{D}_s ds + \int_t^{+\infty} ds' \mathcal{D}_{s'} \int_{s'}^{+\infty} (s-t) k_{s-t} ds, \end{aligned}$$

which is the desired result. \square

5. THE EXCLUSION PROCESS AS DYNAMIC RANDOM ENVIRONMENT

As an example of the results of the previous section, we study a random walk evolving in an environment that is a local function of the symmetric simple exclusion process on \mathbb{Z}^d , $d \geq 2$. Recall that the infinitesimal generator \mathcal{L} of the symmetric simple exclusion process is given by

$$(5.1) \quad \mathcal{L}f(\eta) = \sum_{e \in \mathbb{B}} (f(\eta^e) - f(\eta)),$$

where $\eta = (\eta(x))_{x \in \mathbb{Z}^d} \in \{0, 1\}^{\mathbb{Z}^d}$, f is a function that depends on a finite number of coordinates of η , and

$$(\eta^e)(x) = \begin{cases} \eta(y) & \text{if } \{x, y\} = e, \\ \eta(x) & \text{otherwise.} \end{cases}$$

This process is also called the Kawasaki dynamics. (Strictly speaking, the exclusion process is thought of as particles performing independent random walks, but with any jump of one particle onto another being suppressed; since the particles are indistinguishable here, this is equivalent to the process defined by (5.1).) We refer to [30, Part III] for a thorough study of this process.

For every $\rho \in [0, 1]$, let $\mu_\rho = \bigotimes_{\mathbb{Z}^d} \text{Bernoulli}(\rho)$. The measures μ_ρ are the only extremal invariant (and in fact, reversible) measures of the process [30, Corollary III.1.11]. From now on, we fix $\rho \in (0, 1)$. By stationarity, we can build the probability measure \mathbb{P} such that under this measure, the process $(\eta_t)_{t \in \mathbb{R}}$ (defined over the whole time line) evolves as an exclusion process, and η_t is distributed according to μ_ρ for every $t \in \mathbb{R}$.

For concreteness, we define our dynamic environment by

$$a_t(e) = \begin{cases} 1 & \text{if } \eta_t(\underline{e}) = \eta_t(\bar{e}) = 0, \\ 0 & \text{otherwise,} \end{cases}$$

where \underline{e} and \bar{e} are the end-points of e . The specific form of a is not important, as long as it is a function of the *unoriented* edge e , and depends only locally on the configuration η . We let $(p_{s,t}(x, y))_{s \leq t \in \mathbb{R}, x, y \in \mathbb{Z}^d}$ be as in Definition 4.1.

Theorem 5.1 (Heat kernel estimates). *There exists a random variable \mathcal{X} with finite moments of every order and such that*

$$(5.2) \quad \sum_{x \in \mathbb{Z}^d} (p_{0,t}(0, x))^2 \leq \frac{\mathcal{X}}{t^{d/2}}.$$

Moreover, there exists a stationary process $(\mathcal{Y}_t)_{t \in \mathbb{R}}$ satisfying

$$(5.3) \quad p_{0,t}(0, 0) \leq \frac{\mathcal{Y}_t}{t^{d/2}},$$

and such that for every $r > 0$ and $\varepsilon > 0$,

$$(5.4) \quad \mathbb{E}[(\mathcal{Y}_t)^r] < \infty,$$

$$(5.5) \quad \mathbb{E} \left[\left(\sup_{t \geq 1} \frac{\mathcal{Y}_t}{t^\varepsilon} \right)^r \right] < \infty.$$

Remark 5.2. We believe that $\sup_t t^{d/2} p_{0,t}(0, 0)$ should have finite moments of every order, but proving this would require more work, so we do not pursue this question further here.

We begin by showing that $a_t(e)$ is unlikely to remain equal to 0 for a long period of time.

Lemma 5.3. *For every $\kappa < 1$, there exists $C < \infty$ such that*

$$\mathbb{P} \left[\int_0^t a_s(e) ds \leq 1 \right] \leq C \exp(-t^\kappa).$$

Proof. Throughout the proof, the value of the constants $c > 0$ and $C < \infty$ may vary in each occurrence. Let $\delta \in (1/d, 1)$, and denote by $A_{s,t}$ the event

$$\text{there exist } x \neq y \in B_{\log^\delta t} \text{ s.t. } \eta_s(x) = \eta_s(y) = 0.$$

Recall that under \mathbb{P} and for every time s , the configuration η_s is distributed according to the invariant measure μ_ρ . In other words, for any fixed s , the random variables

$(\eta_s(x))_{x \in \mathbb{Z}^d}$ are independent Bernoulli random variables of parameter $\rho \in (0, 1)$. As a consequence, for every $s \in \mathbb{R}$ and $t \geq 2$,

$$\begin{aligned} 1 - \mathbb{P}[A_{s,t}] &\leq \sum_{x \in B_{\log^\delta t}} \mathbb{P}[\eta_s \equiv 1 \text{ on } B_{\log^\delta t} \setminus \{x\}] \\ &= |B_{\log^\delta t}| \exp[(|B_{\log^\delta t}| - 1) \log(\rho)] \\ &\leq C \exp(-c \log^{\delta d} t). \end{aligned}$$

By a union bound (and since $\delta > 1/d$), we obtain

$$1 - \mathbb{P}\left[\forall k \text{ integer, } 0 \leq k \leq \frac{t}{\log^\delta t} - 1, A_{k \log^\delta t, t} \text{ holds}\right] \leq C \exp(-c \log^{\delta d} t).$$

Let e be an edge adjacent to the origin. Conditionally on $A_{0,t}$, we can estimate the probability that $a_s(e)$ becomes 1 by considering a particular choice of moves that bring the two holes in $B_{\log^\delta t}$ onto the two endpoints of e , and then asking that they stay put for one unit of time. More precisely, for every $t \geq 2$,

$$\mathbb{P}\left[\int_0^{\log^\delta t} a_s(e) ds \geq 1 \mid A_{0,t}\right] \geq \exp(-C \log^\delta t)$$

(the bound being very crude). By the Markov property,

$$\begin{aligned} \mathbb{P}\left[\int_0^t a_s(e) ds \leq 1\right] &\leq C \exp(-c \log^{\delta d} t) + \left[1 - \exp(-C \log^\delta t)\right]^{\lfloor \frac{t}{\log^\delta t} \rfloor} \\ &\leq C \exp(-c \log^{\delta d} t) + \exp\left(\left\lfloor \frac{t}{\log^\delta t} \right\rfloor \log\left[1 - \exp(-C \log^\delta t)\right]\right). \end{aligned}$$

The first term in the sum above is controlled since $\delta > 1/d$. For the second term, the fact that $\delta < 1$ ensures that the logarithmic term decays to 0 slower than any negative power of t , so the proof is complete. \square

Proof of Theorem 5.1. We let $k_t = (1+t)^{-4}$, and K, w be defined as in Proposition 4.6, so that the dynamic environment a is (w, cK) -moderate. For every $u \in (0, 1)$,

$$\begin{aligned} \mathbb{P}[w_t(e) \leq u] &= \mathbb{P}\left[\int_0^\infty k_s a_s(e) ds \leq \sqrt{u}\right] \\ &\leq \mathbb{P}\left[(1+t)^{-4} \int_0^t a_s(e) ds \leq \sqrt{u}\right], \end{aligned}$$

for arbitrary $t \geq 0$. We choose t such that $(1+t)^{-4} = \sqrt{u}$, and apply Lemma 5.3 to obtain

$$(5.6) \quad \mathbb{P}[w_t(e) \leq u] \leq C \exp(-u^{-1/9}).$$

In particular, $w_t(e)^{-1}$ has finite moments of every order. By Corollary A.2, the maximal function $\mathcal{M}_q(w_t)$ has finite moments of every order. The same property holds for $\mathcal{M}_q(w)$ itself by Remark A.3. (Strictly speaking, Remark A.3 applies only to processes indexed by a discrete time, but it is not difficult to check that this is sufficient for our purpose.) By Theorem 4.2, we thus obtain (5.2) (the choice of $q \in (d, \infty)$ and of α, β, γ as in Theorem 2.1 is arbitrary). Recall that we denote the time reversals around time t of a and w by $a^{(t)}$ and $w^{(t)}$ respectively. Clearly, $a^{(t)}$ is $(w^{(t)}, cK)$ -moderate, so Corollary 4.4 yields that

$$p_{0,t}(0, 0) \leq \frac{\mathcal{Y}_t}{t^{d/2}},$$

with

$$\mathcal{Y}_t = C \left(\mathcal{M}_q(w) \cdot \mathcal{M}_q(w^{(t)}) \right)^{\alpha/\beta}$$

for some constant $C < \infty$. Since $w^{(t)}$ and w have the same law, every moment of $\mathcal{M}_q(w^{(t)})$ is finite, and (5.4) is proved. In order to obtain (5.5), it suffices to consider a supremum over integer times, since for any $u \in [t-1, t]$, we have $\mathcal{M}_q(w^{(u)}) \leq 2\mathcal{M}_q(w^{(t)})$. We then note that for any $r > 1/\varepsilon$ and $y > 0$,

$$\mathbb{P} \left[\sup_{n \in \mathbb{N} \setminus \{0\}} n^{-\varepsilon} \mathcal{Y}_n \geq y \right] \leq \sum_{n=1}^{+\infty} \frac{\mathbb{E}[(\mathcal{Y}_n)^r]}{(yn^\varepsilon)^r} \leq \frac{C}{y^r},$$

so that (5.5) follows. \square

APPENDIX A. MAXIMAL INEQUALITIES

In this appendix, we recall classical properties of maximal functions in a multi-dimensional setting. Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a probability space, and $(\theta_x)_{x \in \mathbb{Z}^d}$ be a measure-preserving action of \mathbb{Z}^d on this space. For every measurable function $f : \Omega \rightarrow \mathbb{R}$, we define the maximal function

$$\mathbf{M}f(\omega) = \sup_{r \in \mathbb{N}} \frac{1}{|B_r|} \sum_{x \in B_r} f(\theta_x \omega).$$

The following result is [1, Theorem 3.2].

Proposition A.1 (Weak type (1,1) estimate). *For every $f \in L^1(\Omega)$ and $\lambda > 0$,*

$$\mathbb{P}[\mathbf{M}f \geq \lambda] \leq \frac{3^d \|f\|_{L^1(\Omega)}}{\lambda}.$$

Obviously, the maximal function defines a bounded operator from L^∞ to L^∞ . By the Marcinkiewicz interpolation theorem (see e.g. [40, Appendix D]), we thus have the following.

Corollary A.2 (L^p estimate). *For every $p \in (1, \infty]$, there exists $C_p < \infty$ such that*

$$\|\mathbf{M}f\|_{L^p(\Omega)} \leq C_p \|f\|_{L^p(\Omega)}.$$

Remark A.3. If we let

$$\mathbf{m}g(\omega) = \inf_{r \in \mathbb{N}} \frac{1}{|B_r|} \sum_{x \in B_r} g(\theta_x \omega),$$

then for every $p \in (1, \infty]$ and $g > 0$, we also have

$$\|(\mathbf{m}g)^{-1}\|_{L^p(\Omega)} \leq C_p \|g^{-1}\|_{L^p(\Omega)},$$

since by Jensen's inequality, $(\mathbf{m}g)^{-1} \leq \mathbf{M}(g^{-1})$.

REFERENCES

- [1] M. A. Akcoglu and U. Krengel. Ergodic theorems for superadditive processes. *J. Reine Angew. Math.*, 323:53–67, 1981.
- [2] S. Andres, M. T. Barlow, J.-D. Deuschel, and B. M. Hambly. Invariance principle for the random conductance model. *Probab. Theory Related Fields*, 156(3-4):535–580, 2013.
- [3] S. Andres, J.-D. Deuschel, and M. Slowik. Invariance principle for the random conductance model in a degenerate ergodic environment. Preprint, arXiv:1306.2521, 2013.
- [4] S. Andres, J.-D. Deuschel, and M. Slowik. Harnack inequalities on weighted graphs and some applications to the random conductance model. Preprint, arXiv:1312.5473, 2013.
- [5] S. Andres, J.-D. Deuschel, and M. Slowik. Heat kernel estimates for random walks with degenerate weights. Preprint, arXiv:1412.4338, 2014.
- [6] H. Bahouri, J.-Y. Chemin, and R. Danchin. *Fourier analysis and nonlinear partial differential equations*, volume 343 of *Grundlehren der Mathematischen Wissenschaften*. Springer, Heidelberg, 2011.

- [7] M. T. Barlow. Random walks on supercritical percolation clusters. *Ann. Probab.*, 32(4):3024–3084, 2004.
- [8] M. T. Barlow and B. M. Hambly. Parabolic Harnack inequality and local limit theorem for percolation clusters. *Electron. J. Probab.*, 14:no. 1, 1–27, 2009.
- [9] N. Berger and M. Biskup. Quenched invariance principle for simple random walk on percolation clusters. *Probab. Theory Related Fields*, 137(1-2):83–120, 2007.
- [10] N. Berger, M. Biskup, C. E. Hoffman, and G. Kozma. Anomalous heat-kernel decay for random walk among bounded random conductances. *Ann. Inst. Henri Poincaré Probab. Stat.*, 44(2):374–392, 2008.
- [11] L. Bertini and B. Zegarlinski. Coercive inequalities for Kawasaki dynamics. The product case. *Markov Process. Related Fields*, 5(2):125–162, 1999.
- [12] M. Biskup. Recent progress on the random conductance model. *Probab. Surv.*, 8:294–373, 2011.
- [13] M. Biskup and O. Boukhadra. Subdiffusive heat-kernel decay in four-dimensional i.i.d. random conductance models. *J. Lond. Math. Soc. (2)*, 86(2):455–481, 2012.
- [14] M. Biskup, O. Luidor, A. Rozinov, and A. Vandenberg-Rodes. Trapping in the random conductance model. *J. Stat. Phys.*, 150(1):66–87, 2013.
- [15] M. Biskup and T. M. Prescott. Functional CLT for random walk among bounded random conductances. *Electron. J. Probab.*, 12:no. 49, 1323–1348, 2007.
- [16] O. Boukhadra. Heat-kernel estimates for random walk among random conductances with heavy tail. *Stochastic Process. Appl.*, 120(2):182–194, 2010.
- [17] O. Boukhadra, T. Kumagai, and P. Mathieu. Harnack inequalities and local central limit theorem for the polynomial lower tail random conductance model. Preprint, arXiv:1308.1067, 2015.
- [18] S. Buckley. Anomalous heat kernel behaviour for the dynamic random conductance model. *Electron. Commun. Probab.*, 18:no. 1, 11, 2013.
- [19] E. A. Carlen, S. Kusuoka, and D. W. Stroock. Upper bounds for symmetric Markov transition functions. *Ann. Inst. H. Poincaré Probab. Statist.*, 23(2, suppl.):245–287, 1987.
- [20] X. Chen. Pointwise upper estimates for transition probability of continuous time random walks on graphs. Preprint, arXiv:1310.2680, 2013.
- [21] T. Coulhon, A. Grigor’yan, and F. Zucca. The discrete integral maximum principle and its applications. *Tohoku Math. J. (2)*, 57(4):559–587, 2005.
- [22] A. Fannjiang and T. Komorowski. A martingale approach to homogenization of unbounded random flows. *Ann. Probab.*, 25(4):1872–1894, 1997.
- [23] A. Fannjiang and T. Komorowski. An invariance principle for diffusion in turbulence. *Ann. Probab.*, 27(2):751–781, 1999.
- [24] M. Folz. Gaussian upper bounds for heat kernels of continuous time simple random walks. *Electron. J. Probab.*, 16:no. 62, 1693–1722, 2011.
- [25] A. Gloria, S. Neukamm, and F. Otto. Quantification of ergodicity in stochastic homogenization: optimal bounds via spectral gap on Glauber dynamics. *Invent. Math.*, 199(2):455–515, 2015.
- [26] A. Grigor’yan. Gaussian upper bounds for the heat kernel on arbitrary manifolds. *J. Differential Geom.*, 45(1):33–52, 1997.
- [27] E. Janvresse, C. Landim, J. Quastel, and H. T. Yau. Relaxation to equilibrium of conservative dynamics. I. Zero-range processes. *Ann. Probab.*, 27(1):325–360, 1999.
- [28] T. Kumagai. *Random walks on disordered media and their scaling limits*, volume 2101 of *Lecture Notes in Mathematics*. Springer, Cham, 2014. Lecture notes from the 40th Probability Summer School held in Saint-Flour, 2010, École d’Été de Probabilités de Saint-Flour.
- [29] A. Lamacz, S. Neukamm, and F. Otto. Moment bounds for the corrector in stochastic homogenization of a percolation model. Preprint, arXiv:1406.5723, 2014.
- [30] T. M. Liggett. *Stochastic interacting systems: contact, voter and exclusion processes*, volume 324 of *Grundlehren der Mathematischen Wissenschaften*. Springer-Verlag, Berlin, 1999.
- [31] P. Mathieu. Quenched invariance principles for random walks with random conductances. *J. Stat. Phys.*, 130(5):1025–1046, 2008.
- [32] P. Mathieu and A. Piatnitski. Quenched invariance principles for random walks on percolation clusters. *Proc. R. Soc. Lond. Ser. A Math. Phys. Eng. Sci.*, 463(2085):2287–2307, 2007.
- [33] P. Mathieu and E. Remy. Isoperimetry and heat kernel decay on percolation clusters. *Ann. Probab.*, 32(1A):100–128, 2004.
- [34] J.-C. Mourrat. Lyapunov exponents, shape theorems and large deviations for the random walk in random potential. *ALEA Lat. Am. J. Probab. Math. Stat.*, 9:165–211, 2012.
- [35] J. Nash. Continuity of solutions of parabolic and elliptic equations. *Amer. J. Math.*, 80:931–954, 1958.
- [36] Y. Peres, A. Stauffer, and J. Steif. Random walks on dynamical percolation: mixing times, mean squared displacement and hitting times. *Probab. Theory Related Fields*, to appear.

- [37] E. Procaccia, R. Rosenthal, and A. Sapozhnikov. Quenched invariance principle for simple random walk on clusters in correlated percolation models. Preprint, arXiv:1310.4764, 2013.
- [38] L. Saloff-Coste. Lectures on finite Markov chains. In *Lectures on probability theory and statistics (Saint-Flour, 1996)*, volume 1665 of *Lecture Notes in Math.*, pages 301–413. Springer, Berlin, 1997.
- [39] V. Sidoravicius and A.-S. Sznitman. Quenched invariance principles for walks on clusters of percolation or among random conductances. *Probab. Theory Related Fields*, 129(2):219–244, 2004.
- [40] M. E. Taylor. *Measure theory and integration*, volume 76 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, 2006.

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