

Anomalous fluctuations for a perturbed hamiltonian system with exponential interactions

Cedric Bernardin, P. Gonçalves

► **To cite this version:**

Cedric Bernardin, P. Gonçalves. Anomalous fluctuations for a perturbed hamiltonian system with exponential interactions. Communications in Mathematical Physics, Springer Verlag, 2013, pp.44. <ensl-00695381v2>

HAL Id: ensl-00695381

<https://hal-ens-lyon.archives-ouvertes.fr/ensl-00695381v2>

Submitted on 29 Nov 2013

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Anomalous fluctuations for a perturbed Hamiltonian system with exponential interactions

Cédric Bernardin¹, Patrícia Gonçalves²

¹ Université de Lyon and CNRS, UMPA, UMR-CNRS 5669, ENS-Lyon, 46, allée d'Italie, 69364 Lyon Cedex 07 - France.

E-mail: Cedric.Bernardin@ens-lyon.fr

² CMAT, Centro de Matemática da Universidade do Minho, Campus de Gualtar, 4710-057 Braga, Portugal.

E-mail: patg@math.uminho.pt

The date of receipt and acceptance will be inserted by the editor

Abstract: A one-dimensional Hamiltonian system with exponential interactions perturbed by a conservative noise is considered. It is proved that energy superdiffuses and upper and lower bounds describing this anomalous diffusion are obtained.

1. Introduction

Over the last decade, transport properties of one-dimensional Hamiltonian systems consisting of coupled oscillators on a lattice have been the subject of many theoretical and numerical studies, see the review papers [9, 11, 24]. Despite many efforts, our knowledge of the fundamental mechanisms necessary and/or sufficient to have a normal diffusion remains very limited. Nevertheless, it has been recognized that conservation of momentum plays a major role and numerical simulations provide a strong evidence of the fact that one dimensional chains of anharmonic oscillators conserving momentum are usually ¹ superdiffusive.

An interesting area of current research consists in studying this problem for hybrid models where a stochastic perturbation is superposed to the deterministic evolution. Even if the problem is considerably simplified, several open challenging questions can be addressed for these systems. In [2] it is proved that the thermal conductivity of an unpinned harmonic chain of oscillators perturbed by an energy-momentum conservative noise is infinite while if a pinning potential (destroying momentum conservation) is added it is finite. In the same paper, diverging upper bounds are provided when some nonlinearities are added. This does not, however, exclude the possibility of having a finite conductivity. Therefore much more interesting would be to obtain lower bounds showing that the

¹ See however the coupled-rotor model which displays normal behavior (see [24], Section 6.4).

conductivity is infinite and that energy superdiffuses, but this problem is left open in [2].

In [7], has been introduced and studied numerically, a class of Hamiltonian models for which anomalous diffusion is observed. There, the investigated systems present strong analogies with standard chains of oscillators. They can be described as follows. Let V and U be two non-negative potentials on \mathbb{R} and consider the Hamiltonian system $(\mathbf{r}(t), \mathbf{p}(t))_{t \geq 0}$ whose equations of motion are given by

$$\frac{dp_x}{dt} = V'(r_{x+1}) - V'(r_x), \quad \frac{dr_x}{dt} = U'(p_x) - U'(p_{x-1}), \quad x \in \mathbb{Z}, \quad (1)$$

where p_x is the momentum of the particle x , q_x its position and $r_x = q_x - q_{x-1}$ is the “deformation” of the lattice at x . Standard chains of oscillators are recovered for a quadratic kinetic energy $U(p) = p^2/2$. Now, take $V = U$, and call $\eta_{2x-1} = r_x$ and $\eta_{2x} = p_x$. The dynamics can be rewritten as:

$$d\eta_x(t) = \left(V'(\eta_{x+1}) - V'(\eta_{x-1}) \right) dt. \quad (2)$$

Notice that with these new variables the energy of the system is simply given by $\sum_{x \in \mathbb{Z}} V(\eta_x)$. In [7] an anomalous diffusion of energy is numerically observed for a generic potential V . Then, following the spirit of [2], the deterministic evolution is perturbed by adding a noise which consists to exchange η_x with η_{x+1} at random exponential times, independently for each bond $\{x, x+1\}$. The dynamics still conserves the energy $\sum_{x \in \mathbb{Z}} V(\eta_x)$ and the “volume” $\sum_{x \in \mathbb{Z}} \eta_x$ and destroys all other conserved quantities. As argued in [7], the volume conservation law is responsible for the anomalous energy diffusion observed for this class of energy-volume conserving dynamics. This can be shown for quadratic interactions ([7]) with a behavior similar to the one observed in [2]. For nonlinear interactions the problem is much more difficult.

The aim of this paper is to show that if the interacting potential is of exponential type then the energy superdiffuses. Therefore, for this class of related models, in a particular case, we answer to the open question stated in [2]. With some additional technical work we think that our methods could be carried out to the Toda lattice perturbed by an energy-momentum conserving noise (considered e.g. in [19]). The exponential form of the potential V makes the deterministic dynamics given by (2) completely integrable. Nevertheless our proofs do not rely on this exceptional property of the dynamics and could be potentially generalized to other potentials V . The main ingredient used is the existence of explicit orthogonal polynomials for the equilibrium measures (see Section 5).

The paper is organized as follows. In Section 2 we define precisely the model. The results are stated in Section 3. To prove the theorems we first perform a microscopic change of variables (Section 4) which permits to use a nice orthogonal decomposition of the generator (Section 5). Roughly speaking the upper bound on the energy superdiffusion is proved in Section 6 and the lower bound in Section 7. Section 8 contains a comment about the possible extensions and comparisons of our model to others. In the Appendix we prove the existence of the infinite dynamics.

Notations: For any $a, b \in \mathbb{R}^2$, $a \cdot b$ stands for the standard scalar product between a and b and $|a| = \sqrt{a \cdot a}$ for the norm of a . The transpose of a matrix M is

denoted by M^T . If $u : \mathbf{x} = (x_1, \dots, x_n)^T \in \mathbb{R}^n \rightarrow u(\mathbf{x}) = (u_1(\mathbf{x}), \dots, u_d(\mathbf{x}))^T \in \mathbb{R}^d$ is a differentiable function then $\partial_{x_j} u_i(\mathbf{x})$ denotes the partial derivative of u_j with respect to the j -th coordinate at \mathbf{x} and $\nabla u(\mathbf{x})$ denotes the differential matrix (the gradient if $d = 1$) of u at \mathbf{x} , i.e. the $n \times d$ matrix whose (i, j) -th entry is $\partial_{x_j} u_i(\mathbf{x})$; if $u := (u_1, \dots, u_d)^T : \mathbb{Z} \rightarrow \mathbb{R}^d$ then we adopt the same notation to denote the discrete gradient of u defined by $\nabla u := (\nabla u_1, \dots, \nabla u_d)^T$ with $\nabla u_i(x) = u_i(x+1) - u_i(x)$.

2. The model

Let $b > 0$ and $V_b(q) = e^{-bq} - 1 + bq$. We consider the system $\eta(t) = \{\eta_x(t) : x \in \mathbb{Z}\}$ on $\mathbb{R}^{\mathbb{Z}}$ defined by its generator $L = A + \gamma S$, $\gamma > 0$, where for local² differentiable functions $f : \mathbb{R}^{\mathbb{Z}} \rightarrow \mathbb{R}$ we have that

$$(Af)(\eta) = \sum_{x \in \mathbb{Z}} \left(V'_b(\eta_{x+1}) - V'_b(\eta_{x-1}) \right) (\partial_{\eta_x} f)(\eta)$$

and

$$(Sf)(\eta) = \sum_{x \in \mathbb{Z}} \left(f(\eta^{x,x+1}) - f(\eta) \right),$$

where $\eta^{x,x+1}$ is obtained from η by exchanging the variables η_x and η_{x+1} , namely

$$\eta_y^{x,x+1} = \begin{cases} \eta_{x+1}, & \text{if } y = x, \\ \eta_x, & \text{if } y = x+1, \\ \eta_y, & \text{otherwise.} \end{cases} \quad (3)$$

The deterministic system (2) with potential V_b is well known in the integrable systems literature. It has been introduced in [21] by Kac and van Moerbeke and was shown to be completely integrable. Consequently, the energy transport is ballistic ([7, 32]). As we will see this is different when the noise is added: the energy transport is no more ballistic but superdiffusive.

The existence of the dynamics generated by L is proved in the Appendix for a large set of initial conditions and in particular for a set of full measure w.r.t. any invariant state $\mu_{\bar{\beta}, \bar{\lambda}}$ (see below for its definition).

The system conserves the energy $\sum_{x \in \mathbb{Z}} V_b(\eta_x)$ and the volume $\sum_{x \in \mathbb{Z}} \eta_x$. In fact, we have

$$L(V_b(\eta_x)) = -\nabla \bar{j}_{x-1,x}(\eta), \quad L(\eta_x) = -\nabla \bar{j}'_{x-1,x}(\eta),$$

where the microscopic currents are given by

$$\bar{j}_{x,x+1}(\eta) = -b^2 e^{-b(\eta_x + \eta_{x+1})} + b^2 (e^{-b\eta_x} + e^{-b\eta_{x+1}}) - \gamma \nabla V_b(\eta_x)$$

and

$$\bar{j}'_{x,x+1}(\eta) = b e^{-b\eta_x} + b e^{-b\eta_{x+1}} - \gamma \nabla \eta_x.$$

² A function f defined on an infinite product space is said to be local if it depends only on its variable through a finite number of coordinates.

Every product probability measure $\mu_{\bar{\beta}, \bar{\lambda}}$ on $\mathbb{R}^{\mathbb{Z}}$ in the form

$$\mu_{\bar{\beta}, \bar{\lambda}}(d\eta) = \prod_{x \in \mathbb{Z}} \bar{Z}^{-1}(\bar{\beta}, \bar{\lambda}) \exp\{-\bar{\beta}e^{-b\eta_x} - \bar{\lambda}\eta_x\} d\eta_x, \quad \bar{\beta} > 0, \bar{\lambda} > 0$$

is invariant for the dynamics.

Let $\langle \cdot \rangle_{\mu_{\bar{\beta}, \bar{\lambda}}}$ denote the average with respect to $\mu_{\bar{\beta}, \bar{\lambda}}$. We define $\bar{e} := \bar{e}(\bar{\beta}, \bar{\lambda}), \bar{v} := \bar{v}(\bar{\beta}, \bar{\lambda})$ as the averages of the conserved quantities $V_b(\eta_x), \eta_x$ with respect to $\mu_{\bar{\beta}, \bar{\lambda}}$, respectively, namely $\bar{e} = \langle V_b(\eta_x) \rangle_{\mu_{\bar{\beta}, \bar{\lambda}}}$ and $\bar{v} = \langle \eta_x \rangle_{\mu_{\bar{\beta}, \bar{\lambda}}}$.

A simple computation shows that

$$\langle \bar{J}_{x, x+1} \rangle_{\mu_{\bar{\beta}, \bar{\lambda}}} = -b^2(\bar{e} - b\bar{v})^2 + b^2 \quad \text{and} \quad \langle \bar{J}'_{x, x+1} \rangle_{\mu_{\bar{\beta}, \bar{\lambda}}} = 2b(\bar{e} - b\bar{v} + 1). \quad (4)$$

Hence, in the hyperbolic scaling, the hydrodynamical equations are given by

$$\begin{cases} \partial_t \mathbf{e} - b^2 \partial_q ((\mathbf{e} - b\mathbf{v})^2) = 0 \\ \partial_t \mathbf{v} + 2b \partial_q (\mathbf{e} - b\mathbf{v}) = 0 \end{cases} \quad (5)$$

and can be written in the compact form $\partial_t \tilde{\mathbf{x}} + \partial_q \tilde{\mathbf{J}}(\tilde{\mathbf{x}}) = 0$ with

$$\tilde{\mathbf{x}} = \begin{pmatrix} \mathbf{e} \\ \mathbf{v} \end{pmatrix}, \quad \text{and} \quad \tilde{\mathbf{J}}(\tilde{\mathbf{x}}) = \begin{pmatrix} -b^2(\mathbf{e} - b\mathbf{v})^2 \\ 2b(\mathbf{e} - b\mathbf{v}) \end{pmatrix}. \quad (6)$$

This can be proved before the appearance of the shocks (see [7]). The differential matrix of $\tilde{\mathbf{J}}$ is given by

$$\nabla \tilde{\mathbf{J}}(\tilde{\mathbf{x}}) = 2b \begin{pmatrix} -b(\mathbf{e} - b\mathbf{v}) & b^2(\mathbf{e} - b\mathbf{v}) \\ 1 & -b \end{pmatrix}.$$

For given (\bar{e}, \bar{v}) we denote by $(\bar{T}_t^+)_{t \geq 0}$ (resp. $(\bar{T}_t^-)_{t \geq 0}$) the semigroup on $S(\mathbb{R}) \times S(\mathbb{R})$ generated by

$$\partial_t \varepsilon + \bar{M}^T \partial_q \varepsilon = 0, \quad (\text{resp. } \partial_t \varepsilon - \bar{M}^T \partial_q \varepsilon = 0), \quad (7)$$

where

$$\bar{M} := \bar{M}(\bar{e}, \bar{v}) = [\nabla \tilde{\mathbf{J}}](\bar{\omega}), \quad \bar{\omega} = \begin{pmatrix} \bar{e} \\ \bar{v} \end{pmatrix}.$$

We omit the dependence of these semigroups on (\bar{e}, \bar{v}) for lightness of the notations. Above $S(\mathbb{R})$ denotes the Schwartz space of smooth rapidly decreasing functions.

3. Statement of the results

For each integer $z \geq 0$, let $H_z(x) = (-1)^z e^{x^2} \frac{d^z}{dx^z} e^{-x^2}$ be the Hermite polynomial and $h_z(x) = (z! \sqrt{2\pi})^{-1} H_z(x) e^{-x^2}$ the Hermite function. The set $\{h_z, z \geq 0\}$ is an orthonormal basis of $\mathbb{L}^2(\mathbb{R})$. Consider in $\mathbb{L}^2(\mathbb{R})$ the operator $K_0 = x^2 - \Delta$, Δ being the Laplacian on \mathbb{R} . For an integer $k \geq 0$, denote by \mathbb{H}_k the Hilbert space induced by $S(\mathbb{R})$ and the scalar product $\langle \cdot, \cdot \rangle_k$ defined by $\langle f, g \rangle_k = \langle f, K_0^k g \rangle_0$, where $\langle \cdot, \cdot \rangle_0$ denotes the inner product of $\mathbb{L}^2(\mathbb{R})$ and denote by \mathbb{H}_{-k} the dual of \mathbb{H}_k , relatively to this inner product. Let $\langle \cdot \rangle$ represent the average with respect to the Lebesgue measure.

We take the infinite system at equilibrium under the Gibbs measure $\mu_{\bar{\beta}, \bar{\lambda}}$ corresponding to a mean energy \bar{e} and a mean volume \bar{v} . Our goal is to study the energy-volume fluctuation field in the time-scale $tn^{1+\alpha}$, $\alpha \geq 0$:

$$\mathcal{Y}_t^{n,\alpha}(\mathbf{G}) = \frac{1}{\sqrt{n}} \sum_{x \in \mathbb{Z}} \mathbf{G}(x/n) \cdot (\bar{\omega}_x(tn^{1+\alpha}) - \bar{\omega}), \quad (8)$$

where for $q \in \mathbb{R}$, $x \in \mathbb{Z}$,

$$\mathbf{G}(q) = \begin{pmatrix} G_1(q) \\ G_2(q) \end{pmatrix}, \quad \bar{\omega}_x = \begin{pmatrix} V_b(\eta_x) \\ \eta_x \end{pmatrix}$$

and G_1, G_2 are test functions belonging to $S(\mathbb{R})$.

If E is a Polish space then $D(\mathbb{R}^+, E)$ (resp. $C(\mathbb{R}^+, E)$) denotes the space of E -valued functions, right continuous with left limits (resp. continuous), endowed with the Skorohod (resp. uniform) topology. Let $Q^{n,\alpha}$ be the probability measure on $D(\mathbb{R}^+, \mathbb{H}_{-k} \times \mathbb{H}_{-k})$ induced by the fluctuation field $\mathcal{Y}_t^{n,\alpha}$ and $\mu_{\bar{\beta}, \bar{\lambda}}$. Let $\mathbb{P}_{\mu_{\bar{\beta}, \bar{\lambda}}}$ denote the probability measure on $D(\mathbb{R}^+, \mathbb{R}^{\mathbb{Z}})$ induced by $(\eta(t))_{t \geq 0}$ and $\mu_{\bar{\beta}, \bar{\lambda}}$. Let $\mathbb{E}_{\mu_{\bar{\beta}, \bar{\lambda}}}$ denote the expectation with respect to $\mathbb{P}_{\mu_{\bar{\beta}, \bar{\lambda}}}$.

Theorem 1 *Fix an integer $k > 2$. Denote by Q the probability measure on $C(\mathbb{R}^+, \mathbb{H}_{-k} \times \mathbb{H}_{-k})$ corresponding to a stationary Gaussian process with mean 0 and covariance given by*

$$\mathbb{E}_Q[\mathcal{Y}_t(\mathbf{H}) \mathcal{Y}_s(\mathbf{G})] = \langle \bar{T}_t^- \mathbf{H} \cdot \bar{\chi} \bar{T}_s^- \mathbf{G} \rangle$$

for every $0 \leq s \leq t$ and \mathbf{H}, \mathbf{G} in $\mathbb{H}_k \times \mathbb{H}_k$. Here $\bar{\chi} := \bar{\chi}(\bar{\beta}, \bar{\lambda})$ is the equilibrium covariance matrix³ of $\bar{\omega}_0$. Then, the sequence $(Q^{n,0})_{n \geq 1}$ converges weakly, as $n \rightarrow \infty$, to the probability measure Q .

A byproduct of Theorem 1 is a Central Limit Theorem for the energy flux and for the volume flux through a fixed bond. Despite it is not directly related to the problem of anomalous diffusion it has a probabilistic interest. For that purpose, fix a site $x \in \mathbb{Z}$, let $\mathcal{E}_{x,x+1}^n(t)$ (resp. $\mathcal{V}_{x,x+1}^n(t)$) denote the energy (resp. volume) flux through the bond $\{x, x+1\}$ during the time interval $[0, tn]$. By conservation laws, for any $x \in \mathbb{Z}$ it holds that:

$$\begin{aligned} \mathcal{E}_{x-1,x}^n(t) - \mathcal{E}_{x,x+1}^n(t) &:= V_b(\eta_x(tn)) - V_b(\eta_x(0)) \\ \left(\text{resp. } \mathcal{V}_{x-1,x}^n(t) - \mathcal{V}_{x,x+1}^n(t) &:= \eta_x(tn) - \eta_x(0) \right). \end{aligned}$$

This, together with the previous result allow us to conclude that

³ See (19) for an explicit expression.

Corollary 1. Fix $x \in \mathbb{Z}$ and let $Z_t^{n,e} := \frac{1}{\sqrt{n}}\{\mathcal{E}_{x,x+1}^n(t) - \mathbb{E}_{\mu_{\bar{\beta},\bar{\lambda}}}[\mathcal{E}_{x,x+1}^n(t)]\}$. For every $t \geq 0$, $(Z_t^{n,e})_{n \geq 1}$ converges in law in the sense of finite-dimensional distributions, as $n \rightarrow \infty$, to a Brownian motion Z_t^e with mean zero and covariance given by

$$\mathbb{E}_{\mathbb{Q}}[Z_t^e Z_s^e] = \frac{2}{\bar{\beta}^3}(\bar{\lambda} - b\bar{\beta})^2 s,$$

for all $s \leq t$.

Corollary 2. Fix $x \in \mathbb{Z}$ and let $Z_t^{n,v} := \frac{1}{\sqrt{n}}\{\mathcal{V}_{x,x+1}^n(t) - \mathbb{E}_{\mu_{\bar{\beta},\bar{\lambda}}}[\mathcal{V}_{x,x+1}^n(t)]\}$. For every $t \geq 0$, $(Z_t^{n,v})_{n \geq 1}$ converges in law in the sense of finite-dimensional distributions, as $n \rightarrow \infty$, to a Brownian motion Z_t^v with mean zero and covariance given by

$$\mathbb{E}_{\mathbb{Q}}[Z_t^v Z_s^v] = \frac{2}{\bar{\beta}} s,$$

for all $s \leq t$.

We notice that, according to Corollary 1, the limiting energy flux Z_t^e has a vanishing variance for $\bar{\lambda} = b\bar{\beta}$ which is equivalent to $\bar{e} = b\bar{v}$. Last equivalence is a consequence of (17) and (16).

The theorem above means that in the hyperbolic scaling the fluctuations are trivial: the initial fluctuations are transported by the linearized system of (5). To see a nontrivial behavior we have to study, in the transport frame, the fluctuations at a longer time scale $tn^{1+\alpha}$, with $\alpha > 0$. Thus, we consider the fluctuation field $\hat{\mathcal{Y}}_t^{n,\alpha}$, $\alpha > 0$, defined, for any $\mathbf{G} \in S(\mathbb{R}) \times S(\mathbb{R})$, by

$$\hat{\mathcal{Y}}_t^{n,\alpha}(\mathbf{G}) = \mathcal{Y}_t^{n,\alpha}(\bar{T}_{tn^\alpha}^+ \mathbf{G}). \quad (9)$$

According to the fluctuating hydrodynamics theory ([28], pp. 85-96), in the case of a normal (diffusive) behavior $\alpha = 1$, the field $(\hat{\mathcal{Y}}_t^{n,\alpha})_{t \geq 0}$ should converge to the stationary field $(\hat{\mathcal{Y}}_t)_{t \geq 0}$ simply related to the solution $(\hat{\mathcal{Z}}_t)_{t \geq 0}$ of the linear two dimensional vector valued (infinite-dimensional) stochastic partial differential equation

$$\partial_t \hat{\mathcal{Z}}_t = \nabla \cdot \left(\mathcal{D} \nabla \hat{\mathcal{Z}}_t \right) + \sqrt{2\mathcal{D}\bar{\chi}} \nabla \cdot W_t. \quad (10)$$

Here W_t is a standard two-dimensional vector valued space-time white noise and the coefficient $\mathcal{D} := \mathcal{D}(\bar{e}, \bar{v})$ is expressed by a Green-Kubo formula (see (12)). As above, let $\hat{\mathcal{Q}}^{n,\alpha}$ be the probability measure on $D(\mathbb{R}^+, \mathbb{H}_{-k} \times \mathbb{H}_{-k})$ induced by the fluctuation field $\hat{\mathcal{Y}}_t^{n,\alpha}$ and $\mu_{\bar{\beta},\bar{\lambda}}$. Our second main theorem shows that the correct scaling exponent α is greater or equal than $1/3$:

Theorem 2 Fix an integer $k > 1$ and $\alpha < 1/3$. Denote by Q the probability measure on $C(\mathbb{R}^+, \mathbb{H}_{-k} \times \mathbb{H}_{-k})$ corresponding to a stationary Gaussian process with mean 0 and covariance given by

$$\mathbb{E}_Q[\mathcal{Y}_t(\mathbf{H}) \mathcal{Y}_s(\mathbf{G})] = \langle \mathbf{H} \cdot \bar{\chi} \mathbf{G} \rangle$$

for every $0 \leq s \leq t$ and \mathbf{H}, \mathbf{G} in $\mathbb{H}_k \times \mathbb{H}_k$. Then, the sequence $(\hat{\mathcal{Q}}^{n,\alpha})_{n \geq 1}$ converges weakly, as $n \rightarrow \infty$, to the probability measure Q .

As in the hyperbolic time scale from the previous result we obtain limiting results for the energy flux and volume flux. In this case, we need to define the energy and volume flux through the time dependent bond $\{u_t^{x,\alpha}(n), u_t^{x,\alpha}(n) + 1\}$, where $u_t^{x,\alpha}(n) := \lfloor x - \frac{2b\lambda}{\beta}tn^{1+\alpha} \rfloor$ and $\lfloor u \rfloor$ denotes the biggest integer number smaller or equal to u . The justification for taking this reference frame with precisely this velocity will be given ahead in Remark 1. Now, fix a site $x \in \mathbb{Z}$ and let $\mathcal{E}_{u_t^{x,\alpha}(n)}^n$ (resp. $\mathcal{V}_{u_t^{x,\alpha}(n)}^n(t)$) denote the energy (resp. volume) flux through the bond $\{u_t^{x,\alpha}(n), u_t^{x,\alpha}(n) + 1\}$ during the time interval $[0, tn^{1+\alpha}]$. Then, from the previous result we conclude that

Corollary 3. *Fix $t \geq 0$, $x \in \mathbb{Z}$ and $\alpha < 1/3$. Then*

$$\lim_{n \rightarrow \infty} \mathbb{E}_{\mu_{\bar{\beta}, \bar{\lambda}}} \left[\left(\frac{1}{\sqrt{n}} \left\{ \mathcal{E}_{u_t^{x,\alpha}(n)}^n(t) - \mathbb{E}_{\mu_{\bar{\beta}, \bar{\lambda}}} [\mathcal{E}_{u_t^{x,\alpha}(n)}^n(t)] \right\} \right)^2 \right] = 0.$$

and

$$\lim_{n \rightarrow \infty} \mathbb{E}_{\mu_{\bar{\beta}, \bar{\lambda}}} \left[\left(\frac{1}{\sqrt{n}} \left\{ \mathcal{V}_{u_t^{x,\alpha}(n)}^n(t) - \mathbb{E}_{\mu_{\bar{\beta}, \bar{\lambda}}} [\mathcal{V}_{u_t^{x,\alpha}(n)}^n(t)] \right\} \right)^2 \right] = 0.$$

Similar results have been obtained in [18] by one of the authors for the asymmetric simple exclusion. The proof of Corollaries 1, 2 and 3 follows the same arguments as in [18] once the previous theorems are proved. For that reason we will only give a sketch of their proof. The proof of the theorems is more problematic since the multi-scale analysis performed in [18] relies crucially on the existence of a spectral gap so that we cannot follow [18]. Therefore we propose an alternative approach based on computations of some resolvent norms.

Theorem 2 does not exclude the possibility of normal fluctuations, i.e. $\alpha = 1$. In order to show that the system we consider is really superdiffusive we will show that the transport coefficient \mathcal{D} which appears in (10) is infinite so that the correct scaling exponent α is strictly smaller than 1. Our third result, stated below, shows it is in fact less than $3/4$.

With the notations introduced in the previous section, the *normalized* currents are defined by

$$\hat{J}_{x,x+1}(\eta) = \begin{pmatrix} \bar{j}_{x,x+1}(\eta) \\ \bar{j}'_{x,x+1}(\eta) \end{pmatrix} - \bar{\mathfrak{J}}(\bar{\omega}) - (\nabla \bar{\mathfrak{J}})(\bar{\omega}) \begin{pmatrix} V_b(\eta_x) - \bar{v} \\ \eta_x - \bar{v} \end{pmatrix}. \quad (11)$$

Up to a constant matrix coming from a martingale term (due to the noise) and thus irrelevant for us (see [2], [7]), the coefficient \mathcal{D} is defined by the Green-Kubo formula

$$\mathcal{D} = \int_0^\infty C(t) dt, \quad (12)$$

where

$$C(t) := \mathbb{E}_{\mu_{\bar{\beta}, \bar{\lambda}}} \left[\sum_{x \in \mathbb{Z}} \hat{J}_{x,x+1}(\eta(t)) \left[\hat{J}_{0,1}(\eta(0)) \right]^T \right]$$

is the current-current correlation function. The signature of the superdiffusive behavior of the system is seen in the divergence of the integral defining \mathcal{D} , i.e. in a slow decay of the current-current correlation function. We introduce the

Laplace transform function $\mathcal{F}(\gamma, \cdot)$ of the current-current correlation function. It is defined, for any $z > 0$ by

$$\mathcal{F}(\gamma, z) = \int_0^\infty e^{-zt} C(t) dt.$$

Our third theorem is the following lower bound on $\mathcal{F}(\gamma, z)$. Observe that $\mathcal{F}(\gamma, z)$ is a square matrix of size 2 whose (i, j) -th entry is denoted by $\mathcal{F}_{i,j}$.

Theorem 3 *Fix $\gamma > 0$. For any $(i, j) \neq (1, 1)$ and any $z > 0$ we have*

$$\mathcal{F}_{i,j}(\gamma, z) = 0.$$

There exists a positive constant $c := c(\gamma) > 0$ such that for any $z > 0$,

$$\mathcal{F}_{1,1}(\gamma, z) \geq cz^{-1/4}.$$

Moreover, there exists a positive constant $C := C(\gamma)$ such that for any $z > 0$,

$$C^{-1}\mathcal{F}_{1,1}(1, z/\gamma) \leq \mathcal{F}_{1,1}(\gamma, z) \leq C\mathcal{F}_{1,1}(1, z/\gamma). \quad (13)$$

The lower bound $\mathcal{F}_{1,1}(\gamma, z) \geq cz^{1/4}$ means roughly that the current-current correlation function $C(t)$ is bounded by below by a constant times $t^{-3/4}$. The last part of the theorem is easy to prove but has an important consequence. In [7] numerical simulations are performed to detect the anomalous diffusion of energy. Since it is difficult to estimate numerically the time autocorrelation functions of the currents because of their expected long-time tails, a more tenable approach consists in studying a non equilibrium system in its steady state, i.e. considering a finite system in contact with two thermostats which fix the value of the energy at the boundaries. Then we estimate the dependence of the energy transport coefficient $\kappa(N)$ with the system size N . The latter is defined as N times the average energy current. It turns out that $\kappa(N) \sim N^\delta$ with a parameter $\delta := \delta(\gamma) > 0$ increasing with the noise intensity γ (except for the singular value $\delta = 1$ when $\gamma = 0$ which is a manifestation of the ballistic behavior of the Kac-van Moerbeke system). This result is very surprising since the more stochasticity in the model is introduced, the less the system is diffusive. The same has been observed for other anharmonic potentials in [7] and also for the Toda lattice perturbed by an energy-momentum conservative noise ([19]). It has been argued in [19] that this may be explained by the fact that some diffusive phenomena due to non-linearities, like localized breathers, are destroyed by the noise. In [3] simulations have been performed directly with the Green-Kubo formula for other standard anharmonic chains with the same conclusion: current-current correlation function decreases slower when the noise intensity increases. If all these numerical simulations reproduce correctly the real behavior of the models investigated, they dismiss the theories which pretend that some universality holds, e.g. [31]. It is therefore very important to decide if the phenomena numerically observed are correct or not.

Assuming that the current-current correlation function $C(t)$ has the time decay $C(t) \sim_{t \rightarrow \infty} t^{-\delta'(\gamma)}$, the inequality (13) shows that the exponent $\delta' := \delta'(\gamma)$ is independent of γ (up to possible slowly varying functions corrections, i.e. in a Tauberian sense). It is usually argued but not proved (see e.g. the end of Section 5.3 in [24]) that the exponent δ defined by the non-equilibrium stationary state

is related to δ' by the relation $\delta = 1 - \delta'$, and is consequently independent of γ too. Therefore the numerical simulations do not seem to reflect the correct behavior of the system ⁴. A possible explanation of the inconsistency between the numerical observation and our result is simply that the relation $\delta = 1 - \delta'$ is not satisfied. Nevertheless, notice that the last part of our theorem is in fact valid for all the models cited above. It applies in particular to the models studied in [3] and shows that the numerical observations of that paper, which are performed for the Green-Kubo formula, are not consistent with the real behavior of the system.

4. A change of variables

To study the energy-volume fluctuation field $\mathcal{Y}^{n,\alpha}$, we introduce the following change of variables $\xi_x = e^{-b\eta_x}$, for each $x \in \mathbb{Z}$. Then, the previous Markovian system $(\eta(t))_{t \geq 0}$ defines a new Markovian system $(\xi(t))_{t \geq 0}$ with state space $(0, +\infty)^{\mathbb{Z}}$ whose generator \mathcal{L} is equal to $b^2\mathcal{A} + \gamma\mathcal{S}$, where for local differentiable functions $f : (0, +\infty)^{\mathbb{Z}} \rightarrow \mathbb{R}$ we have that

$$(\mathcal{A}f)(\xi) = \sum_{x \in \mathbb{Z}} \xi_x (\xi_{x+1} - \xi_{x-1}) (\partial_{\xi_x} f)(\xi)$$

and

$$(\mathcal{S}f)(\xi) = \sum_{x \in \mathbb{Z}} \left(f(\xi^{x,x+1}) - f(\xi) \right),$$

where $\xi^{x,x+1}$ is defined as in (3).

Observe that the energy and volume conservation laws correspond, for the process $(\xi(t))_{t \geq 0}$, to the conservation of the two following quantities $\sum_{x \in \mathbb{Z}} \xi_x$ and $\sum_{x \in \mathbb{Z}} \log(\xi_x)$. The corresponding microscopic currents are defined by the conservation law equations:

$$\mathcal{L}(\xi_x) = -\nabla j_{x-1,x}(\xi), \quad \mathcal{L}(\log \xi_x) = -\nabla j'_{x-1,x}(\xi),$$

where

$$j_{x,x+1}(\xi) = -b^2 \xi_x \xi_{x+1} - \gamma \nabla \xi_x,$$

and

$$j'_{x,x+1}(\xi) = -b^2 (\xi_x + \xi_{x+1}) - \gamma \nabla \log(\xi_x).$$

We will use the compact notation

$$J_{x,x+1}(\xi) = \begin{pmatrix} j_{x,x+1}(\xi) \\ j'_{x,x+1}(\xi) \end{pmatrix}. \quad (14)$$

Since $V_b(\eta_x) = \xi_x - \log(\xi_x) + 1$ and $\eta_x = -\frac{1}{b} \log(\xi_x)$, we have the following relations between the microscopic currents

$$\bar{j}_{x,x+1}(\eta) = j_{x,x+1}(\xi) - j'_{x,x+1}(\xi), \quad \text{and} \quad \bar{j}'_{x,x+1}(\eta) = -\frac{1}{b} j'_{x,x+1}(\xi). \quad (15)$$

⁴ It would be very interesting to understand why the numerical simulations are so sensitive to the noise.

If η is distributed according to $\mu_{\bar{\beta}, \bar{\lambda}}$ then ξ defined by $\xi_x = e^{-b\eta_x}$ is distributed according to the probability measure $\nu_{\beta, \lambda}$ on $(0, +\infty)^{\mathbb{Z}}$ given by

$$\nu_{\beta, \lambda}(d\xi) = \prod_{x \in \mathbb{Z}} Z^{-1}(\beta, \lambda) \mathbf{1}_{\{\xi_x > 0\}} \exp\{-\beta \xi_x + \lambda \log(\xi_x)\} d\xi_x$$

with $Z(\beta, \lambda)$ the partition function,

$$\beta = \bar{\beta}, \quad \text{and} \quad \lambda = -1 + \bar{\lambda}/b. \quad (16)$$

Remark that $\nu_{\beta, \lambda}$ is nothing but a product probability measure whose marginal follows a Gamma distribution $\gamma_{\lambda+1, \beta^{-1}}$ with parameter $(\lambda + 1, \beta^{-1})$. In particular, we have $Z := Z(\beta, \lambda) = \beta^{-(\lambda+1)} \Gamma(\lambda + 1)$, where Γ is the usual Gamma function.

Thus, the process $(\xi(t))_{t \geq 0}$ has a family of translation invariant measures $\nu_{\beta, \lambda}$ parameterized by the chemical potentials $(\beta, \lambda) \in (0, +\infty) \times (-1, +\infty)$.

Let $\mathbb{P}_{\nu_{\beta, \lambda}}$ be the probability measure on $D(\mathbb{R}^+, (0, +\infty)^{\mathbb{Z}})$ induced by $(\xi(t))_{t \geq 0}$ and $\nu_{\beta, \lambda}$ and let $\mathbb{E}_{\nu_{\beta, \lambda}}$ denote the expectation with respect to $\mathbb{P}_{\nu_{\beta, \lambda}}$.

Let $\langle \cdot \rangle_{\nu_{\beta, \lambda}}$ denote the average with respect to $\nu_{\beta, \lambda}$. The averages $\rho := \rho(\beta, \lambda)$ and $\theta := \theta(\beta, \lambda)$ of the conserved quantities for $(\xi(t))_{t \geq 0}$ at equilibrium under $\nu_{\beta, \lambda}$ are defined by $\rho = \langle \xi_x \rangle_{\nu_{\beta, \lambda}}$ and $\theta = \langle \log(\xi_x) \rangle_{\nu_{\beta, \lambda}}$. By a direct computation we get

$$\rho = 1 + \bar{e} - b\bar{v} = \frac{\lambda + 1}{\beta}, \quad \text{and} \quad \theta = -b\bar{v} = \frac{\Gamma'(\lambda + 1)}{\Gamma(\lambda + 1)} - \log(\beta). \quad (17)$$

It is understood, here and in the whole paper, that (β, λ) are related to $(\bar{\beta}, \bar{\lambda})$ through (16). We will use the following compact notation, for each $x \in \mathbb{Z}$,

$$\omega_x = \begin{pmatrix} \xi_x \\ \log(\xi_x) \end{pmatrix}, \quad \text{and} \quad \omega = \begin{pmatrix} \rho \\ \theta \end{pmatrix}.$$

Observe that $\bar{\omega}_x = \Lambda \omega_x - \begin{pmatrix} 1 \\ 0 \end{pmatrix}$, where

$$\Lambda = \begin{pmatrix} 1 & -1 \\ 0 & -1/b \end{pmatrix}. \quad (18)$$

The covariance matrix $\chi := \chi(\beta, \lambda)$ of ω_0 under $\nu_{\beta, \lambda}$ is given by

$$\chi = \begin{pmatrix} \langle (\xi_0 - \rho)^2 \rangle_{\nu_{\beta, \lambda}} & \langle (\xi_0 - \rho)(\log(\xi_0) - \theta) \rangle_{\nu_{\beta, \lambda}} \\ \langle (\xi_0 - \rho)(\log(\xi_0) - \theta) \rangle_{\nu_{\beta, \lambda}} & \langle (\log(\xi_0) - \theta)^2 \rangle_{\nu_{\beta, \lambda}} \end{pmatrix}.$$

A simple computation shows that

$$\chi = \begin{pmatrix} \frac{\lambda+1}{\beta^2} & \frac{1}{\beta} \\ \frac{1}{\beta} & (\log \Gamma)'(\lambda + 1) \end{pmatrix} = \begin{pmatrix} \partial_{\beta}^2 \log(Z) & -\partial_{\beta, \lambda} \log(Z) \\ -\partial_{\beta, \lambda} \log(Z) & \partial_{\lambda}^2 \log(Z) \end{pmatrix}.$$

Denote the covariance matrix of $\bar{\omega}_0$ under $\mu_{\bar{\beta}, \bar{\lambda}}$ by $\bar{\chi} := \bar{\chi}(\bar{\beta}, \bar{\lambda})$, which is defined by

$$\bar{\chi} = \begin{pmatrix} \langle (V_b(\eta_0) - \bar{e})^2 \rangle_{\mu_{\bar{\beta}, \bar{\lambda}}} & \langle (V_b(\eta_0) - \bar{e})(\eta_0 - \bar{v}) \rangle_{\mu_{\bar{\beta}, \bar{\lambda}}} \\ \langle (V_b(\eta_0) - \bar{e})(\eta_0 - \bar{v}) \rangle_{\mu_{\bar{\beta}, \bar{\lambda}}} & \langle (\eta_0 - \bar{v})^2 \rangle_{\mu_{\bar{\beta}, \bar{\lambda}}} \end{pmatrix}.$$

Thus, the covariance matrix χ of ω_0 under $\nu_{\beta,\lambda}$ is related to the covariance matrix $\bar{\chi}$ of $\bar{\omega}_0$ under $\mu_{\bar{\beta},\bar{\lambda}}$, by

$$\bar{\chi} = \Lambda \chi \Lambda^T = \begin{pmatrix} \frac{\lambda+1}{\beta^2} + \frac{2}{\beta} + (\log \Gamma)''(\lambda+1) & \frac{1}{b\beta} + \frac{(\log \Gamma)''(\lambda+1)}{b} \\ \frac{1}{b\beta} + \frac{(\log \Gamma)''(\lambda+1)}{b} & \frac{(\log \Gamma)''(\lambda+1)}{b^2} \end{pmatrix}. \quad (19)$$

A simple computation shows that $\langle j_{x,x+1} \rangle_{\nu_{\beta,\lambda}} = -b^2 \rho^2$ and $\langle j'_{x,x+1} \rangle_{\nu_{\beta,\lambda}} = -2b^2 \rho$. The hydrodynamical equations for the process $(\xi(t))_{t \geq 0}$ are given by

$$\begin{cases} \partial_t \rho - b^2 \partial_q (\rho^2) = 0 \\ \partial_t \theta - 2b^2 \partial_q \rho = 0 \end{cases} \quad (20)$$

and can be written in the compact form $\partial_t \mathfrak{X} + \partial_q \mathfrak{J}(\mathfrak{X}) = 0$ with

$$\mathfrak{X} = \begin{pmatrix} \rho \\ \theta \end{pmatrix}, \quad \text{and} \quad \mathfrak{J}(\mathfrak{X}) = \begin{pmatrix} -b^2 \rho^2 \\ -2b^2 \rho \end{pmatrix}.$$

The differential matrix of \mathfrak{J} is given by

$$\nabla \mathfrak{J}(\mathfrak{X}) = \begin{pmatrix} -2b^2 \rho & 0 \\ -2b^2 & 0 \end{pmatrix}.$$

As above, let $(T_t^+)_{t \geq 0}$ (resp. $(T_t^-)_{t \geq 0}$) denote the semigroup on $S(\mathbb{R}) \times S(\mathbb{R})$ generated by

$$\partial_t \varepsilon + M^T \partial_q \varepsilon = 0, \quad (\text{resp. } \partial_t \varepsilon - M^T \partial_q \varepsilon = 0). \quad (21)$$

where

$$M := M(\rho, \theta) = (\nabla \mathfrak{J})(\omega),$$

ρ and θ are given by (17). We omit the dependence of these semigroups on (ρ, θ) for lightness of the notations.

We remark that the transposed linearized system of (20) around the constant profiles (ρ, θ) is given by the first equation on the left hand side of (21). It is easy to show that $\bar{M} = \Lambda M \Lambda^{-1}$ and $\Lambda^T \bar{T}_t^- = T_t^- \Lambda^T$.

5. Orthogonal decomposition

Observe that $\nu_{\beta,\lambda}$ is a product of Gamma distributions. Let us recall that the Gamma distribution $\gamma_{\alpha,k}$ with parameter (α, k) is the probability distribution on $(0, +\infty)$ absolutely continuous with respect to the Lebesgue measure with density $f_{\alpha,k}$ given by

$$f_{\alpha,k}(q) = \left(k^\alpha \Gamma(\alpha) \right)^{-1} q^{\alpha-1} e^{-q/k}, \quad q > 0. \quad (22)$$

Thus, we have $\nu_{\beta,\lambda}(d\xi) = \prod_{x \in \mathbb{Z}} (f_{\lambda+1,\beta^{-1}}(\xi_x) d\xi_x) = \prod_{x \in \mathbb{Z}} (\beta f_{\lambda+1,1}(\beta \xi_x) d\xi_x)$.

The generalized Laguerre polynomials $(H_n^{(\lambda)})_{n \geq 0}$ form an orthogonal basis of the space $\mathbb{L}^2(\gamma_{\lambda+1,1})$. They satisfy the following equations:

$$\begin{aligned} H_0^{(\lambda)} &= 1, \\ q \frac{d}{dq} H_n^{(\lambda)} &= n H_n^{(\lambda)} - (n + \lambda) H_{n-1}^{(\lambda)}, \\ \left(q \frac{d^2}{dq^2} + (\lambda + 1 - q) \frac{d}{dq} + n \right) H_n^{(\lambda)} &= 0, \\ (n + 1) H_{n+1}^{(\lambda)}(q) &= (2n + 1 + \lambda - q) H_n^{(\lambda)}(q) - (n + \lambda) H_{n-1}^{(\lambda)}(q) \end{aligned} \quad (23)$$

and the normalization condition

$$\int_0^\infty \left(H_n^{(\lambda)}(q) \right)^2 f_{\lambda+1,1}(q) dq = \frac{\Gamma(\lambda + n + 1)}{\Gamma(\lambda + 1)} \frac{1}{n!}.$$

In particular, we have

$$\begin{aligned} H_1^{(\lambda)}(q) &= -q + (\lambda + 1), \\ H_2^{(\lambda)}(q) &= \frac{(2 + \lambda)(1 + \lambda)}{2} - (\lambda + 2)q + \frac{q^2}{2}. \end{aligned} \quad (24)$$

Let Σ be the set composed of configurations $\sigma = (\sigma_x)_{x \in \mathbb{Z}} \in \mathbb{N}^{\mathbb{Z}}$ such that $\sigma_x \neq 0$ only for a finite number of x . The number $\sum_{x \in \mathbb{Z}} \sigma_x$ is called the size of σ and is denoted by $|\sigma|$. Let $\Sigma_n = \{\sigma \in \Sigma ; |\sigma| = n\}$. On the set of n -tuples $\mathbf{x} := (x_1, \dots, x_n)$ of \mathbb{Z}^n , we introduce the equivalence relation $\mathbf{x} \sim \mathbf{y}$ if there exists a permutation p on $\{1, \dots, n\}$ such that $x_{p(i)} = y_i$ for all $i \in \{1, \dots, n\}$. The class of \mathbf{x} for the relation \sim is denoted by $[\mathbf{x}]$ and its cardinal by $c(\mathbf{x})$. Then the set of configurations of Σ_n can be identified with the set of n -tuples classes for \sim by the one-to-one application:

$$[\mathbf{x}] = [(x_1, \dots, x_n)] \in \mathbb{Z}^n / \sim \rightarrow \sigma^{[\mathbf{x}]} \in \Sigma_n$$

where for any $y \in \mathbb{Z}$, $(\sigma^{[\mathbf{x}]})_y = \sum_{i=1}^n \mathbf{1}_{y=x_i}$. We will identify $\sigma \in \Sigma_n$ with the occupation numbers of a configuration with n particles, and $[\mathbf{x}]$ will correspond to the positions of those n particles.

To any $\sigma \in \Sigma$, we associate the polynomial function $H_\sigma^{\beta,\lambda}$ given by

$$H_\sigma^{\beta,\lambda}(\xi) = \prod_{x \in \mathbb{Z}} H_{\sigma_x}^{(\lambda)}(\beta \xi_x).$$

Then, the family $\{H_\sigma^{\beta,\lambda} ; \sigma \in \Sigma\}$ forms an orthogonal basis of $\mathbb{L}^2(\nu_{\beta,\lambda})$ such that

$$\int H_\sigma^{\beta,\lambda} H_{\sigma'}^{\beta,\lambda} d\nu_{\beta,\lambda} = \delta_{\sigma=\sigma'} \prod_{x \in \mathbb{Z}} \frac{\Gamma(\lambda + \sigma_x + 1)}{\Gamma(\lambda + 1)} \frac{1}{\sigma_x!} = \delta_{\sigma=\sigma'} \mathcal{W}^\lambda(\sigma), \quad (25)$$

where

$$\mathcal{W}^\lambda(\sigma) := \prod_{x \in \mathbb{Z}} \frac{\Gamma(\lambda + \sigma_x + 1)}{\Gamma(\lambda + 1)} \frac{1}{\sigma_x!} \quad (26)$$

and δ denotes the Kronecker function, so that $\delta_{\sigma=\sigma'} = 1$ if $\sigma = \sigma'$, otherwise it is equal to zero.

A function $F : \Sigma \rightarrow \mathbb{R}$ such that $F(\sigma) = 0$ if $\sigma \notin \Sigma_n$ is called a degree n function. Thus, such a function is sometimes considered as a function defined only on Σ_n . A local function $f \in \mathbb{L}^2(\nu_{\beta,\lambda})$ whose decomposition on the orthogonal basis $\{H_\sigma^{\beta,\lambda}; \sigma \in \Sigma\}$ is given by $f = \sum_\sigma F(\sigma)H_\sigma^{\beta,\lambda}$ is called of degree n if and only if F is of degree n . A function $F : \Sigma_n \rightarrow \mathbb{R}$ is nothing but a symmetric function $F : \mathbb{Z}^n \rightarrow \mathbb{R}$ through the identification of σ with $[\mathbf{x}]$. We denote by $\langle \cdot, \cdot \rangle$ the scalar product on $\oplus \mathbb{L}^2(\Sigma_n)$, each Σ_n being equipped with the counting measure. Hence, if $F, G : \Sigma \rightarrow \mathbb{R}$, we have

$$\langle F, G \rangle = \sum_{n \geq 0} \sum_{\sigma \in \Sigma_n} F_n(\sigma)G_n(\sigma) = \sum_{n \geq 0} \sum_{\mathbf{x} \in \mathbb{Z}^n} \frac{1}{c(\mathbf{x})} F_n(\mathbf{x})G_n(\mathbf{x}),$$

with F_n, G_n the restrictions of F, G to Σ_n . We recall that $c(\mathbf{x})$ is the cardinal of $[\mathbf{x}]$. Since (β, λ) are fixed through the paper we denote $H_\sigma^{\beta,\lambda}$ by H_σ and $\mathcal{W}^\lambda(\sigma)$ by $\mathcal{W}(\sigma)$.

If a local function $f \in \mathbb{L}^2(\nu_{\beta,\lambda})$ is written in the form $f(\xi) = \sum_{\sigma \in \Sigma} F(\sigma)H_\sigma(\xi)$ then we have

$$(\mathcal{A}f)(\xi) = \sum_{\sigma \in \Sigma} (\mathfrak{A}F)(\sigma)H_\sigma(\xi), \quad (\mathcal{S}f)(\xi) = \sum_{\sigma \in \Sigma} (\mathfrak{S}F)(\sigma)H_\sigma(\xi)$$

with

$$(\mathfrak{S}F)(\sigma) = \sum_{x \in \mathbb{Z}} (F(\sigma^{x,x+1}) - F(\sigma)),$$

where $\sigma^{x,x+1}$ is obtained from σ by exchanging the occupation numbers σ_x and σ_{x+1} .

Let us now compute the operator \mathfrak{A} . We have

$$(\mathcal{A}H_\sigma)(\xi) = \sum_{x \in \mathbb{Z}} \xi_x (\xi_{x+1} - \xi_{x-1}) \partial_{\xi_x} H_\sigma(\xi).$$

By the definition of H_σ and by the second equality in (23), it follows that

$$(\mathcal{A}H_\sigma)(\xi) = \beta \sum_{x \in \mathbb{Z}} (\xi_{x+1} - \xi_{x-1}) \left(\sigma_x H_\sigma(\xi) - (\sigma_x + \lambda) H_{\sigma - \delta_x}(\xi) \right),$$

where $\sigma - \delta_x$ is the configuration where a particle has been deleted at site x (if there was no particle on site x , then $\sigma - \delta_x = \sigma$).

Now, noticing that the fourth equality in (23) can be written as

$$\beta q H_n^{(\lambda)}(\beta q) = (2n + 1 + \lambda) H_n^{(\lambda)}(\beta q) - (n + \lambda) H_{n-1}^{(\lambda)}(\beta q) - (n + 1) H_{n+1}^{(\lambda)}(\beta q)$$

and performing some change of variables, we have that

$$\begin{aligned} (\mathcal{A}H_\sigma)(\xi) &= \sum_{\substack{x, y \in \mathbb{Z} \\ |x-y|=1}} a(y-x)(\sigma_x + \lambda)(\sigma_y + 1) H_{\sigma + \delta_y - \delta_x}(\xi) \\ &\quad - \sum_{x \in \mathbb{Z}} (\sigma_x + \lambda)(\sigma_{x+1} - \sigma_{x-1}) H_{\sigma - \delta_x}(\xi) \\ &\quad + \sum_{x \in \mathbb{Z}} (\sigma_x + 1)(\sigma_{x+1} - \sigma_{x-1}) H_{\sigma + \delta_x}(\xi). \end{aligned}$$

Here, $a(z) = -1$ if $z = -1$, $a(z) = 1$ if $z = 1$ and 0 otherwise. It follows that

$$\mathfrak{A} = \mathfrak{A}_0 + \mathfrak{A}_- + \mathfrak{A}_+$$

with

$$\begin{aligned} (\mathfrak{A}_0 F)(\sigma) &= - \sum_{\substack{x, y \in \mathbb{Z} \\ |x-y|=1}} a(y-x) \sigma_x (\sigma_y + 1 + \lambda) F(\sigma + \delta_y - \delta_x), \\ (\mathfrak{A}_+ F)(\sigma) &= - \sum_{x \in \mathbb{Z}} \sigma_x (\sigma_{x+1} - \sigma_{x-1}) F(\sigma - \delta_x), \\ (\mathfrak{A}_- F)(\sigma) &= \sum_{x \in \mathbb{Z}} (\sigma_x - 1 + \lambda) (\sigma_{x+1} - \sigma_{x-1}) F(\sigma + \delta_x). \end{aligned}$$

Observe that if F vanishes outside of Σ_n then $\mathfrak{A}_\pm F$ vanishes outside of $\Sigma_{n \mp 1}$ and \mathfrak{A}_0 vanishes outside of Σ_n . In other words, \mathfrak{A}_0 keeps fixed the degree of a function, \mathfrak{A}_+ raises the degree by one while \mathfrak{A}_- lowers the degree by one.

The Dirichlet form $\mathcal{D}(f)$ of a local function $f \in \mathbb{L}^2(\nu_{\beta, \lambda})$ is defined by

$$\mathcal{D}(f) = \langle f, (-\mathcal{S}f) \rangle_{\nu_{\beta, \lambda}} = \frac{1}{2} \sum_{x \in \mathbb{Z}} \int (f(\xi^{x, x+1}) - f(\xi))^2 \nu_{\beta, \lambda}(d\xi).$$

Recall that $\langle \cdot, \cdot \rangle_{\nu_{\beta, \lambda}}$ denotes the inner product of $\mathbb{L}^2(\nu_{\beta, \lambda})$.

Since f has the decomposition $f = \sum_{\sigma \in \Sigma} F(\sigma) H_\sigma$ then

$$\mathcal{D}(f) = \frac{1}{2} \sum_{x \in \mathbb{Z}} \sum_{\sigma \in \Sigma} \mathcal{W}(\sigma) (F(\sigma^{x, x+1}) - F(\sigma))^2. \quad (27)$$

Let $\Delta_+ = \{(x, y) \in \mathbb{Z}^2; y \geq x + 1\}$, $\Delta_- = \{(x, y) \in \mathbb{Z}^2; y \leq x - 1\}$ and $\Delta_0 = \{(x, x); x \in \mathbb{Z}\}$. We denote by \mathbb{D}_1 the Dirichlet form of a symmetric simple one dimensional random walk, i.e.

$$\mathbb{D}_1(F) = \frac{1}{2} \sum_{x \in \mathbb{Z}} (F(x+1) - F(x))^2,$$

where $F : \mathbb{Z} \rightarrow \mathbb{R}$ is such that $\sum_{x \in \mathbb{Z}} F^2(x) < \infty$.

We denote by \mathbb{D}_2 the Dirichlet form of a symmetric simple random walk on \mathbb{Z}^2 where jumps from Δ_\pm to Δ_0 and from Δ_0 to Δ_\pm have been suppressed and jumps from $(x, x) \in \Delta_0$ to $(x \pm 1, x \pm 1) \in \Delta_0$ have been added, i.e.

$$\mathbb{D}_2(F) = \frac{1}{2} \sum_{|\mathbf{e}|=1} \sum_{\mathbf{x} \in \Delta_\pm, \mathbf{x}+\mathbf{e} \in \Delta_\pm} (F(\mathbf{x} + \mathbf{e}) - F(\mathbf{x}))^2 + \frac{1}{2} \sum_{\mathbf{x} \in \Delta_0} (F(\mathbf{x} \pm (1, 1)) - F(\mathbf{x}))^2,$$

where $F : \mathbb{Z}^2 \rightarrow \mathbb{R}$ is a symmetric function such that $\sum_{\mathbf{x} \in \mathbb{Z}^2} F^2(\mathbf{x}) < \infty$.

Lemma 1. *Let $f = \sum_{n=1}^2 \sum_{\sigma \in \Sigma_n} F_n(\sigma) H_\sigma$ be a local function such that F_1 (resp. F_2) is of degree 1 (resp. degree 2). There exists a positive constant $C := C(\lambda)$, independent of f , such that*

$$C^{-1} [\mathbb{D}_1(F_1) + \mathbb{D}_2(F_2)] \leq \mathcal{D}(f) \leq C [\mathbb{D}_1(F_1) + \mathbb{D}_2(F_2)].$$

Proof. Observe that

- If $\sigma \in \Sigma_1$, then $\mathcal{W}(\sigma) = (\lambda + 1)$.
- If $\sigma \in \Sigma_2$, $\sigma = \delta_x + \delta_y$, $x \neq y$, then $\mathcal{W}(\sigma) = (\lambda + 1)^2$; if $\sigma \in \Sigma_2$, $\sigma = 2\delta_x$, then $\mathcal{W}(\sigma) = [(\lambda + 2)(\lambda + 1)]/2$.

This follows from the relation $\Gamma(z + 1) = z\Gamma(z)$. Then, by using (27) and the identification of functions $F : \Sigma_n \rightarrow \mathbb{R}$ of degree n with their representations as symmetric functions on \mathbb{Z}^n , the claim follows.

6. Triviality of the fluctuations

In this section we prove Theorems 1 and 2 and Corollaries 1, 2 and 3 above. The proof of Theorems 1 and 2 is standard and relies on a careful analysis of martingales associated to the respective density fields. For this reason we present only the sketch of their proofs. For the interested reader we refer to chapter 11 of [22]. We notice that the restrictions on k appearing in the statement of those theorems come from tightness estimates, that we do not prove here since they follow from very similar computations to those presented in [22].

To approach the proof of theorems we notice that since $V_b(\eta_x) - 1 = \xi_x - \log(\xi_x)$, $\eta_x = -b^{-1} \log(\xi_x)$, the problem is reduced to study the fluctuation field of the conserved quantities for the process $(\xi(t))_{t \geq 0}$ at equilibrium under the probability measure $\nu_{\beta, \lambda}$. The fluctuation field for $(\xi(t))_{t \geq 0}$ is defined by

$$\mathcal{Z}_t^{n, \alpha}(\mathbf{G}) = \frac{1}{\sqrt{n}} \sum_{x \in \mathbb{Z}} \mathbf{G}(x/n) \cdot (\omega_x(tn^{1+\alpha}) - \omega), \quad (28)$$

where \mathbf{G} is a test function belonging to $S(\mathbb{R}) \times S(\mathbb{R})$. Recalling (18) we have

$$\mathcal{Y}_t^{n, \alpha}(\mathbf{G}) = \frac{1}{\sqrt{n}} \sum_{x \in \mathbb{Z}} (\Lambda^T \mathbf{G})(x/n) \cdot (\omega_x(tn^{1+\alpha}) - \omega) = \mathcal{Z}_t^{n, \alpha}(\Lambda^T \mathbf{G}). \quad (29)$$

By the relation $\bar{M} = \Lambda M \Lambda^{-1}$, we are able to translate any result about the convergence of $\mathcal{Z}^{n, \alpha}$ into a corresponding result for $\mathcal{Y}^{n, \alpha}$.

6.1. The hyperbolic scaling. For any local function $g := g(\xi)$ we define the projection $\mathcal{P}_{\rho, \theta} g$ of g on the fields of the conserved quantities by

$$(\mathcal{P}_{\rho, \theta} g)(\xi) = (\nabla \tilde{g})(\rho, \theta) \cdot (\omega_0 - \omega)$$

where $\tilde{g}(\rho, \theta) = \langle g \rangle_{\nu_{\beta, \lambda}}$ and $\nabla \tilde{g}$ is the gradient of the function \tilde{g} .

We have that

Proposition 1 (Boltzmann-Gibbs principle I) *For every $\mathbf{H} \in S(\mathbb{R}) \times S(\mathbb{R})$ and every $t > 0$,*

$$\lim_{n \rightarrow \infty} \mathbb{E}_{\nu_{\beta, \lambda}} \left[\left(\int_0^t \frac{1}{\sqrt{n}} \sum_{x \in \mathbb{Z}} \mathbf{H}(x/n) \cdot [\tau_x V_{J_{0,1}}(\xi(sn))] ds \right)^2 \right] = 0,$$

where for a local function g we define $V_g(\xi) := g(\xi) - \tilde{g}(\rho, \theta) - \mathcal{P}_{\rho, \theta} g(\xi)$ and for $\xi \in (0, +\infty)^{\mathbb{Z}}$, $\tau_x g(\xi) := g(\tau_x \xi)$, $\tau_x \xi(y) := \xi(x + y)$ and $J_{0,1}$ is given in (14).

Since we prove a refined version of this proposition we omit its proof. As a consequence of last result, we get that the fluctuation field $(\mathcal{Z}^{n,0})_{n \geq 1}$ converges in law (in the sense of Theorem 1) to \mathcal{Z}^0 solution of the equation at the right hand side of (21). Theorem 1 is a simple consequence of this fact.

In order to prove Corollary 1 and Corollary 2 we follow the approach first presented in [25] and considered also in [20] (resp. [18]) for the symmetric (resp. asymmetric) simple exclusion. For that reason we sketch the main steps of the proof. For more details we refer the reader to, for example, the proof of Theorem 4.2 of [18]. The main goal is to related the energy and volume currents with the density field and to use the result of Theorem 1. For that purpose and whenever the total energy (resp. volume) at η is finite we can write down the energy (resp. volume) flux through the bond $\{x, x+1\}$ during the time interval $[0, tn]$, as:

$$\mathcal{E}_{x,x+1}^n(t) := \sum_{y \geq x+1} \left\{ V_b(\eta_y(tn)) - V_b(\eta_y(0)) \right\}$$

$$\left(\text{resp. } \mathcal{V}_{x,x+1}^n(t) := \sum_{y \geq x+1} \left\{ \eta_y(tn) - \eta_y(0) \right\} \right).$$

In such case, we can relate the energy (resp. volume) flux given above with the energy-volume fluctuation field as

$$\mathcal{E}_{x,x+1}^n(t) := \mathcal{Y}_t^n(H_x^1) - \mathcal{Y}_0^n(H_x^1)$$

$$\left(\text{resp. } \mathcal{V}_{x,x+1}^n(t) := \mathcal{Y}_t^n(H_x^2) - \mathcal{Y}_0^n(H_x^2), \right.$$

where

$$H_x^1(y) = \begin{pmatrix} \mathbf{1}_{\{y \geq x\}} \\ 0 \end{pmatrix}, \quad H_x^2(y) = \begin{pmatrix} 0 \\ \mathbf{1}_{\{y \geq x\}} \end{pmatrix}.$$

Since the function $\mathbf{1}_{\{y \geq x\}}$ does not belong to our space of test functions for which we derived Theorem 1 we first show that

Proposition 1. *For every $t \geq 0$,*

$$\lim_{\ell \rightarrow \infty} \mathbb{E}_{\nu_{\beta, \lambda}} \left[\left(\mathcal{E}_{x,x+1}^n(t) - (\mathcal{Y}_t^n(G_{\ell,x}^1) - \mathcal{Y}_0^n(G_{\ell,x}^1)) \right)^2 \right] = 0,$$

$$\left(\text{resp. } \lim_{\ell \rightarrow \infty} \mathbb{E}_{\nu_{\beta, \lambda}} \left[\left(\mathcal{V}_{x,x+1}^n(t) - (\mathcal{Y}_t^n(G_{\ell,x}^2) - \mathcal{Y}_0^n(G_{\ell,x}^2)) \right)^2 \right] = 0, \right.$$

where

$$G_{\ell,x}^1(y) = \begin{pmatrix} G_{\ell,x}(y) \\ 0 \end{pmatrix}, \quad G_{\ell,x}^2(y) = \begin{pmatrix} 0 \\ G_{\ell,x}(y) \end{pmatrix}$$

and $G_{\ell,x}(y) := (1 - y/\ell) \mathbf{1}_{\{x \leq y \leq x+\ell\}}$.

The proof of last result follows the same lines as in the proof of Proposition 4.1 of [18] and for that reason we omitted it. We notice that, at this point we are still not able to apply Theorem 1 since $G_{\ell,x}^1$ and $G_{\ell,x}^2$ are not functions in $\mathcal{S}(\mathbb{R})$. Therefore, we approximate in $L^2(\mathbb{R})$ each one of these functions by smooth functions for which Theorem 1 holds. Then, the proof of Corollary 1 and

2 follows combining the previous proposition with Theorem 1. For more details on this argument, we refer the reader to [18].

Finally, in order to compute the limiting variance, for example for the energy flux, we do the following. Here we take $x = 0$ to simplify the notation

$$\begin{aligned}\mathbb{E}_Q[Z_t^e Z_s^e] &= \mathbb{E}_Q[\{\mathcal{Y}_t(H_0^1) - \mathcal{Y}_0(H_0^1)\}\{\mathcal{Y}_s(H_0^1) - \mathcal{Y}_0(H_0^1)\}] \\ &= \lim_{\ell \rightarrow \infty} \mathbb{E}_Q \left[\mathcal{Y}_t(G_{\ell,0}^1) \mathcal{Y}_s(G_{\ell,0}^1) - \mathcal{Y}_t(G_{\ell,0}^1) \mathcal{Y}_0(G_{\ell,0}^1) \right. \\ &\quad \left. - \mathcal{Y}_s(G_{\ell,0}^1) \mathcal{Y}_0(G_{\ell,0}^1) + \mathcal{Y}_0(G_{\ell,0}^1) \mathcal{Y}_0(G_{\ell,0}^1) \right]\end{aligned}$$

Now, to compute last expectation we use the change of variables. Notice that for $H, G \in S(\mathbb{R})$ we have that $\mathbb{E}_Q[\mathcal{Z}_t^0(H) \mathcal{Z}_0^0(G)] := \langle T_t^- H \cdot \chi G \rangle$. Combining this with (29), it follows that $\mathbb{E}_Q[\mathcal{Y}_t(H) \mathcal{Y}_0(G)] := \langle T_t^- (\Lambda^T H) \cdot \chi \Lambda^T G \rangle$. By the definition of $(T_t^-)_{t \geq 0}$ we have for G_1, G_2 test functions in $S(\mathbb{R})$:

$$T_t^- \begin{pmatrix} G_1(x) \\ G_2(x) \end{pmatrix} = \begin{pmatrix} \frac{1}{\rho} (G_2(x - 2b^2 \rho t) - G_2(x)) + G_1(x - 2b^2 \rho t) \\ G_2(x) \end{pmatrix}.$$

As a consequence we obtain that

$$\begin{aligned}\mathbb{E}_Q[Z_t^e Z_s^e] &= \left(1 - \frac{1}{\rho}\right)^2 \left(\frac{\lambda + 1}{\beta^2}\right) \\ &\quad \times \lim_{\ell \rightarrow \infty} \int_{\mathbb{R}} \left(G_\ell^t(x) G_\ell^s(x) - G_\ell^t(x) G_\ell(x) - G_\ell^s(x) G_\ell(x) + G_\ell(x) G_\ell(x) \right) dx,\end{aligned}$$

where for $t \geq 0$, $G_\ell^t(x) := G_{\ell,0}(x - 2b^2 \rho t)$. Now, using (16) and (17) the proof ends. Analogously, repeating the computations above, replacing H_0^1 by H_0^2 we get the covariance for the volume flux.

6.2. The longer time scale. Since in the hyperbolic time scale the initial fluctuations for the field $\mathcal{Z}^{n,\alpha}$ are transported by the transposed linearized system given on the right hand side of (21), we redefine the fluctuation field $\widehat{\mathcal{Z}}^{n,\alpha}$, $\alpha > 0$, on $\mathbf{G} \in S(\mathbb{R}) \times S(\mathbb{R})$, by

$$\widehat{\mathcal{Z}}_t^{n,\alpha}(\mathbf{G}) = \mathcal{Z}_t^{n,\alpha}(T_{tn^\alpha}^+ \mathbf{G}).$$

By Dynkin's formula, see for example Appendix 1, section 5 of [22]

$$\mathcal{M}_t^{n,\alpha}(\mathbf{G}) = \widehat{\mathcal{Z}}_t^{n,\alpha}(\mathbf{G}) - \widehat{\mathcal{Z}}_0^{n,\alpha}(\mathbf{G}) - \int_0^t \left\{ n^{1+\alpha} \mathcal{L} \left(\widehat{\mathcal{Z}}_s^{n,\alpha}(\mathbf{G}) \right) + \partial_s \widehat{\mathcal{Z}}_s^{n,\alpha}(\mathbf{G}) \right\} ds$$

is a martingale with quadratic variation given by

$$\langle \mathcal{M}^{n,\alpha} \rangle_t = \int_0^t n^{1+\alpha} \mathcal{L} \left(\widehat{\mathcal{Z}}_s^{n,\alpha}(\mathbf{G}) \right)^2 - 2n^{1+\alpha} \left(\widehat{\mathcal{Z}}_s^{n,\alpha}(\mathbf{G}) \right) \mathcal{L} \left(\widehat{\mathcal{Z}}_s^{n,\alpha}(\mathbf{G}) \right) ds.$$

A simple computation shows that $\mathbb{E}_{\nu_{\beta,\lambda}}[\langle \mathcal{M}^{n,\alpha} \rangle_t]$ vanishes as n goes to ∞ for $\alpha < 1$. This is equivalent to saying that the martingale $\mathcal{M}_t^{n,\alpha}$ vanishes as n goes to ∞ in $\mathbb{L}^2(\mathbb{P}_{\nu_{\beta,\lambda}})$, for $\alpha < 1$. Observe that, by definition of $(T_t^+)_{t \geq 0}$, we have

$$\begin{aligned} \partial_s \widehat{\mathcal{Z}}_s^{n,\alpha}(\mathbf{G}) &= -\frac{n^\alpha}{\sqrt{n}} \sum_{x \in \mathbb{Z}} M^T [\partial_q (T_{sn^\alpha}^+ \mathbf{G})(x/n)] \cdot (\omega_x(sn^{1+\alpha}) - \omega) \\ &= -\frac{n^\alpha}{\sqrt{n}} \sum_{x \in \mathbb{Z}} [\partial_q (T_{sn^\alpha}^+ \mathbf{G})(x/n)] \cdot M(\omega_x(sn^{1+\alpha}) - \omega). \end{aligned}$$

On the other hand, the first term in the integral part of the martingale $\mathcal{M}_t^{n,\alpha}(\mathbf{G})$ is equal to

$$\frac{n^\alpha}{\sqrt{n}} \sum_{x \in \mathbb{Z}} n \left((T_{sn^\alpha}^+ \mathbf{G}) \left(\frac{x+1}{n} \right) - (T_{sn^\alpha}^+ \mathbf{G}) \left(\frac{x}{n} \right) \right) \cdot (J_{x,x+1}(\xi(sn^{1+\alpha})) - \langle J_{x,x+1} \rangle_{\nu_{\beta,\lambda}}).$$

Performing a Taylor expansion, we can replace this term, up to a term vanishing as n goes to ∞ in $\mathbb{L}^2(\mathbb{P}_{\nu_{\beta,\lambda}})$, by

$$\frac{n^\alpha}{\sqrt{n}} \sum_{x \in \mathbb{Z}} \left((\partial_q T_{sn^\alpha}^+ \mathbf{G})(x/n) \right) \cdot (J_{x,x+1}(\xi(sn^{1+\alpha})) - \langle J_{x,x+1} \rangle_{\nu_{\beta,\lambda}}).$$

Thus, in order to show that

$$\lim_{n \rightarrow \infty} \mathbb{E}_{\nu_{\beta,\lambda}} \left[\left(\widehat{\mathcal{Z}}_t^{n,\alpha}(\mathbf{G}) - \widehat{\mathcal{Z}}_0^{n,\alpha}(\mathbf{G}) \right)^2 \right] = 0, \quad (30)$$

it remains to show that

$$\lim_{n \rightarrow \infty} \mathbb{E}_{\nu_{\beta,\lambda}} \left[\left(\frac{n^\alpha}{\sqrt{n}} \int_0^t ds \sum_{x \in \mathbb{Z}} (\partial_q T_{sn^\alpha}^+ \mathbf{G})(x/n) \cdot \Theta_x(\xi(sn^{1+\alpha})) \right)^2 \right] = 0$$

where for $\xi \in (0, +\infty)^{\mathbb{Z}}$

$$\Theta_x(\xi) = J_{x,x+1}(\xi) - \langle J_{x,x+1} \rangle_{\nu_{\beta,\lambda}} - M(\omega_x - \omega).$$

Observe that in this formula, $M := M(\rho, \theta)$ is the differential with respect to (ρ, θ) of the function $\langle J_{x,x+1} \rangle_{\nu_{\beta,\lambda}}$ as computed below (21). A simple computation shows that for $\xi \in (0, +\infty)^{\mathbb{Z}}$

$$\Theta_x(\xi) = \begin{pmatrix} -b^2(\xi_{x+1} - \rho)(\xi_x - \rho) - (\gamma + b^2\rho)\nabla \xi_x \\ -\nabla(b^2\xi_x + \gamma \log(\xi_x)) \end{pmatrix}.$$

The discrete gradient terms appearing in the previous expression, permit to perform another discrete integration by parts and the resulting terms vanish in $\mathbb{L}^2(\mathbb{P}_{\nu_{\beta,\lambda}})$ as n goes to ∞ , for $\alpha < 1$. Using the smoothness of the function \mathbf{G} , we see that it only remains to show the following theorem with $\varphi(s, q)$ equal to the first component of the column vector $\partial_q T_s^+ \mathbf{G}$.

Theorem 4 (Boltzmann-Gibbs principle II) Fix $\alpha < 1/3$ and let $\varphi : \mathbb{R}^+ \times \mathbb{R} \rightarrow \mathbb{R}$ be such that for any $t \geq 0$, $\varphi(t, \cdot) \in S(\mathbb{R})$. For every $t > 0$

$$\lim_{n \rightarrow \infty} \mathbb{E}_{\nu_{\beta, \lambda}} \left[\left(\int_0^t \frac{n^\alpha}{\sqrt{n}} \sum_{x \in \mathbb{Z}} \varphi(sn^\alpha, x/n) (\xi_x(sn^{1+\alpha}) - \rho) (\xi_{x+1}(sn^{1+\alpha}) - \rho) ds \right)^2 \right] = 0$$

Proof. In the following, C, C_0, C_1, \dots denote constants independent of n whose values can change from line to line.

Let $f_s(\xi)$ be the function defined by

$$f_s(\xi) = \sum_{x \in \mathbb{Z}} \varphi(s, x/n) H_{\delta_x + \delta_{x+1}}(\xi) = -\beta^2 \sum_{x \in \mathbb{Z}} \varphi(s, x/n) (\xi_x - \rho) (\xi_{x+1} - \rho).$$

The last equality follows from (24) and (17).

We have the following upper bound

$$\begin{aligned} \mathbb{E}_{\nu_{\beta, \lambda}} \left[\left(\int_0^t f_{sn^\alpha}(\xi(sn^{1+\alpha})) ds \right)^2 \right] &\leq C \int_0^t \langle f_{sn^\alpha}, (s^{-1} - n^{1+\alpha} \mathcal{L})^{-1} f_{sn^\alpha} \rangle_{\nu_{\beta, \lambda}} ds \\ &= \frac{C}{n^{1+\alpha}} \int_0^t \left\langle f_{sn^\alpha}, \left(\frac{1}{sn^{1+\alpha}} - \mathcal{L} \right)^{-1} f_{sn^\alpha} \right\rangle_{\nu_{\beta, \lambda}} ds \\ &\leq \frac{C}{n^{1+\alpha}} \int_0^t \left\langle f_{sn^\alpha}, \left(\frac{1}{sn^{1+\alpha}} - \gamma \mathcal{S} \right)^{-1} f_{sn^\alpha} \right\rangle_{\nu_{\beta, \lambda}} ds. \end{aligned}$$

In the first inequality above we used Lemma 3.9 of [27] applied to this setting. We notice that since our test functions depend on time, the lemma of [27] has to be modified as written here. To prove the last result one can simply adapt the proof of Lemma 4.3 of [10] to this case.

In order to simplify notations, let us define $\varepsilon = 1/sn^{1+\alpha}$.

We denote by Σ_2^0 the set of configurations σ of Σ_2 such that $\sigma = 2\delta_x$, $x \in \mathbb{Z}$, and Σ_2^\pm the complementary set of Σ_2^0 in Σ_2 , i.e. the set of configurations $\sigma \in \Sigma_2$ such that $\sigma = \delta_x + \delta_y$, $y \neq x \in \mathbb{Z}$. Observe that f_{sn^α} is a function of degree 2 with a decomposition in the form $f_{sn^\alpha} = \sum_{\sigma \in \Sigma_2} \Phi_{sn^\alpha}(\sigma) H_\sigma$ which satisfies $\Phi_{sn^\alpha}(\sigma) = 0$ if $\sigma \in \Sigma_2^0$. We have that (see e.g. [27])

$$\left\langle f_{sn^\alpha}, (\varepsilon - \gamma \mathcal{S})^{-1} f_{sn^\alpha} \right\rangle_{\nu_{\beta, \lambda}} = \sup_g \left\{ 2 \langle f_{sn^\alpha}, g \rangle_{\nu_{\beta, \lambda}} - \varepsilon \langle g, g \rangle_{\nu_{\beta, \lambda}} - \gamma \mathcal{D}(g) \right\}$$

where the supremum is taken over local functions $g \in \mathbb{L}^2(\nu_{\beta, \lambda})$. Decompose g appearing in this variational formula as $g = \sum_{\sigma} G(\sigma) H_\sigma$. Recall that $\{H_\sigma; \sigma \in \Sigma\}$ are orthogonal, that the function f_{sn^α} is a degree 2 function such that $\Phi_{sn^\alpha}(\sigma) = 0$ for any $\sigma \notin \Sigma_2^\pm$ and formula (27) for the Dirichlet form $\mathcal{D}(g)$. Thus, we can restrict this supremum over degree 2 functions g such that $G(\sigma) = 0$ if $\sigma \in \Sigma_2^0$.

Then, by Lemma 1, we have

$$\left\langle f_{sn^\alpha}, (\varepsilon - \gamma \mathcal{S})^{-1} f_{sn^\alpha} \right\rangle_{\nu_{\beta, \lambda}} \leq \sup_G \left\{ \sum_{x \neq y} \Phi_{sn^\alpha}(x, y) G(x, y) - \varepsilon \sum_{\substack{(x, y) \in \mathbb{Z}^2 \\ x \neq y}} G^2(x, y) \right. \\ \left. - C \sum_{|\mathbf{e}|=1} \sum_{\substack{(x, y) \in \Delta^\pm \\ (x, y) + \mathbf{e} \in \Delta^\pm}} \left(G((x, y) + \mathbf{e}) - G(x, y) \right)^2 \right\}$$

where $C := C(\lambda, \gamma)$, $\Delta_\pm = \{(x, y) \in \mathbb{Z}^2; x \neq y\}$ and as usual we identify the functions defined on Σ_n with symmetric functions defined on \mathbb{Z}^n .

In order to get rid of the geometric constraints appearing in the last term of the variational formula, for any symmetric function G defined on the set Δ_\pm , we denote by \tilde{G} its extension to \mathbb{Z}^2 defined by

$$\tilde{G}(x, y) = G(x, y) \text{ if } x \neq y, \quad \tilde{G}(x, x) = \frac{1}{4} \sum_{|\mathbf{e}|=1} G((x, x) + \mathbf{e}).$$

It is trivial that

$$\sum_{(x, y) \in \mathbb{Z}^2} \tilde{G}^2(x, y) \leq C \sum_{\substack{(x, y) \in \mathbb{Z}^2 \\ x \neq y}} G^2(x, y),$$

and

$$\sum_{|\mathbf{e}|=1} \sum_{(x, y) \in \mathbb{Z}^2} \left(\tilde{G}((x, y) + \mathbf{e}) - \tilde{G}(x, y) \right)^2 \\ \leq C \sum_{|\mathbf{e}|=1} \sum_{\substack{(x, y) \in \Delta^\pm \\ (x, y) + \mathbf{e} \in \Delta^\pm}} \left(G((x, y) + \mathbf{e}) - G(x, y) \right)^2.$$

Thus, we have

$$\left\langle f_{sn^\alpha}, (\varepsilon - \gamma \mathcal{S})^{-1} f_{sn^\alpha} \right\rangle_{\nu_{\beta, \lambda}} \\ \leq C_0 \sup_G \left\{ \sum_{(x, y) \in \mathbb{Z}^2} \Phi_{sn^\alpha}(x, y) G(x, y) - C_1 \varepsilon \sum_{(x, y) \in \mathbb{Z}^2} G^2(x, y) \right. \\ \left. - C_2 \sum_{|\mathbf{e}|=1} \sum_{(x, y) \in \mathbb{Z}^2} \left(G((x, y) + \mathbf{e}) - G(x, y) \right)^2 \right\}$$

where the supremum is now taken over all symmetric local functions $G : \mathbb{Z}^2 \rightarrow \mathbb{R}$. Notice that the last variational formula is equal to the resolvent norm, for a

simple symmetric two dimensional random walk, of the function Φ_{sn^α} . By using Fourier transform one can easily show that this supremum is equal to

$$\frac{C_0}{4} \int_{[0,1]^2} \frac{|\hat{\Phi}_{sn^\alpha}(\mathbf{k})|^2}{C_1\varepsilon + 4C_2 \sum_{i=1}^2 \sin^2(\pi k_i)} d\mathbf{k}$$

where the Fourier transform $\hat{\Phi}_{sn^\alpha}$ of Φ_{sn^α} is given by

$$\hat{\Phi}_{sn^\alpha}(\mathbf{k}) = \sum_{(x,y) \in \mathbb{Z}^2} \Phi_{sn^\alpha}(x,y) e^{2i\pi(k_1x+k_2y)}, \quad \mathbf{k} = (k_1, k_2) \in [0,1]^2.$$

By definition of f_{sn^α} , we have $\Phi_{sn^\alpha}(x,y) = \frac{1}{2}(\varphi(sn^\alpha, x/n) + \varphi(sn^\alpha, y/n))$ if $|x-y|=1$ and 0 otherwise. Consequently, we have

$$\begin{aligned} \left\langle f_{sn^\alpha}, (\varepsilon - \gamma\mathcal{S})^{-1} f_{sn^\alpha} \right\rangle_{\nu_{\beta,\lambda}} &\leq \frac{C_0}{16} \int_{[0,1]^2} \frac{|\sum_{x \in \mathbb{Z}} \varphi(sn^\alpha, x/n) e^{2i\pi x(k_1+k_2)}|^2}{C_1\varepsilon + 4C_2 \sum_{i=1}^2 \sin^2(\pi k_i)} d\mathbf{k} \\ &= \frac{C_0}{16} \int_{[0,2]} \left(\int_{[0,1]} \frac{\mathbf{1}_{[\sup(1-p,1), \inf(1,0)]}(p) |\sum_{x \in \mathbb{Z}} \varphi(sn^\alpha, x/n) e^{2i\pi xp}|^2}{C_1\varepsilon + 4C_2 \sin^2(\pi k_1) + 4C_2 \sin^2(\pi(p-k_1))} dk_1 \right) dp \\ &= \frac{C_0}{16} \int_{[0,1]} \left(\int_{[0,1]} \frac{|\sum_{x \in \mathbb{Z}} \varphi(sn^\alpha, x/n) e^{2i\pi xp}|^2}{C_1\varepsilon + 4C_2 \sin^2(\pi k_1) + 4C_2 \sin^2(\pi(p-k_1))} dk_1 \right) dp \end{aligned}$$

where we used the change of variables $p = k_2 + k_1$ for the first equality and the periodicity of the functions involved for the second one. It follows that

$$\begin{aligned} \left\langle f_{sn^\alpha}, (\varepsilon - \gamma\mathcal{S})^{-1} f_{sn^\alpha} \right\rangle_{\nu_{\beta,\lambda}} &\leq \frac{C_0}{16} \int_{[0,1]} \left| \sum_{x \in \mathbb{Z}} \varphi(sn^\alpha, x/n) e^{2i\pi xp} \right|^2 dp \int_{[0,1]} \frac{dk_1}{C_1\varepsilon + 4C_2 \sin^2(\pi k_1)} \\ &\leq \frac{C}{\sqrt{\varepsilon}} \int_{[0,1]} \left| \sum_{x \in \mathbb{Z}} \varphi(sn^\alpha, x/n) e^{2i\pi xp} \right|^2 dp. \end{aligned}$$

Observe now that

$$\int_{[0,1]} \left| \sum_{x \in \mathbb{Z}} \varphi(sn^\alpha, x/n) e^{2i\pi xp} \right|^2 dp = \sum_{x \in \mathbb{Z}} \varphi^2(sn^\alpha, x/n) \leq Cn.$$

Putting everything together, we get that

$$\begin{aligned} \mathbb{E}_{\nu_{\beta,\lambda}} \left[\left(\int_0^t \frac{n^\alpha}{\sqrt{n}} \sum_{x \in \mathbb{Z}} \varphi(sn^\alpha, x/n) (\xi_x(sn^{1+\alpha}) - \rho) (\xi_{x+1}(sn^{1+\alpha}) - \rho) ds \right)^2 \right] \\ \leq \frac{Ctn^{2\alpha-1}}{n^{1+\alpha}} \int_0^t \frac{n}{\sqrt{\varepsilon}} ds. \end{aligned}$$

Since $\varepsilon := 1/sn^{1+\alpha}$ last expression vanishes as n goes to ∞ , if $\alpha < 1/3$.

Now, in order to prove Corollary 3 we follow the same arguments as in the proof of Proposition 9.3 of [18] and we proceed as follows. Whenever the total energy (resp. volume) at η is finite we can write down:

$$\begin{aligned} \mathcal{E}_{u_t^{x,\alpha}(n)}^n(t) &:= \sum_{y \geq u_t^{x,\alpha}(n)} \left\{ V_b(\eta_y(tn^{1+\alpha})) - V_b(\eta_y(0)) \right\}, \\ \left(\text{resp. } \mathcal{V}_{u_t^{x,\alpha}(n)}^n(t) &:= \sum_{y \geq u_t^{x,\alpha}(n)} \left\{ \eta_y(tn^{1+\alpha}) - \eta_y(0) \right\} \right). \end{aligned} \quad (31)$$

In order to justify the previous equalities one can repeat the same arguments as used in the hyperbolic scaling. Now, we use the change of variables to define the energy (resp. volume) flux through the time-dependent bond $\{u_t^{x,\alpha}(n), u_t^{x,\alpha}(n) + 1\}$ during the time interval $[0, tn^{1+\alpha}]$. For that purpose, we define the flux fields in terms of ξ_x such that

$$\begin{aligned} \tilde{\mathcal{E}}_{x-1,x}^n(t) - \tilde{\mathcal{E}}_{x,x+1}^n(t) &:= \xi_x(tn^{1+\alpha}) - \xi_x(0) \\ \left(\text{resp. } \tilde{\mathcal{V}}_{x-1,x}^n(t) - \tilde{\mathcal{V}}_{x,x+1}^n(t) &:= \log(\xi_x(tn^{1+\alpha})) - \log(\xi_x(0)). \right. \end{aligned}$$

As above, when it makes sense, we have that

$$\begin{aligned} \tilde{\mathcal{E}}_{u_t^{x,\alpha}(n)}^n(t) &:= \sum_{y \geq u_t^{x,\alpha}(n)} \left\{ \xi_y(tn^{1+\alpha}) - \xi_y(0) \right\} \\ \left(\text{resp. } \tilde{\mathcal{V}}_{u_t^{x,\alpha}(n)}^n(t) &:= \sum_{y \geq u_t^{x,\alpha}(n)} \left\{ \log(\xi_y(tn^{1+\alpha})) - \log(\xi_y(0)) \right\} \right) \end{aligned}$$

and in this case we can write the previous fields in terms of $\widehat{Z}_t^{n,\alpha}$. A simple computation shows that Proposition 1 can similarly be stated for last fields. Combining this with (15) we have that

$$\mathcal{E}_{u_t^{x,\alpha}(n)}^n(t) := \tilde{\mathcal{E}}_{u_t^{x,\alpha}(n)}^n(t) - \tilde{\mathcal{V}}_{u_t^{x,\alpha}(n)}^n(t), \quad \mathcal{V}_{u_t^{x,\alpha}(n)}^n(t) := -\frac{1}{b} \tilde{\mathcal{V}}_{u_t^{x,\alpha}(n)}^n(t).$$

Then, applying (30) to $G_{\ell,x}^1(y) = \begin{pmatrix} G_{\ell,x}(y) \\ 0 \end{pmatrix}$ we obtain that

$$\lim_{n \rightarrow \infty} \mathbb{E}_{\nu_{\beta,\lambda}} \left[\left(\frac{1}{\sqrt{n}} \left\{ \tilde{\mathcal{E}}_{u_t^{x,\alpha}(n)}^n(t) - \mathbb{E}_{\nu_{\beta,\lambda}}[\tilde{\mathcal{E}}_{u_t^{x,\alpha}(n)}^n(t)] \right\} \right)^2 \right] = 0.$$

On the other hand, applying (30) to

$$\tilde{G}_{\ell,x}(y) = \begin{pmatrix} \frac{1}{\rho} G_{\ell,x}(y - u_t^{x,\alpha}(n)) \\ -G_{\ell,x}(y - u_t^{x,\alpha}(n)) \end{pmatrix}$$

we obtain that

$$\begin{aligned} \lim_{n \rightarrow \infty} \mathbb{E}_{\nu_{\beta,\lambda}} \left[\left(\frac{1}{\rho} \frac{1}{\sqrt{n}} \left\{ \tilde{\mathcal{E}}_{u_t^{x,\alpha}(n)}^n(t) - \mathbb{E}_{\nu_{\beta,\lambda}}[\tilde{\mathcal{E}}_{u_t^{x,\alpha}(n)}^n(t)] \right\} \right. \right. \\ \left. \left. - \frac{1}{\sqrt{n}} \left\{ \tilde{\mathcal{V}}_{u_t^{x,\alpha}(n)}^n(t) - \mathbb{E}_{\nu_{\beta,\lambda}}[\tilde{\mathcal{V}}_{u_t^{x,\alpha}(n)}^n(t)] \right\} \right)^2 \right] = 0. \end{aligned}$$

Now, Corollary 3 follows easily from the previous results.

Remark 1. From (20), the hydrodynamic equation of ρ is independent of θ and it can be rewritten as $\partial_t \rho - 2b^2 \rho \partial_q \rho = 0$. Following the system along the characteristics for ρ , that is, removing the velocity $2b^2 \rho$ from the system, we do not see a time evolution for ρ , and since $1/\rho \partial_t \rho - \partial_t \theta = 0$, nor for θ . Therefore, translating the velocity $2b^2 \rho$ in terms of the original variables it corresponds to $2b\bar{\lambda}/\beta$ and that is the reason why we took the time dependent bond as written in Corollary 3.

7. Diffusivity

In this section we prove Theorem 3. Our proof is based on the resolvent methods introduced in [4, 23] and developed in few other contexts (e.g. [5, 26, 30]). Some differences with these previous works are the presence of two and not only one conserved quantity and the degeneracy of the symmetric part of the generator.

The main steps of the proof are the following. First we use the microscopic change of variables and express the Laplace transform of the current-current correlation function as a resolvent norm in a suitable Hilbert space (see (33)). Then, we rewrite this resolvent norm as the supremum over the set of local functions of a functional acting on these functions (see (36)). To get a lower bound we restrict the supremum over degree two functions. The estimate of the value of the functional for a given degree two function remains in general very difficult. Thus we replace the functional restricted to the set of degree two functions by an equivalent functional simpler to estimate. This is accomplished through Lemma 3, Lemma 4 and Lemma 5. In the context of the asymmetric simple exclusion, this replacement step is called the “free particles approximation” ([4]) or the “hard core removal” ([23]). It is then possible to estimate the value of this equivalent functional for a suitable degree two test function.

We fix $\rho \geq 0, \theta \in \mathbb{R}$ and denote by β, λ the chemical potentials given by (17). Let also $(\bar{\beta}, \bar{\lambda})$ be given in terms of (β, λ) by (16).

Recall the definition of $\hat{J}_{x,x+1}$ given in (11). We introduce the normalized currents $\mathbf{j}_{x,x+1}, \mathbf{j}'_{x,x+1}$ and $\mathbf{J}_{x,x+1}$ corresponding to the process $(\xi(t))_{t \geq 0}$, which are defined by

$$\begin{aligned} \mathbf{j}_{x,x+1}(\xi) &= j_{x,x+1}(\xi) - \langle j_{x,x+1} \rangle_{\nu_{\beta,\lambda}} - \partial_\rho \langle j_{x,x+1} \rangle_{\nu_{\beta,\lambda}} (\xi_x - \rho) - \partial_\theta \langle j_{x,x+1} \rangle_{\nu_{\beta,\lambda}} (\log(\xi_x) - \theta), \\ \mathbf{j}'_{x,x+1}(\xi) &= j'_{x,x+1}(\xi) - \langle j'_{x,x+1} \rangle_{\nu_{\beta,\lambda}} - \partial_\rho \langle j'_{x,x+1} \rangle_{\nu_{\beta,\lambda}} (\xi_x - \rho) - \partial_\theta \langle j'_{x,x+1} \rangle_{\nu_{\beta,\lambda}} (\log(\xi_x) - \theta), \\ \mathbf{J}_{x,x+1}(\xi) &= \mathbf{j}_{x,x+1}(\xi) - \mathbf{j}'_{x,x+1}(\xi). \end{aligned}$$

Since $\langle j_{x,x+1} \rangle_{\nu_{\beta,\lambda}} = -b^2 \rho^2$ and $\langle j'_{x,x+1} \rangle_{\nu_{\beta,\lambda}} = -2b^2 \rho$, we get

$$\begin{aligned} \mathbf{j}_{x,x+1}(\xi) &= -b^2 (\xi_x - \rho) (\xi_{x+1} - \rho) - (\gamma + b^2 \rho) \nabla \xi_x \\ \mathbf{j}'_{x,x+1}(\xi) &= -\nabla (b^2 \xi_x + \gamma \log(\xi_x)). \end{aligned} \tag{32}$$

For any local compactly supported functions $f, g : (0, +\infty)^{\mathbb{Z}} \rightarrow \mathbb{R}$ we define the semi-inner product $\ll f, g \gg := \ll f, g \gg_{\beta, \lambda}$ of f and g by

$$\begin{aligned}
& \ll f, g \gg \\
&= \sum_{x \in \mathbb{Z}} (\langle \tau_x f g \rangle_{\nu_{\beta, \lambda}} - \langle f \rangle_{\nu_{\beta, \lambda}} \langle g \rangle_{\nu_{\beta, \lambda}}) \\
&= \lim_{k \rightarrow \infty} \sum_{|x| \leq k} (\langle \tau_x f g \rangle_{\nu_{\beta, \lambda}} - \langle f \rangle_{\nu_{\beta, \lambda}} \langle g \rangle_{\nu_{\beta, \lambda}}) \\
&= \lim_{k \rightarrow \infty} \frac{1}{2k+1} \sum_{|x| \leq k} \left\{ \sum_{|y-x| \leq k} (\langle \tau_{x+y} f \tau_y g \rangle_{\nu_{\beta, \lambda}} - \langle f \rangle_{\nu_{\beta, \lambda}} \langle g \rangle_{\nu_{\beta, \lambda}}) \right\} \\
&= \lim_{k \rightarrow \infty} \left\langle \left(\frac{1}{\sqrt{2k+1}} \sum_{|x| \leq k} (\tau_x f - \langle f \rangle_{\nu_{\beta, \lambda}}) \right) \left(\frac{1}{\sqrt{2k+1}} \sum_{|x| \leq k} (\tau_x g - \langle g \rangle_{\nu_{\beta, \lambda}}) \right) \right\rangle_{\nu_{\beta, \lambda}}
\end{aligned}$$

where the third equality follows from the invariance of $\nu_{\beta, \lambda}$ by the shift. Observe also that the first sum on \mathbb{Z} is in fact a finite sum since f and g are assumed to be local functions. We denote by \mathcal{H}_0 the space generated by the local compactly supported functions and the semi-inner product $\ll \cdot, \cdot \gg$. Observe that any constant or gradient functions are equal to 0 in \mathcal{H}_0 .

By (32), the normalized current associated to the volume is a gradient and this shows that $\mathcal{F}_{i,j}(\gamma, z) = 0$ if $(i, j) \neq (1, 1)$. By the definition of $\hat{J}_{x, x+1}$ and by (15), we are only interested in the behavior, as $z \rightarrow 0$, of

$$\mathfrak{L}(z) = \ll \mathbf{J}_{0,1}, (z - \mathcal{L})^{-1} \mathbf{J}_{0,1} \gg = \int_0^\infty e^{-zt} \ll \mathbf{J}_{0,1}(t), \mathbf{J}_{0,1}(0) \gg dt.$$

Since gradient functions are equal to 0 in \mathcal{H}_0 , this is equivalent to estimate

$$\mathfrak{L}(z) = b^4 \ll W_{0,1}, (z - \mathcal{L})^{-1} W_{0,1} \gg \quad (33)$$

where $W_{x,y}$ is the local function $W_{x,y} = (\xi_x - \rho)(\xi_y - \rho)$.

In this section we prove that there exists a constant $C > 0$ such that

$$\ll W_{0,1}, (z - \mathcal{L})^{-1} W_{0,1} \gg \geq Cz^{-1/4}. \quad (34)$$

But before proving (34) let us show (13) which is a direct consequence of the following lemma.

Lemma 2. *For any $\gamma > 0$, there exists a constant $C := C(\gamma)$ such that*

$$\ll W_{0,1}, (z/\gamma - b^2 \mathcal{A} - \mathcal{S})^{-1} W_{0,1} \gg \leq C \ll W_{0,1}, (z - b^2 \mathcal{A} - \gamma \mathcal{S})^{-1} W_{0,1} \gg$$

and

$$\ll W_{0,1}, (z - b^2 \mathcal{A} - \gamma \mathcal{S})^{-1} W_{0,1} \gg \leq C \ll W_{0,1}, (z/\gamma - b^2 \mathcal{A} - \mathcal{S})^{-1} W_{0,1} \gg.$$

Proof. Assume $\gamma > 1$ the case $\gamma < 1$ being similar. By Lemma 2.1 of [4] we have the variational formula for $\ll W_{0,1}, (z - \mathcal{L})^{-1}W_{0,1} \gg$, where $\mathcal{L} = b^2\mathcal{A} + \gamma\mathcal{S}$, given by

$$\sup_f \left\{ 2 \ll W_{0,1}, f \gg - \ll f, (z - \gamma\mathcal{S})f \gg - b^4 \ll \mathcal{A}f, (z - \gamma\mathcal{S})^{-1}\mathcal{A}f \gg \right\},$$

where the supremum is carried over functions f belonging to the domain of the generator \mathcal{L} or equivalently to a dense subspace included in this domain, say the space of smooth local compactly supported functions. We have that

$$\begin{aligned} & \sup_f \left\{ 2 \ll W_{0,1}, f \gg - \ll f, (z - \gamma\mathcal{S})f \gg - b^4 \ll \mathcal{A}f, (z - \gamma\mathcal{S})^{-1}\mathcal{A}f \gg \right\} \\ &= \sup_f \left\{ 2 \ll W_{0,1}, f \gg - \gamma \ll f, (z/\gamma - \mathcal{S})f \gg - b^4\gamma^{-1} \ll \mathcal{A}f, (z/\gamma - \mathcal{S})^{-1}\mathcal{A}f \gg \right\} \\ &\geq \sup_f \left\{ 2 \ll W_{0,1}, f \gg - \gamma \ll f, (z/\gamma - \mathcal{S})f \gg - b^4\gamma \ll \mathcal{A}f, (z/\gamma - \mathcal{S})^{-1}\mathcal{A}f \gg \right\} \\ &= \sup_f \left\{ 2\gamma^{-1/2} \ll W_{0,1}, f \gg - \ll f, (z/\gamma - \mathcal{S})f \gg - b^4 \ll \mathcal{A}f, (z/\gamma - \mathcal{S})^{-1}\mathcal{A}f \gg \right\} \end{aligned}$$

where the inequality comes from $\gamma > 1$ and last equality is obtained by the change of f into $\gamma^{-1/2}f$. The last term is equal to

$$\gamma^{-1} \ll W_{0,1}, (z/\gamma - b^2\mathcal{A} - \mathcal{S})^{-1}W_{0,1} \gg$$

and this proves the first inequality of the lemma.

For the second one we proceed similarly:

$$\begin{aligned} & \sup_f \left\{ 2 \ll W_{0,1}, f \gg - \ll f, (z - \gamma\mathcal{S})f \gg - b^4 \ll \mathcal{A}f, (z - \gamma\mathcal{S})^{-1}\mathcal{A}f \gg \right\} \\ &= \sup_f \left\{ 2 \ll W_{0,1}, f \gg - \gamma \ll f, (z/\gamma - \mathcal{S})f \gg - b^4\gamma^{-1} \ll \mathcal{A}f, (z/\gamma - \mathcal{S})^{-1}\mathcal{A}f \gg \right\} \\ &\leq \sup_f \left\{ 2 \ll W_{0,1}, f \gg - \gamma^{-1} \ll f, (z/\gamma - \mathcal{S})f \gg - b^4\gamma^{-1} \ll \mathcal{A}f, (z/\gamma - \mathcal{S})^{-1}\mathcal{A}f \gg \right\} \\ &= \sup_f \left\{ 2\gamma^{1/2} \ll W_{0,1}, f \gg - \ll f, (z/\gamma - \mathcal{S})f \gg - b^4 \ll \mathcal{A}f, (z/\gamma - \mathcal{S})^{-1}\mathcal{A}f \gg \right\} \\ &= \gamma \ll W_{0,1}, (z/\gamma - b^2\mathcal{A} - \mathcal{S})^{-1}W_{0,1} \gg. \end{aligned}$$

Recall the orthogonal decomposition described in Section 5. Let $f = \sum_{\sigma} F(\sigma)H_{\sigma}$ and $g = \sum_{\sigma} G(\sigma)H_{\sigma}$ be two centered local functions. The configuration σ shifted by $z \in \mathbb{Z}$ is denoted by $\tau_z\sigma$, that is $\tau_z\sigma(x) = \sigma(x - z)$. We identify F_n, G_n , the restrictions of F, G to Σ_n , with symmetric functions on \mathbb{Z}^n . By (25) we have that

$$\ll f, g \gg = \sum_{z \in \mathbb{Z}} \sum_{\sigma \in \Sigma} F(\tau_z\sigma)G(\sigma)\mathcal{W}(\sigma),$$

where \mathcal{W} was defined in (26).

With some abuse of notations, we denote by $\ll F, G \gg$ the scalar product defined by

$$\ll F, G \gg = \sum_{z \in \mathbb{Z}} \sum_{\sigma \in \Sigma} F(\tau_z\sigma)G(\sigma)\mathcal{W}(\sigma).$$

We also introduce the inner product $\ll \cdot, \cdot \gg_{\text{free}}$ defined by

$$\ll F, G \gg_{\text{free}} = \sum_{y \in \mathbb{Z}} \sum_{\sigma \in \Sigma} F(\tau_y \sigma) G(\sigma).$$

Since the function \mathcal{W} is invariant by the shift, we have a very simple relation between these two inner products:

$$\ll F, G \gg = \ll \mathcal{W}^{1/2} F, \mathcal{W}^{1/2} G \gg_{\text{free}}. \quad (35)$$

On the set Σ_n we introduce the equivalence relation \star defined by $\sigma \star \sigma'$ if and only if there exists $u \in \mathbb{Z}$ such that $\tau_u \sigma = \sigma'$. Let $\Sigma_n^\star = \Sigma_n / \star$ be the set of classes for this relation and $\Sigma^\star = \cup_{n \geq 1} \Sigma_n^\star$. We can rewrite the scalar product $\ll \cdot, \cdot \gg_{\text{free}}$ as

$$\ll F, G \gg_{\text{free}} = \sum_{\bar{\sigma} \in \Sigma^\star} \bar{F}(\bar{\sigma}) \bar{G}(\bar{\sigma}).$$

Here \bar{F} is defined by $\bar{F}(\bar{\sigma}) = \left(\sum_{y \in \mathbb{Z}} \tau_y F \right) (\sigma)$ where σ is any element of $\bar{\sigma}$. The function \mathcal{W} being invariant by the shift, we define $\mathcal{W}(\bar{\sigma})$ by $\mathcal{W}(\sigma)$, $\sigma \in \bar{\sigma}$, $\bar{\sigma} \in \Sigma^\star$. Then, we have

$$\ll F, G \gg = \sum_{\bar{\sigma} \in \Sigma^\star} \mathcal{W}(\bar{\sigma}) \bar{F}(\bar{\sigma}) \bar{G}(\bar{\sigma}).$$

Lemma 3. *There exists a constant $C := C(n, \lambda)$ such that for any local function $F : \Sigma_n \rightarrow \mathbb{R}$ of degree n it holds that*

1.

$$C^{-1} \ll F, F \gg_{\text{free}} \leq \ll F, F \gg \leq C \ll F, F \gg_{\text{free}}.$$

2.

$$C^{-1} \ll F, -\mathfrak{S}F \gg_{\text{free}} \leq \ll F, -\mathfrak{S}F \gg \leq C \ll F, -\mathfrak{S}F \gg_{\text{free}}.$$

Moreover, for any positive real $z > 0$

$$\ll F, (z - \gamma \mathfrak{S})^{-1} F \gg = \ll \mathcal{W}^{1/2} F, (z - \gamma \mathfrak{S})^{-1} \mathcal{W}^{1/2} F \gg_{\text{free}}.$$

Proof. Recall the definition of \mathcal{W} from (26). Thus, \mathcal{W} is bounded from above (resp. from below) by a constant $C(n, \lambda)$ (resp. $C^{-1}(n, \lambda)$) independent of $\sigma \in \Sigma_n$. This is enough to conclude (1). In order to prove (2), it is enough to use (35) and the fact that for any local function $F : \Sigma \rightarrow \mathbb{R}$ we have that $\mathfrak{S}(\mathcal{W}^{1/2} F) = \mathcal{W}^{1/2} \mathfrak{S}F$. Finally, for a local function F of degree n , we have by (35) and the fact that

$$\ll F, (z - \gamma \mathfrak{S})^{-1} F \gg = \sup_{G \text{ of degree } n} \{2 \ll F, G \gg - \ll G, (z - \gamma \mathfrak{S})G \gg\},$$

the following equality

$$\ll F, (z - \gamma \mathfrak{S})^{-1} F \gg = \ll \mathcal{W}^{1/2} F, (z - \gamma \mathfrak{S})^{-1} \mathcal{W}^{1/2} F \gg_{\text{free}},$$

which proves the last assertion.

Our goal is to get a lower bound for $\ll W_{0,1}, (z - \mathcal{L})^{-1}W_{0,1} \gg$ which by Lemma 2.1 of [4] can be rewritten in the variational form

$$\sup_f \left\{ 2 \ll W_{0,1}, f \gg - \ll f, (z - \gamma \mathcal{S})f \gg - b^4 \ll \mathcal{A}f, (z - \gamma \mathcal{S})^{-1} \mathcal{A}f \gg \right\}. \quad (36)$$

Any element $\bar{\sigma} \in \Sigma_n^*$ can be identified with an element of \mathbb{N}^{n-1} through the application which associates to $(\alpha_1, \dots, \alpha_{n-1}) \in \mathbb{N}^{n-1}$ the class of the configuration $\sigma = \delta_0 + \delta_{\alpha_1} + \dots + \delta_{\alpha_1 + \dots + \alpha_{n-1}}$.

Observe also that \mathfrak{S} is a self-adjoint operator with respect to $\ll \cdot, \cdot \gg$ and with respect to $\ll \cdot, \cdot \gg_{\text{free}}$. We restrict the previous supremum over degree 2 functions $f = \sum_{(x,y) \in \mathbb{Z}^2} F([x,y])H_{[x,y]}$. In order to keep notation simple, whenever we identify a configuration $\sigma \in \Sigma_n$ with $[\mathbf{x}] \in \mathbb{Z}^n$ we will simply write $F(x)$, instead of $F([\mathbf{x}])$.

Up to some irrelevant multiplicative constant, a lower bound is given by

$$\sup_{F \text{ of degree 2}} \left\{ 2F(0,1) - \|F\|_{1,z}^2 - b^4 \|\mathfrak{A}_- F\|_{-1,z}^2 - b^4 \|\mathfrak{A}_+ F\|_{-1,z}^2 - b^4 \|\mathfrak{A}_0 F\|_{-1,z}^2 \right\}$$

where $\|F\|_{\pm 1,z}^2 = \ll F, (z - \gamma \mathfrak{S})^{\pm 1} F \gg$. We also introduce the corresponding $H_{\pm 1,z}$ -norms associated to $\ll \cdot, \cdot \gg_{\text{free}}$: $\|F\|_{\pm 1,z,\text{free}}^2 = \ll F, (z - \gamma \mathfrak{S})^{\pm 1} F \gg_{\text{free}}$, for $F : \Sigma \rightarrow \mathbb{R}$.

By Lemma 3, there exists a constant C such that this lower bound is bounded from below by

$$\begin{aligned} & \sup_{F \text{ of degree 2}} \left\{ 2F(0,1) - C \|F\|_{+1,z,\text{free}}^2 \right. \\ & \left. - b^4 \|\mathcal{W}^{1/2} \mathfrak{A}_- F\|_{-1,z,\text{free}}^2 - b^4 \|\mathcal{W}^{1/2} \mathfrak{A}_+ F\|_{-1,z,\text{free}}^2 - b^4 \|\mathcal{W}^{1/2} \mathfrak{A}_0 F\|_{-1,z,\text{free}}^2 \right\}. \end{aligned}$$

Let us first show that if F is of degree 2 then the contributions given by $\|\mathcal{W}^{1/2} \mathfrak{A}_- F\|_{-1,z,\text{free}}^2$ and $\|\mathcal{W}^{1/2} \mathfrak{A}_0 F\|_{-1,z,\text{free}}^2$ are equal to zero.

The function \mathcal{W} is constant and equal to $(\lambda + 1)$ on Σ_1 so that $\mathcal{W}^{1/2} \mathfrak{A}_- F = \sqrt{\lambda + 1} \mathfrak{A}_- F$. It is easy to check that the degree one function $\mathfrak{A}_- F$ satisfies

$$(\mathfrak{A}_- F)(u) = (\lambda - 1) \left(F(u-1, u) - F(u, u+1) \right).$$

For any degree 1 function G , we have

$$\ll \mathfrak{A}_- F, G \gg_{\text{free}} = \sum_{u,y \in \mathbb{Z}} G(u+y) (\lambda - 1) \left(F(u-1, u) - F(u, u+1) \right) = 0$$

by a telescopic sum argument. This shows that $\mathfrak{A}_- F$ is equal to zero in the Hilbert space generated by $\ll \cdot, \cdot \gg_{\text{free}}$.

Recall that if F is a degree 2 function, i.e. a symmetric function on \mathbb{Z}^2 , then F is identified with a function \bar{F} defined on \mathbb{N} by

$$\bar{F}(\alpha) = \sum_{u \in \mathbb{Z}} F(u, u + \alpha)$$

and as a consequence, for \bar{F} and G degree 2 functions it holds that

$$\ll F, G \gg_{\text{free}} = \sum_{\alpha \in \mathbb{N}} \bar{F}(\alpha) \bar{G}(\alpha). \quad (37)$$

Observe that $(\mathfrak{A}_0 F)(u, v)$ is equal to

$$\begin{cases} 2(1 + \lambda) \left(F(u - 1, u) - F(u, u + 1) \right), & \text{if } u = v, \\ (1 + \lambda) \left(F(u - 1, u + 1) - F(u, u + 2) \right) + (2 + \lambda) \left(F(u, u) - F(u + 1, u + 1) \right), \\ \quad \text{if } (u, v) = (u, u + 1), \\ (1 + \lambda) \left(F(u - 1, v) - F(u + 1, v) + F(u, v - 1) - F(u, v + 1) \right), & \text{if } |u - v| \geq 2 \end{cases}$$

and

$$\mathcal{W}(u, u) = \frac{(\lambda + 1)(\lambda + 2)}{2}, \quad \mathcal{W}(u, v) = (\lambda + 1)^2 \quad \text{for } u \neq v. \quad (38)$$

It is then easy to show that

$$\overline{\mathcal{W}^{1/2}(\mathfrak{A}_0 F)}(\alpha) = 0$$

for any $\alpha \in \mathbb{N}$. Putting together the previous result and (37) it follows that:

$$\begin{aligned} \|\mathcal{W}^{1/2} \mathfrak{A}_0 F\|_{-1, z, \text{free}}^2 &= \ll \mathcal{W}^{1/2} \mathfrak{A}_0 F, (z - \gamma \mathfrak{S})^{-1} (\mathcal{W}^{1/2} \mathfrak{A}_0 F) \gg_{\text{free}} \\ &= \sum_{\alpha \in \mathbb{N}} \overline{\mathcal{W}^{1/2}(\mathfrak{A}_0 F)}(\alpha) \overline{(\lambda - \gamma \mathfrak{S})^{-1} [\mathcal{W}^{1/2}(\mathfrak{A}_0 F)]}(\alpha) = 0. \end{aligned}$$

Lemma 4. *There exists a positive constant C such that for every symmetric function F of degree 2, if $\bar{F}(\alpha) = \sum_{z \in \mathbb{Z}} F(z, z + \alpha)$, then*

$$C^{-1} \sum_{\substack{x, y \neq 0, \\ |x - y| = 1}} \left(\bar{F}(y) - \bar{F}(x) \right)^2 \leq \ll F, -\mathfrak{S}F \gg_{\text{free}} \leq C \sum_{\substack{x, y \neq 0, \\ |x - y| = 1}} \left(\bar{F}(y) - \bar{F}(x) \right)^2.$$

Proof. This follows easily from the following equalities together with (37):

$$\begin{aligned} \overline{\mathfrak{S}F}(0) &= \sum_{y \in \mathbb{Z}} (\mathfrak{S}F)(y, y) \\ &= \sum_{y \in \mathbb{Z}} \left(F(y + 1, y + 1) - F(y, y) \right) + \left(F(y - 1, y - 1) - F(y, y) \right) = 0, \end{aligned}$$

$$\begin{aligned} \overline{\mathfrak{S}F}(1) &= \sum_{y \in \mathbb{Z}} (\mathfrak{S}F)(y, y + 1) \\ &= \sum_{y \in \mathbb{Z}} \left(F(y - 1, y + 1) - F(y, y + 1) \right) + \sum_{y \in \mathbb{Z}} \left(F(y, y + 2) - F(y, y + 1) \right) \\ &= 2 \left(\bar{F}(2) - \bar{F}(1) \right), \end{aligned}$$

$$\begin{aligned}
 \overline{\mathfrak{S}F}(\alpha) &= \sum_y (\mathfrak{S}F)(y, y + \alpha) \\
 &= \sum_{y \in \mathbb{Z}} \left(F(y - 1, y + \alpha) - F(y, y + \alpha) \right) + \sum_{y \in \mathbb{Z}} \left(F(y + 1, y + \alpha) - F(y, y + \alpha) \right) \\
 &+ \sum_{y \in \mathbb{Z}} \left(F(y, y + \alpha + 1) - F(y, y + \alpha) \right) + \sum_{y \in \mathbb{Z}} \left(F(y, y + \alpha - 1) - F(y, y + \alpha) \right) \\
 &= 2 \left(\bar{F}(\alpha + 1) - \bar{F}(\alpha) \right) + 2 \left(\bar{F}(\alpha - 1) - \bar{F}(\alpha) \right), \quad \alpha \geq 2.
 \end{aligned}$$

To any degree 3 function G , i.e. a symmetric function G on \mathbb{Z}^3 , the function \bar{G} is identified with a function on \mathbb{N}^2 :

$$\bar{G}(u, v) = \sum_{y \in \mathbb{Z}} G(y, u + y, u + v + y).$$

Since G is symmetric on \mathbb{Z}^3 , then \bar{G} is symmetric on \mathbb{Z}^2 . As above, for F and G degree 3 functions it holds that

$$\ll F, G \gg_{\text{free}} = \sum_{(\alpha, \beta) \in \mathbb{N}^2} \bar{F}(\alpha, \beta) \bar{G}(\alpha, \beta). \quad (39)$$

Let \mathbb{D}_3 , acting on the local functions on \mathbb{N}^2 , be defined by

$$\begin{aligned}
 \mathbb{D}_3(\bar{G}) &= \sum_{u \geq 1} \left(\bar{G}(u + 1, 0) - \bar{G}(u, 0) \right)^2 + \sum_{v \geq 1} \left(\bar{G}(0, v + 1) - \bar{G}(0, v) \right)^2 \\
 &+ \left(\bar{G}(1, 0) - \bar{G}(0, 1) \right)^2 + \sum_{u, v \geq 1} \left(\bar{G}(u + 1, v) - \bar{G}(u, v) \right)^2 \\
 &+ \left(\bar{G}(u, v + 1) - \bar{G}(u, v) \right)^2.
 \end{aligned} \quad (40)$$

This is the Dirichlet form of a symmetric nearest neighbors random walk on \mathbb{N}^2 where all the jumps between $\{0\} \times \mathbb{N}$ and $\mathbb{N}^* \times \mathbb{N}^{*5}$, all the jumps from $\mathbb{N} \times \{0\}$ and $\mathbb{N}^* \times \mathbb{N}^*$ and all the jumps from 0 have been suppressed, and a jump between $(0, 1)$ and $(1, 0)$ has been added.

Lemma 5. *There exists a constant $C > 0$ such that for any symmetric function G on \mathbb{Z}^3*

$$C^{-1} \mathbb{D}_3(\bar{G}) \leq \ll G, -\mathfrak{S}G \gg_{\text{free}} \leq C \mathbb{D}_3(\bar{G}).$$

⁵ Here and in the sequel $\mathbb{N}^* := \mathbb{N} \setminus \{0\}$

Proof. We have the following equalities

$$\begin{aligned}
\overline{\mathfrak{S}G}(0,0) &= 0, \\
\overline{\mathfrak{S}G}(0,1) &= 2\left(\bar{G}(0,2) - \bar{G}(0,1)\right) + \left(\bar{G}(1,0) - \bar{G}(0,1)\right), \\
\overline{\mathfrak{S}G}(0,\beta) &= 2\left(\bar{G}(0,\beta+1) - \bar{G}(0,\beta)\right) + 2\left(\bar{G}(0,\beta-1) - \bar{G}(0,\beta)\right), \quad \beta \geq 2, \\
\overline{\mathfrak{S}G}(1,0) &= 2\left(\bar{G}(2,0) - \bar{G}(1,0)\right) + \left(\bar{G}(0,1) - \bar{G}(1,0)\right) \\
\overline{\mathfrak{S}G}(\alpha,0) &= 2\left(\bar{G}(\alpha+1) - \bar{G}(\alpha,0)\right) + 2\left(\bar{G}(\alpha-1) - \bar{G}(\alpha,0)\right), \quad \alpha \geq 2, \\
\overline{\mathfrak{S}G}(\alpha,\beta) &= \left(\bar{G}(\alpha+1,\beta) - \bar{G}(\alpha,\beta)\right) + \left(\bar{G}(\alpha,\beta+1) - \bar{G}(\alpha,\beta)\right) \\
&+ \mathbf{1}_{\{\alpha \geq 2\}} \left(\bar{G}(\alpha-1,\beta+1) - \bar{G}(\alpha,\beta)\right) + \mathbf{1}_{\{\alpha \geq 2\}} \left(\bar{G}(\alpha-1,\beta) - \bar{G}(\alpha,\beta)\right) \\
&+ \mathbf{1}_{\{\beta \geq 2\}} \left(\bar{G}(\alpha+1,\beta-1) - \bar{G}(\alpha,\beta)\right) + \mathbf{1}_{\{\beta \geq 2\}} \left(\bar{G}(\alpha,\beta-1) - \bar{G}(\alpha,\beta)\right), \quad \alpha, \beta \geq 1.
\end{aligned}$$

We recognize in these expressions the generator of a symmetric nearest neighbors random walk on \mathbb{N}^2 where

- all the jumps between $\{0\} \times \mathbb{N}$ and $\mathbb{N}^* \times \mathbb{N}^*$, all the jumps between $\mathbb{N} \times \{0\}$ and $\mathbb{N}^* \times \mathbb{N}^*$, and all the jumps from 0 have been suppressed;
- a jump between $(0,1)$ and $(1,0)$ with rate 1 has been added;
- jumps between (α,β) and $(\alpha \pm 1, \beta \mp 1)$ for $(\alpha,\beta) \in \mathbb{N}^* \times \mathbb{N}^*$ with rate 1 have been added.
- the non vanishing jumps on $\mathbb{N} \times \{0\}$ and on $\{0\} \times \mathbb{N}$ have been multiplied by 2.

This together with (39), implies the lemma.

We choose a degree 2 symmetric function F such that

$$\begin{aligned}
\bar{F}(\alpha) &= z^{-1/4} e^{-z^{3/4}(\alpha-1)}, \quad \alpha \geq 1, \\
\bar{F}(0) &= \bar{F}(1).
\end{aligned} \tag{41}$$

This function exists since given a function G defined on \mathbb{N} we can find a symmetric function F defined in \mathbb{Z}^2 such that $\bar{F} = G$. For that purpose, take $F(x,y) = G(|y-x|)[\phi(x) + \phi(y)]$ where the function ϕ is defined on \mathbb{Z} and is such that $\sum_{x \in \mathbb{Z}} \phi(x) = 1/2$. Then for any $\alpha \in \mathbb{N}$, $\bar{F}(\alpha) = \sum_{u \in \mathbb{Z}} F(u, u+\alpha) = G(\alpha) \sum_{u \in \mathbb{Z}} [\phi(u) + \phi(u+\alpha)] = G(\alpha)$.

Observe that with this choice, by Lemma 4,

$$\ll F, -\mathfrak{S}F \gg_{\text{free}} \sim z^{1/4}, \quad \bar{F}(1) = z^{-1/4}, \quad z \sum_{\alpha \in \mathbb{N}} \bar{F}^2(\alpha) \sim z^{-1/4}. \tag{42}$$

It remains to estimate the last contribution given by $\|\mathcal{W}^{1/2}G\|_{-1,z,\text{free}}^2$ where $G = \mathfrak{A}_+ F$ is a degree 3 function.

Lemma 6. *Let $G = \mathfrak{A}_+ F$ where F is defined by (41). There exists a constant $C > 0$ such that*

$$\|\mathcal{W}^{1/2}G\|_{-1,z,\text{free}}^2 \geq Cz^{-1/4}.$$

Proof. For any $u, v, w \in \mathbb{Z}$, we have

$$\begin{aligned} G(u, u+1, u+2) &= F(u, u+1) - F(u+1, u+2), \\ G(u, u+1, v) &= F(u, v) - F(u+1, v), \quad v > u+1, \\ G(v, u, u+1) &= F(v, u) - F(v, u+1), \quad v < u, \\ G(u, u, u+1) &= 2\left(F(u, u) - F(u, u+1)\right), \\ G(u, u, u-1) &= 2\left(F(u-1, u) - F(u, u)\right), \\ G(u, v, w) &= 0 \quad \text{otherwise.} \end{aligned}$$

Let us now compute $\bar{G}(u, v)$, $u, v \in \mathbb{N}$. We get

$$\begin{aligned} \bar{G}(0, 1) &= -\bar{G}(1, 0) = 2\bar{F}(0) - 2\bar{F}(1), \\ \bar{G}(1, v) &= \bar{F}(v+1) - \bar{F}(v), \quad v \geq 2, \\ \bar{G}(u, 1) &= \bar{F}(u) - \bar{F}(u+1), \quad u \geq 2, \\ \bar{G}(u, v) &= 0 \quad \text{otherwise.} \end{aligned} \tag{43}$$

By (41) we have that $\bar{G}(0, 1) = \bar{G}(1, 0) = 0$. Also notice that $\bar{G}(u, u) = 0$ and by (38) we have that $\mathcal{W}^{1/2}(u, v) = (1 + \lambda)$ for $u \neq v$.

It follows, by Lemma 5, that $\|\mathcal{W}^{1/2}G\|_{-1, z, \text{free}}^2$ is upper bounded by the variational formula:

$$\begin{aligned} \|\mathcal{W}^{1/2}G\|_{-1, z, \text{free}}^2 &= \sup_R \left\{ 2 \sum_{(u, v) \in \mathbb{N}^2} R(u, v) \mathcal{W}^{1/2}(u, v) \bar{G}(u, v) - C_0 \mathbb{D}_3(R) \right\} \\ &= \sup_R \left\{ 2(1 + \lambda) \sum_{(u, v) \in \mathbb{N}^2} R(u, v) \bar{G}(u, v) - C_0 \mathbb{D}_3(R) \right\} \end{aligned}$$

where the supremum is taken over local functions on \mathbb{N}^2 . By (43), we have that

$$\begin{aligned} &\sum_{(u, v) \in \mathbb{N}^2} R(u, v) \bar{G}(u, v) \\ &= \sum_{v \geq 2} R(1, v) \left(\bar{F}(v+1) - \bar{F}(v) \right) - \sum_{u \geq 2} R(u, 1) \left(\bar{F}(u+1) - \bar{F}(u) \right) \\ &= \sum_{v \geq 3} \bar{F}(v) \left(R(1, v-1) - R(1, v) \right) - \sum_{u \geq 3} \bar{F}(u) \left(R(u-1, 1) - R(u, 1) \right) \quad (44) \\ &+ \bar{F}(2) \left(R(2, 1) - R(1, 2) \right) \\ &= \sum_{v \geq 2} \bar{F}(v) \left(R(1, v-1) - R(1, v) \right) - \sum_{u \geq 2} \bar{F}(u) \left(R(u-1, 1) - R(u, 1) \right). \end{aligned}$$

We use now the following parametrization of R . For $k \geq 1$, $v \in \mathbb{Z}$, let us define

$$R(k, v) = \phi(k-1, v-k), \quad v \geq k, \quad R(u, k) = \phi(k-1, -u+k), \quad u \geq k,$$

where $\{\phi(k, \cdot) ; k \geq 0\}$ are functions from $\mathbb{Z} \rightarrow \mathbb{R}$. We have the following lower bound for $\mathbb{D}_3(R)$:

$$\mathbb{D}_3(R) \geq \sum_{u, v \geq 1} \left(R(u+1, v) - R(u, v) \right)^2 + \left(R(u, v+1) - R(u, v) \right)^2$$

which is nothing but the Dirichlet form of a random walk where only jumps connecting sites of $\mathbb{N}^* \times \mathbb{N}^*$ have been conserved. With the choice of the parametrization for R and this lower bound, it is not difficult to show there exists a constant $C > 0$ such that

$$\mathbb{D}_3(R) \geq C \sum_{k \geq 0} \sum_{v \in \mathbb{Z}} \left(\phi(k, v+1) - \phi(k, v) \right)^2 + \left(\phi(k+1, v) - \phi(k, v) \right)^2.$$

The right hand side of the previous inequality is the Dirichlet form of a symmetric simple random walk on $\mathbb{N} \times \mathbb{Z}$.

By (44), we get

$$\sum_{(u, v) \in \mathbb{N}^2} R(u, v) \bar{G}(u, v) = \sum_{u \in \mathbb{Z}} \phi(0, u) \left(\tilde{F}(u-1) - \tilde{F}(u) \right)$$

where $\tilde{F} : \mathbb{Z} \rightarrow \mathbb{R}$ is defined by $\tilde{F}(u) = -\bar{F}(u+2)\mathbf{1}_{\{u \geq 0\}} - \bar{F}(1-u)\mathbf{1}_{\{u \leq -1\}}$. We extend the function ϕ defined on $\mathbb{N} \times \mathbb{Z}$ to \mathbb{Z}^2 by defining $\phi(-k, u) = \phi(k, u)$, $k \geq 1, u \in \mathbb{Z}$. Observe then that

$$\begin{aligned} \mathbb{D}_3(R) &\geq C \sum_{k \geq 0} \sum_{v \in \mathbb{Z}} \left(\phi(k, v+1) - \phi(k, v) \right)^2 + \left(\phi(k+1, v) - \phi(k, v) \right)^2 \\ &= \frac{C}{2} \sum_{k \in \mathbb{Z}} \sum_{v \in \mathbb{Z}} \left(\phi(k, v+1) - \phi(k, v) \right)^2 + \left(\phi(k+1, v) - \phi(k, v) \right)^2. \end{aligned}$$

Consequently we have, for suitable positive constants C_1, C_2 :

$$\begin{aligned} \|\mathcal{W}^{1/2} G\|_{-1, z, \text{free}}^2 &\leq C_1 \sup_{\phi} \left\{ 2 \sum_{u \in \mathbb{Z}} \phi(0, u) \left(\tilde{F}(u-1) - \tilde{F}(u) \right) \right. \\ &\quad \left. - C_2 \sum_{\substack{(\mathbf{u}, \mathbf{v}) \in \mathbb{Z}^2 \\ |\mathbf{u} - \mathbf{v}| = 1}} \left(\phi(\mathbf{u}) - \phi(\mathbf{v}) \right)^2 \right\}. \end{aligned} \quad (45)$$

A standard Fourier computation shows this supremum is of order $z^{-1/4}$. Indeed, let \hat{u} be the Fourier transform of the function $u : \mathbb{Z}^n \rightarrow \mathbb{R}$, defined by

$$\hat{u}(\mathbf{k}) = \sum_{\mathbf{x} \in \mathbb{Z}^n} e^{2i\pi \mathbf{x} \cdot \mathbf{k}} u(\mathbf{x}), \quad \mathbf{k} = (k_1, \dots, k_n),$$

and denote by $\hat{u}^*(\mathbf{k})$ the complex conjugate of $\hat{u}(\mathbf{k})$. Using the expression of the sum of a convergent geometric series, we obtain the following expression for the Fourier transform $\Psi(k_1, k_2)$ of the function $(x, y) \in \mathbb{Z}^2 \rightarrow \delta_0(y) \bar{F}(x)$:

$$\Psi(k_1, k_2) = -z^{-1/4} e^{-z^{3/4}} \left\{ \frac{1}{1 - e^{2i\pi k_1} e^{-z^{3/4}}} - \frac{e^{-2i\pi k_1}}{1 - e^{-2i\pi k_1} e^{-z^{3/4}}} \right\}$$

which satisfies

$$|\Psi(k_1, k_2)| \leq \frac{C_3 \sqrt{z}}{z^{3/2} + C_4 \sin^2(\pi k_1)}$$

for some positive constants C_3, C_4 . The supremum appearing in (45) is then given by

$$C_2^{-1} \int_{[0,1]^2} \frac{|\Psi(k_1, k_2)|^2}{z + 4 \sin^2(\pi k_1) + 4 \sin^2(\pi k_2)} dk_1 dk_2.$$

Then the result follows by a standard study of this integral.

To obtain (34), by (42) and Lemma 6, it suffices to take a test function in the form aF with F given by (41) and a sufficiently small.

8. Stochastic perturbations of Hamiltonian systems

In this section we discuss some other possible stochastic perturbations and make some connections with the recent models considered in [8]. Let us start with the Hamiltonian system (2) with potential V and generator A given by

$$A = \sum_{x \in \mathbb{Z}} \left(V'(\eta_{x+1}) - V'(\eta_{x-1}) \right) \partial_{\eta_x}.$$

The energy $\sum_{x \in \mathbb{Z}} V(\eta_x)$ and the volume $\sum_{x \in \mathbb{Z}} \eta_x$ are conserved by these dynamics. Remark that in fact $\sum_{x \in \mathbb{Z}} \eta_{2x}$ and $\sum_{x \in \mathbb{Z}} \eta_{2x+1}$ are also conserved and that we cannot exclude the case that still many others exist. This is the case for example for the exponential interaction for which an infinite number of conserved quantities can be explicitly identified. Anyway, we are only interested in these two first quantities. The product probability measures $\mu_{\beta, \lambda}$ defined by

$$\mu_{\beta, \lambda}(d\eta) = \prod_{x \in \mathbb{Z}} Z(\beta, \lambda)^{-1} \exp \{ -\beta V(\eta_x) - \lambda \eta_x \} d\eta_x,$$

where

$$Z(\beta, \lambda) = \int_{-\infty}^{+\infty} \exp(-\beta V(r) - \lambda r) dr.$$

are invariant for the infinite dynamics.

In [7] we proposed to perturb this deterministic dynamics by the Poissonian noise considered in this paper and conserving both the energy and the volume. One could also consider the ‘‘Brownian’’ noise whose generator S is given by $S = \sum_{x \in \mathbb{Z}} Y_x^2$ where

$$Y_x = (V'(\eta_{x+1}) - V'(\eta_{x-1})) \partial_{\eta_x} + (V'(\eta_{x-1}) - V'(\eta_x)) \partial_{\eta_{x+1}} + (V'(\eta_x) - V'(\eta_{x+1})) \partial_{\eta_{x-1}},$$

is the vector field tangent to the curve

$$\left\{ (\eta_{x-1}, \eta_x, \eta_{x+1}) \in \mathbb{R}^3; \sum_{y=x-1}^{x+1} \eta_y = 0, \sum_{y=x-1}^{x+1} V(\eta_y) = 1 \right\}.$$

It is easy to see that the process with generator $L = A + S$ conserves the energy and the volume and has $\mu_{\beta, \lambda}$ as invariant measures. A priori, it should

be possible to extend our result to this system for V of exponential type but the noise S seems to have a quite complicated expression in the orthogonal basis we used in this paper. The advantage of the Poissonian noise is its very simple form. Notice also that the Poissonian noise is a weaker perturbation of the Hamiltonian dynamics than the Brownian noise in the sense it is less mixing. Indeed, consider the discrete torus \mathbb{T}^N of length N and the Brownian noise $S^N = \sum_{x \in \mathbb{T}^N} Y_x^2$ restricted to the manifold $\mathcal{M}_{\pi, E}^N$ defined by

$$\mathcal{M}_{\pi, E}^N = \left\{ \eta \in \mathbb{R}^{\mathbb{T}^N} ; \sum_{y \in \mathbb{T}^N} \eta_y = \pi, \sum_{y \in \mathbb{T}^N} V(\eta_y) = E \right\}, \quad E > 0, \pi \in \mathbb{R}.$$

Then S^N is ergodic on $\mathcal{M}_{\pi, E}^N$ but this is not true for the restriction of the Poissonian noise restricted to $\mathcal{M}_{\pi, E}^N$.

We could also decide to conserve energy and not the volume by adding a suitable perturbation. The invariant states are then given by $\mu_{\beta, 0}$, $\beta > 0$. If V is even, a simple Poissonian noise consists to change the sign of η_x independently on each site x at random exponential times. In this case one can prove, as in [6], that the energy diffuses in the sense that the Green-Kubo formula converges to a well defined finite value. For a generic V a Brownian noise with generator S given by $S = \sum_{x \in \mathbb{Z}} K_x^2$ with $K_x = V'(\eta_{x+1})\partial_{\eta_x} - V'(\eta_x)\partial_{\eta_{x+1}}$ makes the job.

Consider now the case where we want to add a stochastic perturbation conserving only the volume. It does not seem to be easy to define a simple Poissonian noise with such a property. A Brownian noise is obtained by the following scheme. Fix $\beta > 0$, consider the vector field $X_x = \partial_{\eta_{x+1}} - \partial_{\eta_x}$ which is tangent to the hyperplane $\{(\eta_x, \eta_{x+1}) \in \mathbb{R}^2 ; \eta_x + \eta_{x+1} = 1\}$ and define the Langevin operator S_β by

$$\begin{aligned} S_\beta &= \frac{1}{2} \sum_{x \in \mathbb{Z}} e^{-\mathcal{H}_{\beta, \lambda}} X_x (e^{\mathcal{H}_{\beta, \lambda}} X_x) \\ &= \frac{1}{2} \sum_{x \in \mathbb{Z}} X_x^2 + \frac{\beta}{2} \sum_{x \in \mathbb{Z}} \left(V'(\eta_{x+1}) - V'(\eta_x) \right) X_x \end{aligned}$$

where $\mathcal{H}_{\beta, \lambda} = \beta \sum_{x \in \mathbb{Z}} V(\eta_x) + \lambda \sum_{x \in \mathbb{Z}} \eta_x$. Observe that S_β depends on β but is independent of λ . The operator S_β is a nonpositive self-adjoint operator in $\mathbb{L}^2(\mu_{\beta, \lambda})$ for any λ and $S_\beta(\sum_{x \in \mathbb{Z}} \eta_x) = 0$. Then, the perturbed volume-conserving model has a generator L_β^V given by

$$L_\beta^V = A + \gamma S_\beta \tag{46}$$

where $\gamma > 0$ is a parameter fixing the strength of the noise. By construction, the Markov process generated by L_β^V has $\mu_{\beta, \lambda}$ as a set of invariant probability measures. In fact, using the same methods as in [7, 17] one can prove that the only space-time invariant probability measures with finite local entropy density are mixtures of the $(\mu_{\beta, \lambda})_\lambda$. We can also rewrite L_β^V as

$$\begin{aligned} L_\beta^V &= \sum_{x \in \mathbb{Z}} \left\{ \left(1 - \frac{\gamma\beta}{2} \right) V'(\eta_{x+1}) + \gamma\beta V'(\eta_x) - \left(1 + \frac{\gamma\beta}{2} \right) V'(\eta_{x-1}) \right\} \partial_{\eta_x} \\ &\quad + \gamma \sum_{x \in \mathbb{Z}} (\partial_{\eta_x}^2 - \partial_{\eta_x, \eta_{x+1}}^2). \end{aligned}$$

The microscopic flux $j_{x,x+1}$ associated to the volume conservation law is defined by

$$L_\beta^V(\eta_x) = -\nabla j_{x-1,x}, \quad j_{x-1,x} = -\left(1 + \frac{\gamma\beta}{2}\right) V'(\eta_{x-1}) - \left(1 - \frac{\gamma\beta}{2}\right) V'(\eta_x).$$

The semi-discrete directed polymer model considered in [8] is, up to an irrelevant scaling factor 2, recovered by taking $V(\eta) = e^{-\eta}$, $\beta = 1$ and $\gamma = 2$ (see (3.7) in [29]). In [8] the authors show that for a particular non stationary initial condition (“wedge”), by developing a very nice theory of Macdonald processes, the system belongs to the Kardar-Parisi-Zhang universality class ([29]). Unfortunately one can not use their results or their methods to derive a more precise picture for the model with exponential interactions considered in this paper. For other potentials V the theory developed by Borodin and Corwin in [8] can not be adapted but it would be very interesting to see if one can relate the models generated by L_β^V to the semi-discrete directed polymer and deduce some qualitative information from the latter. The use of the variational formulas considered in this paper could be the way.

A. Existence of the infinite dynamics

In this section we prove existence of the infinite volume dynamics $(\xi(t))_{t \geq 0}$. We focus here on the process ξ but the same proof can be carried for the process η (or just define η in terms of ξ by $\eta_x(t) = -b^{-1} \log \xi_x(t)$, $x \in \mathbb{Z}$. To simplify notations we will assume $b = 1$).

Since the interaction coming from the deterministic part is non-quadratic at infinity, proving the existence of the infinite dynamics is a non trivial task. Nevertheless nice sophisticated techniques have been introduced by Dobrushin and Fritz in [12]. Here, we follow closely the approach of [13] (see also [14,15]) adapted to our case. By itself, the strategy of the proof of existence of solutions is standard: we consider finite subsystems and prove compactness of this family by means of an a priori bound for a quantity \bar{E} which plays the role of an energy density. The obtention of this a priori bound is however non trivial and is the main step to get the existence of the dynamics. The aim of this appendix is to show how to get such an a priori bound. The a priori bound we derive here for the infinite dynamics is also valid for finite subsystems corresponding to a finite set $\Lambda \subset \mathbb{Z}$ with a bound which is independent of the size of Λ . This proves then that the finite subsystems form a compact family from which one can extract a subsequence converging to the infinite dynamics.

We have first to specify the space of allowed configurations $\Omega \subset (0, +\infty)^\mathbb{Z}$. For $x \in \mathbb{Z}$, let $g(x) = 1 + \log(1 + |x|)$ and denote by $E(\xi, \mu, \sigma)$, $\xi \in (0, +\infty)^\mathbb{Z}$, $\mu \in \mathbb{Z}$, $\sigma > 0$, the quantities

$$E(\xi, \mu, \sigma) = \sum_{|x-\mu| \leq \sigma} (1 + 2\xi_x - \log(\xi_x)),$$

$$\bar{E}(\xi) = \sup_{\mu \in \mathbb{Z}} \sup_{\sigma \geq g(\mu)} \sigma^{-1} E(\xi, \mu, \sigma).$$

The quantity \bar{E} is called the logarithmic fluctuation of energy and the set Ω is defined as

$$\Omega := \{\xi \in (0, +\infty)^\mathbb{Z} : \bar{E}(\xi) < +\infty\}.$$

The configuration space Ω is equipped with the product topology and with the associated Borel structure. It is easy to see that $\nu_{\beta,\lambda}(\Omega) = 1$ for any $\beta > 0$ and $\lambda > -1$.

Let $\mathbf{N}(t) = \{N_{x,x+1}(t); x \in \mathbb{Z}\}$ be a collection of independent Poisson processes of intensity $\gamma > 0$. The equations of motion corresponding to the generator \mathcal{L} read as

$$d\xi_x = \xi_x(\xi_{x+1} - \xi_{x-1})dt + \nabla((\xi_x - \xi_{x-1})dN_{x-1,x}(t)), \quad x \in \mathbb{Z}. \quad (47)$$

Let $D(\mathbb{R}_+, \mathbb{R})$ denote the space of càdlàg functions of \mathbb{R}_+ into \mathbb{R} with the Skorohod topology and let $\mathbb{D} = [D(\mathbb{R}_+, \mathbb{R})]^\mathbb{Z}$ equipped with the product topology and the associated Borel field \mathcal{B} . The smallest σ -algebra on which all projections restricted to the time interval $[0, t]$ are measurable will be denoted by \mathcal{B}_t . Finally, suppose that we are given a probability measure \mathbf{P} on \mathcal{B} such that our Poisson processes $N_{x,x+1}$ are realized as components of the random element of \mathbb{D} .

Definition 1. A \mathcal{B}_t -adapted mapping $\xi(t) := \xi(t, \mathbf{N})$ of \mathbb{D} into itself is called a tempered solution of (47) with initial configuration $\xi^0 \in \Omega$ if $\xi(0) = \xi^0$, almost each trajectory $\xi(\cdot, \mathbf{N})$ satisfies the integral form of (47), and the logarithmic energy fluctuation $\bar{E}(\xi(t))$ is bounded on finite intervals of time with probability one.

Theorem 5 For any $\xi^0 \in \Omega$, there exists a unique tempered solution of (47) with initial configuration $\xi^0 \in \Omega$.

As explained above, the main step to prove this theorem is to obtain an a priori bound that we prove in Proposition 3. For a complete proof, we refer to [13] (or [14,15]).

Now we notice that the Gibbs state $\nu_{\beta,\lambda}$, $(\beta, \lambda) \in (0, +\infty) \times (-1, +\infty)$ is formally invariant for the infinite dynamics generated by $(\xi(t))_{t \geq 0}$. This can be seen by observing that $\int(\mathcal{L}f)(\xi)d\nu_{\beta,\lambda}(d\xi) = 0$ for nice functions $f : \Omega \rightarrow \mathbb{R}$. Nevertheless, some care has to be taken to prove this. Indeed, we do not know that \mathcal{L} is really the generator of the semigroup generated by $(\xi(t))_{t \geq 0}$ on the space of bounded measurable functions on Ω in the usual Hille-Yosida theory. This can be a very difficult question that we prefer to avoid (see [15]). Instead we use the fact that the infinite dynamics can be approximated by finite subsystems.

Proposition 2 For any $\beta > 0, \lambda > -1$, the probability measure $\nu_{\beta,\lambda}$ is invariant for the process $(\xi(t))_{t \geq 0}$.

Proof. Let $n \geq 2$ and consider the local dynamics generated by the generator $\mathcal{L}_n = \mathcal{A}_n + \gamma\mathcal{S}_n$ where

$$\begin{aligned} (\mathcal{A}_n f)(\xi) &= \sum_{x=-n}^n \xi_x(\xi_{x+1} - \xi_{x-1})\partial_{\xi_x} f(\xi) \\ &\quad - \xi_{n+1} \left(\xi_n + \frac{\lambda+1}{\beta} \right) \partial_{\xi_n} f(\xi) + \xi_{-n-1} \left(\xi_{-n} + \frac{\lambda+1}{\beta} \right) \partial_{\xi_{-n-1}} f(\xi), \\ (\mathcal{S}_n f)(\xi) &= \sum_{x=-n}^n \left(f(\xi^{x,x+1}) - f(\xi) \right) \end{aligned}$$

where $f : \Omega \rightarrow \mathbb{R}$ is a compactly supported continuously differentiable function. The dynamics is essentially finite-dimensional since the particles outside the box $\{-n-1, \dots, n+1\}$ are frozen. Thus, the classical Hille-Yosida theory can be applied. The boundary conditions have been chosen to have

$$\int (\mathcal{L}_n f)(\xi) d\nu_{\beta, \lambda}(\xi) = 0$$

for any compactly supported continuously differentiable function f which shows that $\nu_{\beta, \lambda}$ is invariant for the local dynamics. Since, as a consequence of the a priori bound, the infinite dynamics is obtained as a limit of finite local dynamics, this implies that $\nu_{\beta, \lambda}$ is invariant for the infinite dynamics.

Then this defines a strongly continuous semigroup of contractions $(P_t)_{t \geq 0}$ on the Hilbert space $\mathbb{L}^2(\Omega, \mathcal{B}, \nu_{\beta, \lambda})$. Moreover, Itô's formula shows that its generator is a closable extension of \mathcal{L} given by $\mathcal{A} + \gamma \mathcal{S}$ since for any local compactly supported continuously differentiable function f , we have

$$(P_t f)(\xi) = f(\xi) + \int_0^t (P_s \mathcal{L} f)(\xi) ds, \quad \xi \in \Omega, \quad t \geq 0.$$

A.1. Logarithmic energy fluctuation. We have first to consider a clever smooth modification of \bar{E} . Let $0 < \lambda < 1$ and consider a twice continuously differentiable nonincreasing function $\varphi : \mathbb{R} \rightarrow (0, 1)$ such that $\varphi(u) = e^{\lambda(1-u)}$ if $u \geq 2$, $\varphi(u) = (1 + \lambda + \lambda^2/2)e^{-\lambda}$ if $u \leq 1$, and φ is concave for $u \leq 3/2$, convex if $u \geq 3/2$. Finally, $0 \leq -\varphi'(u) \leq \lambda\varphi(u) \leq e^{\lambda(1-u)}$, $\varphi(u) \geq e^{-\lambda(1+u)}$ and $|\varphi''(u)| \leq \varphi(u)$ for all $u > 0$.

For $x \in \mathbb{Z}$ and $\sigma \geq 1$ we define the function f as

$$f(x, \sigma) = \int_{\mathbb{R}} \varphi(|x-y|/\sigma) e^{-2\lambda|y|} dy.$$

In [15] are proved the following properties on f :

$$\begin{aligned} c_1 \exp(-\lambda|x|/\sigma) &\leq f(x, \sigma) \leq c_2 \exp(-\lambda|x|/\sigma), \\ f(x, \sigma) &\leq f(y, \sigma) e^{2\lambda|x-y|}, \quad \partial_\sigma f(x, \sigma) \leq e^{2\lambda|x-y|} \partial_\sigma f(y, \sigma), \\ |\partial_x f(x, \sigma)| &\leq \min\{\partial_\sigma f(x, \sigma), \sigma^{-1} f(x, \sigma)\}, \\ g(x) |\partial_x f(x - \mu, \sigma)| &\leq 4g(|\mu| + \sigma) (\partial_\sigma f)(x - \mu, \sigma). \end{aligned} \tag{48}$$

Here the constants depend only on λ .

For $\xi \in (0, +\infty)^{\mathbb{Z}}$, $\mu \in \mathbb{Z}$ and $\sigma > 0$, consider the function

$$W(\xi, \mu, \sigma) = \sum_{x \in \mathbb{Z}} f(x - \mu, \sigma) (1 + 2\xi_x - \log \xi_x) \tag{49}$$

and let

$$\bar{W}(\xi) = \sup_{\mu \in \mathbb{Z}} \sup_{\sigma \geq g(\mu)} \left\{ \sigma^{-1} W(\xi, \mu, \sigma) \right\}. \tag{50}$$

Observe that by (48),

$$W(\xi, \mu, \sigma) \geq c_1 e^{-\lambda} E(\xi, \mu, \sigma), \quad (51)$$

for all $\xi \in (0, +\infty)^{\mathbb{Z}}$, $\mu \in \mathbb{Z}$ and $\sigma > 0$.

For $\xi \in (0, +\infty)^{\mathbb{Z}}$, we also consider the function

$$\widehat{W}(\xi) = \sup_{\mu \in \mathbb{Z}} \left\{ \frac{W(\xi, \mu, g(\mu))}{g(\mu)} \right\}. \quad (52)$$

The following lemma shows that these two modifications of the logarithmic energy fluctuation are equivalent to \bar{E} .

Lemma 7. *There exists a constant C such that for all $\xi \in (0, +\infty)^{\mathbb{Z}}$:*

$$C^{-1} \widehat{W}(\xi) \leq \bar{W}(\xi) \leq C \widehat{W}(\xi), \quad C^{-1} \bar{E}(\xi) \leq \bar{W}(\xi) \leq C \bar{E}(\xi).$$

Proof. The inequality $\widehat{W}(\xi) \leq \bar{W}(\xi)$ for all $\xi \in (0, +\infty)^{\mathbb{Z}}$, is trivial. Let us prove the second one by taking $\sigma \geq g(\mu)$, $\mu \in \mathbb{Z}$ and denoting $1 + 2\xi_x - \log \xi_x$ by H_x . By (48), we have

$$\begin{aligned} W(\xi, \mu, \sigma) &\leq c_2 \sum_{x \in \mathbb{Z}} \exp(-\lambda|x - \mu|/\sigma) H_x = c_2 \sum_{n=0}^{\infty} e^{-\lambda n/\sigma} \sum_{|x-\mu|=n} H_x \\ &= c_2 (1 - e^{-\lambda/\sigma}) \sum_{n=0}^{\infty} e^{-\lambda n/\sigma} \sum_{|x-\mu| \leq n} H_x, \end{aligned}$$

where the last equality follows from $\sum_{|x-\mu|=n} H_x = \sum_{|x-\mu| \leq n} H_x - \sum_{|x-\mu| \leq n-1} H_x$ and a discrete integration by parts. Let $r \geq 1$ be the integer such that $r - 1 < g(\mu) \leq r$ and decompose the set $\{x \in \mathbb{Z}; |x - \mu| \leq n\}$ as $\cup_{j=1}^{K+1} \Lambda_j$ where the Λ_j are non intersecting intervals of length r for $j = 1, \dots, K$ and Λ_{K+1} is of length at most $r - 1$. Observe that $K + 1$ is of order $n/g(\mu)$. By using (51), we have easily that

$$\sum_{x \in \Lambda_j} H_x \leq C g(\mu) \widehat{W}(\xi)$$

where C depends only on λ . Thus we get

$$W(\xi, \mu, \sigma) \leq C (1 - e^{-\lambda/\sigma}) \sum_{n=0}^{\infty} e^{-\lambda n/\sigma} n \widehat{W}(\xi) \leq C' \sigma \widehat{W}(\xi)$$

which concludes the proof of the second inequality.

The proof of $C^{-1} \bar{E}(\xi) \leq \bar{W}(\xi) \leq C \bar{E}(\xi)$ for all $\xi \in (0, +\infty)^{\mathbb{Z}}$, is the same. The first inequality follows from (51) and the constant can be taken equal to $c_1 e^{-\lambda}$. The second inequality follows from a similar argument to the one used above.

A.2. The a priori bound.

Proposition 3 (A priori bound) *For each $w \geq 1$ there exists a continuous function $q_w(t)$, $t \geq 0$, such that*

$$\mathbf{P} \left\{ \sup_{0 \leq s \leq t} \bar{W}(\xi(s)) > \exp(q_w(t)g(u)) \right\} \leq e^{-u}$$

for each $u \geq 1$, $t \geq 0$, whenever $\bar{W}(\xi^0) \leq w$ and $(\xi(t))_{t \geq 0}$ is a tempered solution of (47) with initial condition ξ^0 .

Proof. We consider a tempered solution $(\xi(t))_{t \geq 0}$ of (47) with initial configuration $\xi^0 \in \Omega$.

For each $k \geq 1$, $\mu \in \mathbb{Z}$ and $t \geq 0$ we define the stochastic process ρ_k by

$$\rho_k(t) = k g(\mu) - C_0 \int_0^t g(|\mu| + |\rho_k(s)|) Z'(s) ds \quad (53)$$

where $C_0 := C_0(\gamma, \lambda)$ is a positive constant that will be chosen later and

$$Z(t) = \int_0^t \bar{W}(\xi(s)) ds.$$

Since the function $f(\cdot)$ is positive, $\bar{W}(\cdot)$ is also positive and this turns $Z(\cdot)$ positive. The trajectories of ρ_k are differentiable, decreasing and satisfy $\rho_{k+1}(t) - \rho_k(t) \leq g(\mu)$ a.s. for each $t \geq 0$. We consider also the sequence of stopping times $\tau_k = \inf\{t \geq 0; \rho_k(t) \leq g(\mu)\}$ which satisfy $\tau_k < \tau_{k+1} < +\infty$ and $\lim_{k \rightarrow \infty} \tau_k = \infty$ a.s. We evaluate now the stochastic differential of $t \rightarrow W(\xi(t), \mu, \rho_k(t))$ for $t \leq \tau_k$ (so that $\rho_k(t) \geq 1$). This is given by

$$d[W(\xi(t), \mu, \rho_k(t))] = I_0^{(k)}(t) dt - C_0 (\partial_\sigma W)(\xi(t), \mu, \rho_k(t)) g(|\mu| + \rho_k(t)) \bar{W}(\xi(t)) dt + dI_1^{(k)}(t)$$

where

$$I_0^{(k)}(t) = 2 \sum_{x \in \mathbb{Z}} \left(f(x - \mu, \rho_k(t)) - f(x + 1 - \mu, \rho_k(t)) \right) \xi_x(t) \xi_{x+1}(t) + \sum_{x \in \mathbb{Z}} \left(f(x + 1 - \mu, \rho_k(t)) - f(x - 1 - \mu, \rho_k(t)) \right) \xi_x(t) \quad (54)$$

and

$$dI_1^{(k)} = \sum_{x \in \mathbb{Z}} f(x - \mu, \rho_k) \left\{ 2 \nabla \left((\xi_x - \xi_{x-1}) dN_{x-1, x} \right) - \nabla \left((\log \xi_x - \log \xi_{x-1}) dN_{x-1, x} \right) \right\}.$$

We first estimate the term $I_0^{(k)}(t)$ and we show that if C_0 is taken sufficiently large then, for $t \leq \tau_k$ we have that

$$I_0^{(k)}(t) - C_0 (\partial_\sigma W)(\xi(t), \mu, \rho_k(t)) g(|\mu| + \rho_k(t)) \bar{W}(\xi(t)) \leq 0. \quad (55)$$

The second term on the right hand side of (54) can be estimated, by using (48) and (48), to get to

$$\begin{aligned}
\left| f(x+1-\mu, \rho_k(t)) - f(x-1-\mu, \rho_k(t)) \right| &= \left| \int_{-1}^1 (\partial_x f)(x-\mu+\alpha, \rho_k(t)) d\alpha \right| \\
&\leq \int_{-1}^1 \left| (\partial_x f)(x-\mu+\alpha, \rho_k(t)) \right| d\alpha \\
&\leq \int_{-1}^1 (\partial_\sigma f)(x-\mu+\alpha, \rho_k(t)) d\alpha \\
&\leq 2 \sup_{[x-\mu-1, x-\mu+1]} \left\{ \partial_\sigma f(\cdot, \rho_k(t)) \right\} \\
&\leq 2e^{2\lambda} \partial_\sigma f(x-\mu, \rho_k(t))
\end{aligned} \tag{56}$$

which gives us that

$$\begin{aligned}
\sum_{x \in \mathbb{Z}} \left(f(x+1-\mu, \rho_k(t)) - f(x-1-\mu, \rho_k(t)) \right) \xi_x(t) &\leq C \sum_{x \in \mathbb{Z}} \partial_\sigma f(x-\mu, \rho_k(t)) \xi_x(t) \\
&\leq C (\partial_\sigma W)(\xi(t), \mu, \rho_k(t)).
\end{aligned}$$

Now, notice that for any $x \in \mathbb{Z}$ and for all $\xi \in (0, \infty)^\mathbb{Z}$ we have that

$$\bar{W}(\xi) \geq \hat{W}(\xi) \geq \frac{W(\xi, x, g(x))}{g(x)}.$$

On the other hand, by (51) and since for all $x > 0$ it holds that $\log(x) \leq 1+x$, then we have that $W(\xi, x, g(x)) \geq c_1 e^\lambda E(\xi, x, g(x)) \geq c_1 e^{-\lambda} \xi_{x+1}$. Then, we conclude that there exists a constant C such that for all $x \in \mathbb{Z}$ and $\xi \in (0, \infty)^\mathbb{Z}$,

$$\xi_{x+1} \leq C g(x) \bar{W}(\xi). \tag{57}$$

To estimate the first term on the right hand side of (54) we use the previous estimate, (48) and a similar argument as done in (56). It follows that

$$\begin{aligned}
&\left| \sum_{x \in \mathbb{Z}} (f(x-\mu, \rho_k(t)) - f(x+1-\mu, \rho_k(t))) \xi_x(t) \xi_{x+1}(t) \right| \\
&\leq C \bar{W}(\xi(t)) g(|\mu| + \rho_k(t)) \sum_{x \in \mathbb{Z}} \partial_\sigma f(x-\mu, \rho_k(t)) \xi_x(t) \\
&\leq C \bar{W}(\xi(t)) g(|\mu| + \rho_k(t)) (\partial_\sigma W)(\xi(t), \mu, \rho_k(t)).
\end{aligned}$$

Then, (55) follows.

The term $dI_1^{(k)}$ can be written as

$$\begin{aligned}
dI_1^{(k)} &= \sum_{x \in \mathbb{Z}} f(x-\mu, \rho_k) \left\{ 2\nabla((\xi_x - \xi_{x-1}) dN_{x-1, x}) - \nabla((\log \xi_x - \log \xi_{x-1}) dN_{x-1, x}) \right\} \\
&= - \sum_{x \in \mathbb{Z}} (f(x+1-\mu, \rho_k) - f(x-\mu, \rho_k)) \{ 2\nabla \xi_x - \nabla \log \xi_x \} dN_{x, x+1} \\
&= - \sum_{x \in \mathbb{Z}} (f(x+1-\mu, \rho_k) - f(x-\mu, \rho_k)) \{ 2\nabla \xi_x - \nabla \log \xi_x \} (dN_{x, x+1} - \gamma dt) \\
&\quad - \gamma \sum_{x \in \mathbb{Z}} (f(x+1-\mu, \rho_k) - f(x-\mu, \rho_k)) \{ 2\nabla \xi_x - \nabla \log \xi_x \} dt.
\end{aligned}$$

Since the compensated Poisson processes $N_{x,x+1}(t) - \gamma t$ are orthogonal martingales with quadratic variation $\gamma^2 t$, then

$$dM_\mu^{(k)} = - \sum_{x \in \mathbb{Z}} (f(x+1-\mu, \rho_k) - f(x-\mu, \rho_k)) \left\{ 2\nabla \xi_x - \nabla \log \xi_x \right\} (dN_{x,x+1} - \gamma dt)$$

defines a martingale with a quadratic variation equal to

$$d\langle M_\mu^{(k)} \rangle_t = \gamma^2 \sum_{x \in \mathbb{Z}} (f(x+1-\mu, \rho_k) - f(x-\mu, \rho_k))^2 \left\{ 2\nabla \xi_x - \nabla \log \xi_x \right\}^2 dt.$$

Using a similar argument to the one in (56), together with the fact that for all $x, y \in \mathbb{Z}$ such that $|x|, |y| \leq C$ it holds that $|x-y|^2 \leq 2C|x-y|$, the boundedness of the function f , (48), (48), (48) and (57), one has that there exists a constant C such that

$$d\langle M_\mu^{(k)} \rangle_t \leq C g(|\mu| + \rho_k(t)) \bar{W}(\xi(t)) \partial_\sigma W(\xi(t), \mu, \rho_k(t)) dt.$$

Similarly we obtain that

$$\left| \sum_{x \in \mathbb{Z}} [f(x+1-\mu, \rho_k(t)) - f(x-\mu, \rho_k(t))] \left\{ 2\nabla \xi_x(t) - \nabla \log \xi_x(t) \right\} \right| \leq C \partial_\sigma W(\xi(t), \mu, \rho_k(t)).$$

Thus, if the constant C_0 is chosen sufficiently large, we have

$$\sup_{t \geq 0} \left\{ W(\xi(t \wedge \tau_k), \mu, \rho_k(t \wedge \tau_k)) \right\} \leq W(\xi(0), \mu, kg(\mu)) + \sup_{t \geq 0} \left\{ N(\mu, k, t) \right\}$$

where $N(\mu, k, t) = M_\mu^{(k)}(t \wedge \tau_k) - \frac{1}{2} \langle M_\mu^{(k)} \rangle_{t \wedge \tau_k}$. Observe that $\exp(M_\mu^{(k)}(t \wedge \tau_k) - \frac{1}{2} \langle M_\mu^{(k)} \rangle_{t \wedge \tau_k})$ is a martingale with expectation equal to 1. By the exponential supermartingale inequality, we have that

$$\mathbf{P} \left(\sup_{t \geq 0} \left\{ N(\mu, k, t) > u \right\} \right) \leq e^{-u}.$$

Thus we proved that for each $k \geq 1$, $\mu \in \mathbb{Z}$ and $u > 0$,

$$\sup_{t \geq 0} \left\{ W(\xi(t \wedge \tau_k), \mu, \rho_k(t \wedge \tau_k)) \right\} \leq W(\xi(0), \mu, kg(\mu)) + u \quad (58)$$

with a probability greater than $1 - e^{-u}$. Applying (58) for each $\mu \in \mathbb{Z}$ and $k \geq 1$ with u replaced by $u + Akg(\mu)$ where $A \geq 1$ is sufficiently large to have $\sum_{k \geq 1} \sum_{\mu \in \mathbb{Z}} e^{-Akg(\mu)} \leq 1$, we obtain

$$\begin{aligned} \sup_{t \geq 0} \left\{ W(\xi(t \wedge \tau_k), \mu, \rho_k(t \wedge \tau_k)) \right\} &\leq W(\xi(0), \mu, kg(\mu)) + Akg(\mu) + u \\ &\leq kg(\mu) \bar{W}(\xi(0)) + Akg(\mu) + u \end{aligned} \quad (59)$$

with a probability greater than $1 - e^{-u}$ uniformly in k and μ .

Define now $k := k_t$, $t \geq 0$ as the smallest integer $k \geq 1$ for which $\rho_k(t) > g(\mu)$; then $\tau_k > t$ and $\rho_k(t) \leq 2g(\mu)$ as $\rho_{k-1}(t) \leq g(\mu)$; thus choosing $k = k_t$ in (59) and using that $W(\xi, \mu, \sigma)$ is increasing in σ (since $\partial_\sigma f \geq 0$ by the conditions imposed on φ), we get

$$\frac{W(\xi(t), \mu, g(\mu))}{g(\mu)} \leq \frac{W(\xi(t), \mu, \rho_k(t))}{g(\mu)} \leq k\bar{W}(\xi(0)) + Ak + \frac{u}{g(\mu)} \leq k\bar{W}(\xi(0)) + Ak + u,$$

where in the last inequality we used the fact that $g(x) \geq 1$ for all $x \in \mathbb{R}$. Taking the supremum over μ and using Lemma 7, we obtain

$$\bar{W}(\xi(t)) \leq Ck_t\bar{W}(\xi(0)) + u$$

for each $t \geq 0$ with probability at least $1 - e^{-u}$. On the other hand,

$$2g(\mu) \geq k_t g(\mu) - C_0 \int_0^t g(|\mu| + |\rho_k(s)|) Z'(s) ds$$

whence

$$k_t \leq 2 + C_0 \int_0^t \frac{g(|\mu| + |\rho_k(s)|)}{g(\mu)} Z'(s) ds.$$

Since $\rho_k(s) \leq k_t g(\mu)$ for any $s \in [0, t]$ and g is increasing, we have that $g(|\mu| + |\rho_k(s)|) \leq g(\mu + k_t g(\mu))$. On the other hand for $x \geq 2$, $g(x) \leq x$ together with the fact that for $x, y \in \mathbb{R}$ $g(|x||y|) \leq g(|x|)g(|y|)$ and since $g(1+x) \leq 1+g(x)$ for $x \geq 1$, we obtain that $g(\mu + k_t g(\mu)) \leq g(\mu)(1+g(k_t))$. As a consequence we obtain that

$$k_t \leq 2 + C_0 Z(t)(1+g(k_t)). \quad (60)$$

Since for all $x \geq 1$ we have that $g(x) \leq 1 + 2\sqrt{|x|}$, then

$$k_t \leq 2 + C_0 Z(t)(2 + 2\sqrt{k_t}).$$

Finally, it follows that $\sqrt{k_t} \leq 2 + 4C_0 Z(t)$. Then, since g is increasing and by plugging the previous inequality in (60), we obtain that

$$k_t \leq 2 + C_0 Z(t)(1 + g((2 + 4C_0 Z(t))^2)).$$

Recalling that $Z'(t) = \bar{W}(\xi(t))$ we obtained that there exists a constant $M > 0$ depending only on λ such that for any $w \geq 1$ and any initial condition $\xi(0)$ satisfying $\bar{W}(\xi(0)) \leq w$,

$$\mathbf{P} \left[\sup_{t \geq 0} \left\{ M^{-1} Z'(t) - w(1 + Z(t)g(Z(t))) \right\} \leq u \right] \geq 1 - e^{-u}.$$

The a priori bound follows from this last inequality (see [13], Proposition 1).

Acknowledgements

The authors are very grateful to József Fritz for illuminating discussions on the existence of the infinite dynamics. We acknowledge the support of the French Ministry of Education through the grant ANR-10-BLAN 0108 (SHEPI). We are grateful to Égide and FCT for the research project FCT/1560/25/1/2012/S. We are grateful to FCT (Portugal) for support through the research project PTDC/MAT/109844/2009. PG thanks the Research Centre of Mathematics of the University of Minho, for the financial support provided by "FEDER" through the "Programa Operacional Factores de Competitividade COMPETE" and by FCT through the research project PEst-C/MAT/UI0013/2011. PG thanks the warm hospitality of "Courant Institute of Mathematical Sciences", where part of this work was done.

References

1. G. Basile, C. Bernardin, S. Olla, Momentum conserving model with anomalous thermal conductivity in low dimensional systems, *Phys. Rev. Lett.* **96** (2006), 204303.
2. G. Basile, C. Bernardin, S. Olla, Thermal conductivity for a momentum conserving model, *Comm. Math. Phys.* **287** (2009), no. 1, 67–98.
3. G. Basile, L. Delfini, S. Lepri, R. Livi, S. Olla, A. Politi, Anomalous transport and relaxation in classical one-dimensional models, *European Journal of Physics Special Topics*, **151** (2007).
4. C. Bernardin, Fluctuations in the occupation time of a site in the asymmetric simple exclusion process, *Ann. Probab.* **32** (2004), no. 1B, 855–879.
5. C. Bernardin, Superdiffusivity of asymmetric energy model in dimensions 1 and 2, *J. Math. Phys.* **49**, (2008), 103301.
6. C. Bernardin and S. Olla, Transport Properties of a Chain of Anharmonic Oscillators with random flip of velocities, *J. Stat. Phys.* **145**, (2011).
7. C. Bernardin and G. Stoltz, Anomalous diffusion for a class of systems with two conserved quantities, *Nonlinearity* **25**, Num. 4, 10991133.
8. A. Borodin and I. Corwin, Macdonald processes, [arXiv:1111.4408v4](https://arxiv.org/abs/1111.4408v4), (2013).
9. F. Bonetto, J.L. Lebowitz, Rey-Bellet, Fourier's law: A challenge to theorists, in *Mathematical Physics 2000*, A. Fokas et al. (eds.), Imperial College Press, London, (2000), 128–150.
10. C. Chang, C. Landim and S. Olla, Equilibrium Fluctuations of asymmetric simple exclusion processes in dimension $d \geq 3$, *Probability Theory and Related Fields* **119**, Number 3 (2001), 81–409.
11. A. Dhar, Heat Transport in low-dimensional systems, *Adv. Phys.*, **57** (2008), 457.
12. R.L. Dobrushin and J. Fritz, Non-equilibrium dynamics of one-dimensional infinite particle systems, *Commun. Math. Phys.*, **55** (1977), 275–292.
13. J. Fritz, Stochastic dynamics of two-dimensional infinite particle systems, *J. Statist. Phys.* **20** (1979), 351–359.
14. J. Fritz, Some remarks on nonequilibrium dynamics of infinite particle systems, *J. Statist. Phys.* **34** (1984), Number 3-4, 539–556.
15. J. Fritz, Gradient dynamics of infinite point systems, *Ann. Prob.* **15** (1987), 478–514.
16. J. Fritz, Entropy pairs and compensated compactness for weakly asymmetric systems, *Advanced Studies in Pure Mathematics* **39** (2004), 143–172.
17. J. Fritz, T. Funaki and J.L. Lebowitz, Stationary states of random Hamiltonian systems, *Probab. Theory Related Fields*, **99** (1994), 211–236.
18. P. Gonçalves, Central Limit Theorem for a Tagged Particle in Asymmetric Simple Exclusion, *Stochastic Process and their Applications*, **118** (2008), 474–502.
19. A. Iacobucci, F. Legoll, S. Olla and G. Stoltz, Thermal conductivity of the Toda lattice with conservative noise, *J. Stat. Phys.* **140** (2010), Number 2, 336–348.
20. M. Jara and C. Landim, Non Equilibrium Central Limit Theorem for a Tagged Particle in Symmetric Simple Exclusion. *Annals Inst. H. Poincaré (B) Probab. and Statist.* **42** (2006), Number 5, 567–577.

21. M. Kac and P. van Moerbeke, On an explicitly soluble system of nonlinear differential equations related to certain Toda lattices, *Adv. Math.* **16** (1975), 160–169.
22. C. Kipnis and C. Landim, *Scaling limits of interacting particle systems*, Springer-Verlag, Berlin, (1999).
23. C. Landim, C., J. Quastel, M. Salmhofer and H.T. Yau, Superdiffusivity of asymmetric exclusion process in dimensions one and two. *Comm. Math. Phys.* **244** (2004), Number 3, 455–481.
24. S. Lepri, R. Livi, A. Politi, Thermal Conduction in classical low-dimensional lattices, *Phys. Rep.* **377** (2003), 1–80.
25. M. E. Vares and H. Rost, Hydrodynamics of a One-Dimensional Nearest Neighbor Model, *AMS Contemporary Mathematics*, **41** (1985), 329–342.
26. T. Sasamoto and H. Spohn, Superdiffusivity of the 1D Lattice Kardar-Parisi-Zhang Equation, *J. Stat. Phys.*, **137** Number 5-6 (2009), 917–935.
27. S. Sethuraman, Central Limit Theorems for Additive Functionals of the Simple Exclusion Process, *Ann. Probab.*, **28** (2000), 277–302.
28. H. Spohn, *Large Scale Dynamics of Interacting Particles*, Springer (1991).
29. H. Spohn, KPZ Scaling Theory and the Semi-discrete Directed Polymer Model, [arXiv: 1201.0645v1](https://arxiv.org/abs/1201.0645v1), (2012).
30. P. Tarrès, B. Tóth, and B. Valkó, Diffusivity bounds for 1D Brownian polymers, *Ann. Probab.*, **40**, Number 2 (2012), 695–713.
31. H. van Beijeren, Exact results for anomalous transport in one dimensional Hamiltonian systems, *Phy. Rev. Let.*, **28** (2012).
32. X. Zotos, Ballistic transport in classical and quantum integrable systems, *Journal of Low Temperature Physics*, **126** (2002), p. 1185.